

WP 2.2

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common 'ingredients' of transport models:

 \bullet . . .

- •dispersion relation
- •linear mode structures (parallel, radial)
- •linear growth/damping rates
- •orbit/zonal averages over mode structures
- •nl evolution/ saturation amplitudes

(i) the solution of the linearized Eqs. (2) and (3); (ii) the solution of the NLSE-like equation for the nonlinear envelope equations, Eq. (10); (iii) the solution of Eqs. (20), (23) and (25) for the particle response averaged over linear parallel mode structures.

depend on: equilibrium kinetic profiles distribution

can be calculated with different fidelity vs speed

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improvements/extensions of LIGKA are subject of WP 2.2

WP2.2-M3: Generalize fast analytical LIGKA version to non-Maxwellian distribution functions, in particular slowing

WP2.2-M1: Develop (semi-)analytical trapped particle model for LIGKA: end 2022

WP2.2-M2: Test and tune analytical global mode structure model for LIGKA/HAGIS end 2022

down end 2023

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Keeping the *m ±* 1-sidebands, retaining the geodesic curvature and the sound wave

|00 circulating ion appoximation; extension trapped particles [I. Chavdarovski et al, 2014…]

$U(P)$ $\sum_{i=1}^{n}$ expansion in the poloidal angle but keeps the full resonances, i.e. is valid for low *q*. WP2.2-M1: trapped particles:

 x^2 cal re **Tes** ults are available (1 I **H** *F* α , α numerical results are available (LIGKA):

- $\frac{1}{2}$ for the complete in a complete in a complete $\mathcal{L}_{\mathcal{U}}$ $\mathsf{particles}$ is most efficient? •which expression for barely/deeply trapped
- •approximations for wprec and wb,t? $\left|\int_{\mathbb{R}^2}$ was pointed in $\left|\int_{\mathbb{R}^2}$ $\begin{array}{ccc} 1 & 1 \\ \hline \end{array}$
- The aim of this paper is the following: in the first part of the first part of \mathcal{L} •which coordinates? map?
- •when to be neglected? error?
- •plans for DAEPS? Common module?

$$
h = -e_a \sum_m \sum_k \underbrace{\frac{\partial F_0}{\partial E} (\omega - \hat{\omega}_*) e^{-im\theta} J_0}_{=\mathcal{R}_{m,k}} \left[a_{km} \phi_m(r) - (a_{km} - \frac{a_{km}^G \bar{\omega}_d(r)}{\omega} + \frac{a_{km}^G \bar{\omega}_d(r)}{\omega} \right]
$$

coupling by an appropriate approximation of the propagator integrals, leads to: [Zonca PPCF 1996,2009, Gotit lectures, Garbet 2006], [Lauber PPCF 2009]

$$
\sum_{m} \omega^{2} \left(1 - \frac{\omega_{*p}}{\omega}\right) - k_{\parallel}^{2} \omega_{A}^{2} R_{0}^{2} = 2 \frac{v_{thi}^{2}}{R_{0}^{2}} \left(-\left[H(x_{m-1}) + H(x_{m+1})\right] + \frac{N^{m}(x_{m-1})N^{m+1}(x_{m+1})}{D^{m-1}(x_{m-1})} + \frac{N^{m}(x_{m+1})N^{m+1}(x_{m+1})}{D^{m+1}(x_{m+1})}\right)
$$

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Keeping the *m ±* 1-sidebands, retaining the geodesic curvature and the sound wave

|00 circulating ion appoximation; extension trapped particles [I. Chavdarovski et al, 2014…]

x 1 mumerical re **Tes** $x^2 + 1$ $x^2 + 2$ $x^3 + 2$ $x^2 + 2$ $x^2 + 2$ $x^3 + 2$ $x^2 + 2$ $x^2 + 2$ $x^3 + 2$ $x^2 + 2$ $x^2 + 2$ $x^3 + 2$ $x^2 + 2$ $x^2 + 2$ $x^3 + 2$ $x^2 + 2$ $x^2 + 2$ $x^3 + 2$ $x^2 + 2$ $x^2 + 2$ $x^2 + 2$ $x^3 + 2$ $x^2 + 2$ $x^2 + 2$ $x^3 +$ **H** *F* α , α numerical results are available (LIGKA):

- $\frac{1}{2}$ function. Although obtained in a completely dialer was a complete way of the complete similar way. particles is most efficient?

and the asymmetry in the •which expression for barely/deeply trapped
- •approximations for wprec and ω b,t? $\begin{array}{ccc} 1 & 1 \\ \hline \end{array}$ •which coordinates? map?
- •when to be neglected? error? UQ? when to be neglected? error? UQ?
- •plans for DAEPS? Common module?

$U(P)$ $\sum_{i=1}^{n}$ expansion in the poloidal angle but keeps the full resonances, i.e. is valid for low *q*. WP2.2-M1: trapped particles:

$$
h = -e_a \sum_m \sum_k \underbrace{\frac{\partial F_0}{\partial E} (\omega - \hat{\omega}_*) e^{-im\theta} J_0}_{=\mathcal{R}_{m,k}} \left[a_{km} \phi_m(r) - (a_{km} - \frac{a_{km}^G \bar{\omega}_d(r)}{\omega} + \frac{a_{km}^G \bar{\omega}_d(r)}{\omega} \right]
$$

$$
\sum_{m} \omega^{2} \left(1 - \frac{\omega_{*p}}{\omega} \right) - k_{\parallel}^{2} \omega_{A}^{2} R_{0}^{2} = 2 \frac{v_{thi}^{2}}{R_{0}^{2}} \left(- \left[H(x_{m-1}) + H(x_{m+1}) \right] \right)
$$

$$
\tau \left[\frac{N^{m} (x_{m-1}) N^{m-1} (x_{m-1})}{D^{m-1} (x_{m-1})} + \frac{N^{m} (x_{m+1}) N^{m+1} (x_{m+1})}{D^{m+1} (x_{m+1})} \right]
$$

 [Zonca PPCF 1996,2009, Gotit lectures, Garbet 2006], [Lauber PPCF 2009]

motivation:

•analytical FLR and FOW expressions need kr - accurate value improves damping/growth rates •decide, if local or global analysis for particular mode number is needed • direct evaluation of kinetic integrals if e.g. Gaussians are assumed (Qualikiz) •input can be used as initial conditions for initial value codes •available literature on e.g. TAEs: depends on shear, alignment of gaps (radially/frequency); some twists when allowing odd/even coupling along the gap •BAEs, RSAEs straightforward, also EAEs possible

test and verify for many scenarios/time points (WF-LIGKA)

FIG. 10. Visualization of $Z(\zeta)$ and $Z'(\zeta)$ with input function F_{SD} for $v_t = 1$ and $v_c = 4$.

• extension to non-Maxwellian distribution functions according to [Hua-Sheng Xie, PoP 2013], bump-on-tail:

> -400 -600 2

 > 0

- •Rabbit: can Legendre polynomial representation be exploited for (partial) analytical integration? (not COM, however…)
- •LIGKA denominator is expanded in rational polynomials analytical integration?

WP2.2-M3: non-Maxwellian distribution functions

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•define a good test cases for benchmark/comparison with DAEPS •define where shared development of model/implementation is possible •test speed vs accuracy, UQ •interface to transport models - IMAS updates/upgrades

WP 3.3

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common 'ingredients' of transport models:

- •dispersion relation
- •linear mode structures (parallel, radial)
- •linear growth/damping rates
- •orbit/zonal averages over mode structures
- •nl evolution/ saturation amplitudes
- •calculate EP fluxes

 \bullet . . .

depend on: equilibrium kinetic profiles distribution

can be calculated with different fidelity vs speed

M1: Extend unperturbed orbit integration routines and averaging procedures in order to calculate phase space fluxes in HAGIS

M2: Explore methodology and possibly implement RABBIT as EP source into HAGIS

M3: Finish reduced EP transport workflow based in LIGKA/HAGIS within IMAS

Milestones WP 3.3

various diagnostics in HAGIS (thx Guo) are available

EP fluxes based on EP energy exchange:

- •separation wrt particle species (passing/trapped)
- •test particle analysis; global/local
- •extend to multi-resonance/mode cases

•connection to WP 3.1

grals along the markers orbit. These dimensionless expressions are of the n: (1/tau) integral (f(n_tor,m_pol,k,eq,...) dt) from time - tau to time, where is the transit/trapping time of the marker and $f()$ a dimensionless function use factor,drift,etc) of the equilibrium $(e.g. q)$ and perturbation (Fourier nonics n_tor,m_pol and bounce harmonic k) along the particles orbits. In the integrals are taken during the last orbit of each marker at the time le of the time node below

rals/quantities along the markers orbit. These dimensionless expressions are form: $(1/tau)$ integral $(f(eq)$ dt) from time - tau to time_orbit for different values of time_orbit in the interval from time - tau to time, where tau is the transit/ trapping time of the marker and f(eq) a dimensionless function (phase, drift,q,etc) of the equilibrium along the markers orbits. The integrals are taken during the last orbit of each marker at the time value of the time node below

call hagis2(equilibrium_in=equilibrium_ingf,distribution_in= distribution_ingf,distribution_out=distribution_outgf,code_parameters_buffer= buffer_hagis2)

to be done: add perturbations - kick-model equivalent information

- IDS data model: https://sharepoint.iter.org/departments/POP/CM/IMDesign/ Data%20Model/CI/imas-3.32.0/html_documentation.html
- LIGKA: https://git.iter.org/projects/STAB/repos/ligka/commits
- HAGIS: https://git.iter.org/projects/STAB/repos/hagis/browse
- EP-WF: https://git.iter.org/projects/WF/repos/ep-stability-wf/browse
- short documentation EP WF: https://confluence.iter.org/pages/viewpage.action? pageId=289069024

9. Use the first link to keep track of the working versions of each dependency (most

Figure 1: Energetic Particle Stability Workflow general layout of the components.

Now the example that we will use is a MPI actor (mode 4 of LIGKA): Before the actor

Figure 2: Example of a typical actor inside a WF.

An example of a working FC2K is the ligka actor: load modules from EP WF by following the tutorial in the confluence page. Then clone ligka and in root of the dir fc2k command. Then open the file named ligka_WF-PY.xml and check out the parameters/compare them with the ones in the documentation.

resources: IMAS, ACTORS,examples [thx. Alin!]

Resources

Some of the resources I found useful when building the Python version of the WF

- 1. https://confluence.iter.org/display/IMP/Integrated+Modelling+Home+Page -> for keeping track of new version of IMAS/PyAL/FC2K (very important!!)
	- (a) https://jira.iter.org/projects/IMAS?selectedItem=com.atlassian.jira.jira-projects plugin:release-pagestatus=released -> IMAS dictionary changes
	- (b) https://confluence.iter.org/display/IMP/Access+Layer -> HDF5 or MDS+ backend for Python
- 2. https://user.iter.org/?uid=YSQENWaction=get_document -> Backend functions documentation for retrieving/manipulating/storing data (Not only Python but also Fortran, C++ and Java)
- 3. https://docs.psnc.pl/display/WFMS/FC2K+Python+wrapper+redesign -> FC2K actor wrapper design (useful for calling an actor after being wrapped by python)
- 4. https://confluence.iter.org/display/IMP/iWrap+Python+Actor->how to build a pytho actor
- 5. https://confluence.iter.org/display/IMP/4.1+FC2K+Basics->smallFC2K tutorial for kepler, but the same can be used for Python (just select the python generation)
- 6. https://confluence.iter.org/display/IMP/3.2+Fortran+examples -> 4 examples of Fortran code with IDSs
- 7. https://confluence.iter.org/display/IMP/iWrap+-+Fortran+API -> Fortran API (can be used with FC2K to generate an actor that can be used in python wf)
- 8. https://confluence.iter.org/pages/viewpage.action?pageId=289069024 -> working example of the EP WF.

of them do not have backward compatibility!!)

2 Example

In order to be able to connect the numerical tools with IMAS and to be able to perform time-dependent analysis on any scenario, Energetic Particle Stability Workflow was created. This is the first time-dependent workflow which uses IMAS infrastructure to perform Energetic particle analysis. It is written in Python and makes use also of a simple interface which makes parameter configuration easy for both the connection to the IMAS Database (for saving/retrieving data) and for the numerical codes themselves through a series of XML files. A general layout of the components that the workflow uses can be seen in Fig.1.

-
- •define where shared development of model/implementation is possible
- •further collaboration with WP 3.1
- •define connection to transport code IDSs

•implement explicit expressions needed for PSZS model (probably as HAGIS module?)

additional slides

dedicated discussion on DEAPS/LIGKA benchmark/common development dedicated meeting on IMAS in 1-2 weeks? dedicated meeting on role of RABBIT? aligned discussion with TSVV on experimental cases (JET DT)

next general meeting? end of July?

combination of models lead to a staged approach for automated analysisfusion

-
-
-

interactive, gui, no-gui versions available [V.-A. Popa] (C) EUROfusion

linear model equations containing crucial effects for self-**EUROfusion** consistent description of EP driven modes:

reduced MHD as limit

$$
\left[e^{-im\theta} \cdot (\omega - \overline{\omega_*^T} J_0 \cdot [\phi_m(r') - (1 - \frac{\omega_d(\theta')}{\omega}) \psi_m(r')\right]
$$

 μ er energy

quasi-neutrality:

$$
\mathbf{E} = -\nabla \phi - \frac{\partial \mathbf{A}}{\partial t}; \qquad A_{\parallel} = \frac{1}{i\omega} (\nabla \psi)_{\parallel}
$$

$$
-\frac{\partial}{\partial t} \left[\nabla \cdot \left(\frac{1}{v_A^2} \nabla \bot \phi \right) \right] + (\mathbf{B} \cdot \nabla) \frac{\nabla \times \nabla \times \frac{1}{\partial t}}{B^2}
$$

$$
= \left[-\sum_a \mu_0 \int d^3 v (e \mathbf{v}_d \cdot \nabla J_0 f)_a \right] + \frac{3}{4}
$$

gyrokinetic moment equation: shear Alfven law

'pressure' tensor - curvature drift coupling

originally based on H. Qin, PhD Thesis 1998, PoP 1999
 $E_{\ell} = 0$

coefficients can be calculated using the finder programme, included in LIGKA git repository

upgrade analytical expression building on [I. Chavdarovski et al, 2014…]

^S ⁼ ⇣ !*^d* \mathbf{I} ⌘2 \cdot ⁿ \cdot *^D*^ˆ ⁺ *^P ^G*%^ˆ FOW dispersion relation for LIGKA

- equivalent to EGAM FOW equations: Qiu [2009], Miki & Idomura [2015]
-
- LIGKA mode (3/4 , specification of kr needed)
- rationale: implement global effects in local model can be improved by estimating analytically AE mode structures (ongoing…)

$$
\begin{array}{l|l|l} \mathbf{w}^2\big(1-\frac{\omega_{\ast p}}{\omega}\big)-k_{\parallel}^2\omega_{A}^2R_0^2=2\frac{v_{thi}^2}{R_0^2}\bigg(-\big[H(x_{m-1})+H(x_{m+1})\bigg]+\\ &\frac{\rho^2\bigg[D_{0,0}\big[D_{-2,-2}D_{1,1}D_{-1,-1}\big(D_{2,2}\big(D_{-1,-1}\big(Q_{1,0}N^C_{0,1}+N_{1,0}Q^C_{0,1}\big)-\\ &\frac{F_{-1,1}\big(N_{1,0}N^C_{0,-1}-N_{-1,0}N^C_{0,1}\big)\big)}{D^{m-1}(x_{m-1})}+\frac{N^m(x_{m+1})N^{m+1}(x_{m+1})}{D^{m+1}(x_{m+1})}\bigg]\bigg)\\ &\frac{-D_{-1,-1}\big(E_{1,2}P_{2,0}N^C_{0,-1}-E_{-1,0}N^C_{0,1}\big)}{D_{1,1}\big(D_{2,2}\big(E_{-2,-1}N_{-1,0}\big(E_{-1,-2}N^C_{0,-1}-D_{-1,-1}P^C_{0,-2}\big)+\\ &\frac{D_{-1,-1}P_{2,0}\big(D_{-1,-1}P^C_{0,-2}-E_{-1,-2}N^C_{0,-1}-D_{-1,-1}P^C_{0,-2}\big)}{D_{-1,-1}P_{2,0}\big(D_{-1,-1}P^C_{0,-2}-E_{-1,-2}N^C_{0,-1}\big)-F_{-1,-1}N_{-1,0}N^C_{0,-1}\big)}+\frac{D_{-1,-1}P_{2,0}\big(D_{-1,-1}P^C_{0,-2}-E_{-1,-2}N^C_{0,-1}\big)}{D_{-2,-2}\big(D_{-1,-1}\big(Q_{-1,0}N^C_{0,-1}+N_{-1,0}Q^C_{0,-1}\big)-F_{-1,-1}N_{-1,0}N^C_{0,-1}\big)}\big)+\\ &\frac{D_{-2,-2}D_{2,1-1}^2P_{2,0}P^C_{0,2}\big)+D_{-2,-2}D_{2,1-1}^2N_{1,0}N^C_{0,1}\big)}{D_{-1,-1}E_{0,1}N_{1,0}\big(D_{1,1
$$

$$
\omega^2 \left(1 - \frac{\omega_{*p}}{\omega} \right) - k_{\parallel}^2 \omega_A^2 R_0^2 = 2 \frac{v_{thi}^2}{R_0^2} \left(- \left[H(x_{m-1}) + H(x_{m+1}) \right] + \frac{N^m (x_{m-1}) N^{m-1} (x_{m-1})}{D^{m-1} (x_{m-1})} + \frac{N^m (x_{m+1}) N^{m+1} (x_{m+1})}{D^{m+1} (x_{m+1})} \right)
$$

!)

thi = 2*Ti/mi*, !*^p*⇤ = !⇤*ⁿ* + !⇤*^T* = *Ti/*(*eB*)*k*✓(r*n/n*)(1 + ⌘) with 1+*D*˜(*xe,m*) i +⌧ h 1+*D*˜(*xi,m*) , *^N ^m*(*xm*) = *^N*˜ *^m*(*xi,m*)*N*˜ *^m*(*xe,m*), !)*xZ*(*x*) !⇤ ! ⌘ ⇣ *^x*² ⁺ *xZ*(*x*)(*x*² ¹ 2) ⌘ , 2*N*˜ *^m*(*x*) = (1 !*^m x*⁴ + *^x*² ²) + *xZ*(*x*)(¹ ⁴ + *x*⁴) , *H*(*xm*) = *H*˜ (*xm,i*) + ⌧*H*˜ (*xm,e*), !)*F*˜(*xm*) ⌘ !⇤ *G*˜(*xm*) i , 2*F*˜(*x*) = *xZ*(*x*)(¹ no FOW, circulating particle approximation

 A and A a

[Zonca 1996,2009 Lauber 2009]

 $c_{\rm eff}$ to the ballooning formulation results) to the asymmetry in the asymmetry in the asymmetry in the asymmetry in the \sim

² ⁺ *^x*² ⁺ *^x*⁴) + ³*x*²

The diagonal matrix elements are $\frac{1}{2}$ and $\frac{1}{2}$ and $\frac{1}{2}$ are $\frac{1}{2}$ and $\frac{1}{2}$ and $\frac{1}{2$

[Zonca 1998, Z.X. Lu 2017,Lauber JPC 2018]

• fast analytical model for FOW effects: solve equations both locally (scan k_r) and globally

2nd order FOW

analytical values for ω_{t} , ω_{D} , propagator coefficie

analytical expression/result: $a_{k,\sigma}^G = \frac{1}{\tau_t \omega} \int_0^{\tau_t} dt \Big[\omega_d^{\theta} \cos(\theta) + \frac{\omega_d^r}{i} \sin(\theta) +$ geometrical effects 1.0 $_{\odot}$ 0.5 0 16 16 16 17 19 $\overline{16}$ $0.2 \rightarrow 0.4 \rightarrow 0.6 \rightarrow 0.8$ 8 Λ

so far: circulating ion expression used; now: use pre-calculated HAGIS orbits to replace analytical values for ω_t , ω_D , propagator coefficients,...

$$
\left. \cdot \hspace{0.3cm} \omega_{prec} \right] \hspace{0.2cm} e^{i \sigma k \omega_t t} = \frac{\delta_{k,\pm 1}}{2 \omega} (\omega_d^{\theta} \hspace{-0.05cm} \mp \hspace{-0.05cm} \omega_d^r) + \delta_k \omega_{prec}
$$

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dependence on pitch angle, substantial deviations for low frequencies

with

 $\overline{}$ *x*) $\overline{}$ $+$

$$
2\tilde{N}^m(x) = \left(1 - \frac{\omega_*^m}{\omega}\right) \left[x^2 + xZ(x)(x^2 + \frac{1}{2})\right] - \frac{\omega_*^m}{\omega} \eta \left[x^2 + \frac{1}{2}\right]
$$

⇤⌥⌥⇧

P^m

⇤

N

m

/

 $\overline{}$

0

\overline{d} $\overline{\mathsf{in}}$ definition definitions

⇧

with
\n
$$
\tilde{D}^m(x) = (1 - \frac{\omega_*^m}{\omega})xZ(x) - \frac{\omega_*^m}{\omega} \eta \left(x^2 + xZ(x)(x^2 - \frac{1}{2}) \right)
$$
\n
$$
2\tilde{N}^m(x) = (1 - \frac{\omega_*^m}{\omega}) \left[x^2 + xZ(x)(x^2 + \frac{1}{2}) \right] - \frac{\omega_*^m}{\omega} \eta \left[x^2(x^2 + \frac{1}{2}) + xZ(x)(\frac{1}{4} + x^4) \right]
$$

 $\mathop{\text{in}}$ F_{α} with $\partial F_{\alpha}/\partial F_{\alpha} = -F_{\alpha}/T$ and using a maxwellial I'_0 with $OI'_0/U = -I'_0/I$ and using ⌃²*T^a m,p*⌦ *nd rd nd rd* Assuming a Maxwellian F_0 with $\partial F_0 / \partial E =$ $\overline{}$ *F*0*/T* and using ⇧cients *ak,m,* ⌅ , *a k,m,* ⌅ , *Kk,m,p,*

$$
P = \tau (\Gamma_0 - 1) \Big[1 - \frac{\omega_i^*}{\omega} \Big(1 + \eta_i \frac{\Gamma_0 G_0}{\Gamma_0 - 1} \Big) \Big].
$$
\n
$$
H^m(x_m) = \tilde{H}^m(x_{m,i}) + \tau \tilde{H}^m(x_{m,e}) \text{ and } \tilde{H}^m(x) = \frac{1}{2} \Big[(1 - \frac{\omega_i^m}{\omega}) \tilde{F}(x) - \eta \frac{\omega_i^m}{\omega} \tilde{G}(x_{m,i}) + \tau \tilde{H}^m(x_{m,e}) \text{ and } \tilde{H}^m(x) = \frac{1}{2} \Big[(1 - \frac{\omega_i^m}{\omega}) \tilde{F}(x) - \eta \frac{\omega_i^m}{\omega} \tilde{G}(x_{m,i}) + \tau \tilde{H}^m(x_{m,e}) \text{ and } \tilde{H}^m(x) = \frac{1}{2} \Big[(1 - \frac{\omega_i^m}{\omega}) \tilde{F}(x) - \eta \frac{\omega_i^m}{\omega} \tilde{G}(x_{m,i}) + \tau \tilde{H}^m(x_{m,e}) \text{ and } \tilde{H}^m(x) = \frac{1}{2} \Big[(1 - \frac{\omega_i^m}{\omega}) \tilde{F}(x) - \eta \frac{\omega_i^m}{\omega} \tilde{G}(x_{m,i}) + \tau \tilde{H}^m(x_{m,e}) \text{ and } \tilde{H}^m(x) = \frac{1}{2} \Big[(1 - \frac{\omega_i^m}{\omega}) \tilde{F}(x) - \eta \frac{\omega_i^m}{\omega} \tilde{G}(x_{m,i}) + \tau \tilde{H}^m(x_{m,e}) \text{ and } \tilde{H}^m(x) = \frac{1}{2} \Big[(1 - \frac{\omega_i^m}{\omega}) \tilde{F}(x) - \eta \frac{\omega_i^m}{\omega} \tilde{G}(x_{m,i}) + \tau \tilde{H}^m(x_{m,e}) \text{ and } \tilde{H}^m(x) = \frac{1}{2} \Big[(1 - \frac{\omega_i^m}{\omega}) \tilde{F}(x) - \eta \frac{\omega_i^m}{\omega} \tilde{G}(x_{m,i}) + \tau \tilde{H}^m(x_{m,e}) \text{ and } \tilde{H}^m(x) = \frac{1}{2} \Big[(1
$$

$$
P = \tau (1^{\circ}_{0} - 1) \left[1 - \frac{\pi_{i}}{\omega} \left(1 + \eta_{i} \frac{1000}{\Gamma_{0} - 1} \right) \right].
$$
\n
$$
H^{m}(x_{m}) = \tilde{H}^{m}(x_{m,i}) + \tau \tilde{H}^{m}(x_{m,e}) \text{ and } \tilde{H}^{m}(x) = \frac{1}{2} \left[(1 - \frac{\omega_{*}^{m}}{\omega}) \tilde{F}(x) - \eta \frac{\omega_{*}^{m}}{\omega} \tilde{G}(x) \right]
$$
\n
$$
2\tilde{F}(x) = xZ(x)(\frac{1}{2} + x^{2} + x^{4}) + \frac{3x^{2}}{2} + x^{4},
$$
\n
$$
\omega_{d}^{\pm} \approx \frac{\upsilon_{th,i}^{2}}{\Omega_{i}} \frac{1}{R_{0}} (\frac{m}{r} \pm \frac{\partial}{\partial r}) = \omega_{d}^{n} \pm \omega_{d}^{r}
$$
\n
$$
2\tilde{G}(x) = xZ(x)(\frac{3}{4} + x^{2} + \frac{x^{4}}{2} + x^{6}) + 2x^{2} + x^{4} + x^{6}
$$

$$
\int_0^\infty \frac{dt \ e^{-t^2}}{x_m^2 - t^2} = \frac{-\sqrt{\pi}Z(x_m)}{2x_m}; \qquad \int_0^\infty \frac{dt \ t^2 \ e^{-t^2}}{x_m^2 - t^2} = \frac{-\sqrt{\pi}}{2}(x_m + x_m^2 Z(x_m))
$$

where
$$
x_m = \frac{\omega}{|k_{\parallel,m}|v_{th}}; \qquad t = \frac{v_{\parallel}}{v_{th}}; \qquad v_{th} = \sqrt{\frac{2T}{m}} \qquad \qquad \mathsf{T} = \mathsf{Te}/\mathsf{T}
$$

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 k<mark>-off</mark> me

where

adding Maxwellian αs , T_{α} =1MeV, $v_{th,\alpha}/v_{A0}$ =1.0

kr· $pi \sim 0.05$; kr· $p_{\alpha} \sim 0.3$

reduced EP drive in TAE range - as expected, but BAE is stabilised? 2nd order expansion should be valid till kr · $\rho_{\alpha} \sim 1$ difficult system, since local and global effects are both present via kr

