



Recent SOLEDGE3X improvements and first applications, towards a more realistic modelling of the generation of inverse Er

G. Falchetto

*Acknowledgements to Hugo Bufferand, Xavier Garbet, Patrick Tamain,
and SOLEDGE Team: G. Ciraolo, N. Rivals, H. Yang, Y. Marandet*

- ❑ TSVV1 SOLEDGE3X task and deliverables
- ❑ Progress report:
 - ❑ SOLEDGE3X model
 - Considerations on neoclassical viscosity implementation
 - ❑ Status of the implementation
 - ❑ Application to 3D turbulence with fluid neutrals in TCV-like geometry
 - ❑ Next steps
 - Set up of realistic WEST case

□ Explore the impact on the inverse radial electric field formation of :

- neoclassical friction
- favorable versus unfavorable magnetic drift configuration in realistic X-point geometry

- **M2.7 Implement neoclassical friction in SOLEEDGE3X** fluid edge turbulence code. Investigate its impact on the generation and dynamics of the radial electric field well via a **power scan in a limited case**, in comparison to previous results without the friction (of TSVV1 pilot) and to GK results. Investigation of QH or I mode regime. G. Falchetto **06/2022**
- **D2.5** Report including statements on the **relative impact of some separate ingredients playing a role** in the radial electric field formation (orbit losses, ripple, **turbulence, neutrals..**)
Report or paper submitted, conference contribution X. Garbet, R. Varennes, L. Vermare, G. Falchetto **12/2022**
- **M4.2** Compare the generation of an inversed radial electric field in two magnetic configurations (**favourable vs unfavourable magnetic drift direction**) in **SOLEEDGE3X with realistic X-point geometry and neutrals**, compare to experimental findings **on one machine**. G. Falchetto **06/2023**
- **D4.2** Report on the study of the **effect of the direction of the magnetic drift and the level of realism of the edge conditions**, with respect to experimental measurements
Report, paper, or conference contribution G. Falchetto **12/2023**

- SOLEDGE3X merges SOLEDGE2D and TOKAM3X + relaxing some model assumptions
- **Electrostatic drift-fluid** equation system for a **multi-species plasma** based on Zhdanov closure
[Raghunathan et al., PPCF, 2021]

- **Mass balance** for all ion species – quasi-neutrality for electrons

$$[1] \quad \partial_t n + \vec{\nabla} \cdot (n \vec{v}) = S_n^{iz} \quad n_e = \sum_i Z_i n_i$$

Ionization terms
“Anomalous” diffusion
for transport simulations

- Velocity decomposition $\vec{v} = v_{\parallel} \vec{b} + \vec{v}_{\perp drifts} - \frac{D \vec{\nabla}_{\perp} n}{n}$

- **Ionization/recombination sources** involving neutrals are computed by the “neutral solver”:
 - Either fluid neutrals (diffusive): very crude model but fast and robust
good approximation for recycling source and radiation power losses in the divertor
 - Or kinetic neutrals EIRENE MC code : high fidelity

• Momentum balance

$$[2] \quad \partial_t(mn\vec{v}) + \vec{\nabla} \cdot (mn\vec{v} \otimes \vec{v}) = -\vec{\nabla}p - \vec{\nabla} \cdot \bar{\Pi} + Zen(\vec{E} + \vec{v} \times \vec{B}) + \vec{R} + \vec{S}_{v_{\parallel}}^{iz} + \vec{S}_{v_{\perp}}^{ex}$$

Collisional closure terms
 Ionization terms
 External sources
 “Anomalous” diffusion for transport simulations

Parallel projection ($\vec{b} \cdot [2]$)

$$[3] \quad \partial_t(mnv_{\parallel}) + \vec{\nabla} \cdot (mnv_{\parallel}\vec{v}) = -\nabla_{\parallel}p + ZenE_{\parallel} - \vec{b} \cdot (\vec{\nabla} \cdot \bar{\Pi}) + R_{\parallel} + S_{v_{\parallel}}^{iz+ex} + \vec{\nabla} \cdot (mn\eta_{\perp}\vec{\nabla}_{\perp}v_{\parallel})$$

Perpendicular momentum: ($\vec{b} \times [2]$)

$$\vec{v}_{\perp} = \frac{\vec{E} \times \vec{B}}{B^2} + \frac{\vec{B} \times (\vec{\nabla}p + \vec{\nabla} \cdot \bar{\Pi})}{ZenB^2} - \frac{\vec{B} \times (\vec{R}_{\perp} + \vec{S}_{v_{\perp}}^{iz+ex})}{ZenB^2} + \frac{\vec{b}}{n\omega_c} \times [\partial_t(n\vec{v}) + \vec{\nabla} \cdot (n\vec{v} \otimes \vec{v})]$$

Drift ordering (assuming $\varpi \ll \omega_c$):

$$\vec{v}_{\perp}^0 = \frac{\vec{E} \times \vec{B}}{B^2} + \frac{\vec{B} \times \vec{\nabla}p}{ZenB^2} \qquad \vec{v}_{\perp}^1 = \frac{\vec{b}}{n\omega_c} \times [\partial_t(n\vec{v}^0) + \vec{\nabla} \cdot (n\vec{v}^0 \otimes \vec{v}^0)] + \frac{\vec{B} \times \vec{\nabla} \cdot \bar{\Pi}}{ZenB^2} - \frac{\vec{B} \times (\vec{R}_{\perp} + \vec{S}_{v_{\perp}}^{iz+ex})}{ZenB^2}$$

E x B
Diamagnetic
Polarization
Non-linear drifts (function of \vec{v}^0)

- **Current balance** (Equation on ϕ – vorticity equation)

$$[4] \quad \vec{\nabla} \cdot \vec{j} = 0 \quad \text{with} \quad \vec{j} = \sum Z n \vec{v} = j_{\parallel} \vec{b} + \vec{j}^* + \vec{j}_{\Pi} + \vec{j}_{pol} + \vec{j}^{ex}$$

Parallel current obtained from electron parallel momentum balance
(Generalized Ohm's law)

$$j_{\parallel} = -\sigma_{\parallel} \left(\nabla_{\parallel} \phi + \frac{-\nabla_{\parallel} p_e + R_{e,\parallel}^T}{en_e} \right)$$

At 0th in drift ordering, polarization current can be expressed: $\vec{j}_{pol} = -\partial_t \vec{\omega} - \vec{\nabla} \cdot (\vec{v}^0 \otimes \vec{\omega})$

where:

$$\vec{\omega} = \frac{m_i}{B^2} \left(n \vec{\nabla}_{\perp} \phi + \frac{1}{Ze} \vec{\nabla}_{\perp} p_i \right) \quad \vec{\nabla} \cdot \vec{\omega} = \Omega \quad \text{vorticity}$$

ϕ equation takes the form: $\vec{\nabla} \cdot \left(\partial_t \left[\frac{m_i n}{B^2} \vec{\nabla}_{\perp} \phi \right] + \sigma_{\parallel} \nabla_{\parallel} \phi \vec{b} \right) = RHS$

- **Energy balance** $\mathcal{E} = \frac{3}{2} enT + \frac{1}{2} mnv_{\parallel}^2$

$$[5] \quad \partial_t \mathcal{E} + \vec{\nabla} \cdot (\mathcal{E} \vec{v} + p v_{\parallel} \vec{b} + v_{\parallel} \bar{\Pi} \cdot \vec{b} + \vec{q}) = Z n e v_{\parallel} E_{\parallel} + v_{\parallel} R_{\parallel} + Q + S_{\mathcal{E}}^{iz+ex} + \vec{\nabla} \cdot (m n v_{\parallel} \eta_{\perp} \vec{\nabla}_{\perp} v_{\parallel} + n \chi_{\perp} \vec{\nabla}_{\perp} T)$$

Collisional closure terms
Ionization terms
External sources
"Anomalous" diffusion for transport simulations

In SOLEDGE3X, the stress tensor is treated as follows:

$$\bar{\bar{\Pi}} = \bar{\bar{\Pi}}_{\parallel} + \bar{\bar{\Pi}}_{\perp} + \bar{\bar{\Pi}}_{\perp}$$

- The perpendicular part including perpendicular collision is neglected
- $\bar{\bar{\Pi}}_{\perp}$ containing gyroviscous terms which are linked to diamagnetic effects in the fluid description is taken into account **assuming ideal diamagnetic cancellation**, that is considering the divergence of the diamagnetic stress tensor cancels the advection of velocity by diamagnetic velocity. Whereas the advection by second order " $\nabla \mathbf{B}$ " drift is kept.
- The parallel stress tensor embedding the parallel collisional contribution including parallel ion viscosity effects is kept

$$\bar{\bar{\Pi}} = \pi_{\parallel} \left(\vec{b} \otimes \vec{b} - \frac{1}{3} \bar{\bar{I}} \right) \quad \text{where} \quad \pi_{\parallel} = -3\eta_{\parallel} \left(\nabla_{\parallel} v_{\parallel} - \vec{\kappa} \cdot \vec{v} - \frac{1}{3} \vec{\nabla} \cdot \vec{v} \right)$$

$$\vec{\nabla} \cdot \bar{\bar{\Pi}}_{\parallel} = \left[\vec{\nabla} \cdot (\pi_{\parallel} \vec{b}) \right] \vec{b} + \pi_{\parallel} \vec{\kappa} - \frac{1}{3} \vec{\nabla} \pi_{\parallel}$$

Implementation of neoclassical viscosity

Braginskii model pertinent to high collisionality edge plasma should recover Pfirsch – Schluter regime though neoclassical ordering is different : neoclassical $v \sim \varepsilon v_{thi}$ whereas Braginskii & Zhdanov assume $v \sim v_{thi}$

From Helander Sigmar (2002)

$$\langle \mathbf{B} \cdot \nabla \cdot \overline{\Pi}_a \rangle = 3 \langle (\nabla_{\parallel} B)^2 \rangle \left(\mu_{a1} u_{a\vartheta} + \mu_{a2} \frac{2qa_{\theta}}{5pa} \right) \quad \mu_{a1}, \mu_{a2} \text{ parallel viscosity coefficients}$$

$$= \frac{3}{2} \eta_0^a \langle W_{zz} \nabla_{\parallel} B \rangle$$

In the limit $\rho_L \rightarrow 0$

$$W_{zz} = 2(\nabla_{\parallel} v_{\parallel} - \mathbf{v} \cdot \nabla_{\parallel} \mathbf{b}) = 2 u_{a\vartheta} \nabla_{\parallel} B$$

missing term
in TOKAM3X model

$$\langle \mathbf{B} \cdot \nabla \cdot \overline{\Pi}_a \rangle \equiv 3 \eta_0^a \langle (\nabla_{\parallel} B)^2 \rangle u_{a\vartheta}$$

In Braginskii parallel viscosity is determined by particle flux alone (heat flux is smaller order due to the chosen flow ordering)

$$\mu_{a1} = \eta_0^a \quad \eta_0^i = 0.96 n_i T_i \sqrt{2} \tau_{ii}$$

OPEN QUESTION: Is there a better closure than Braginskii/Zhdanov for multi-species edge plasma without assuming a scaling on velocities?

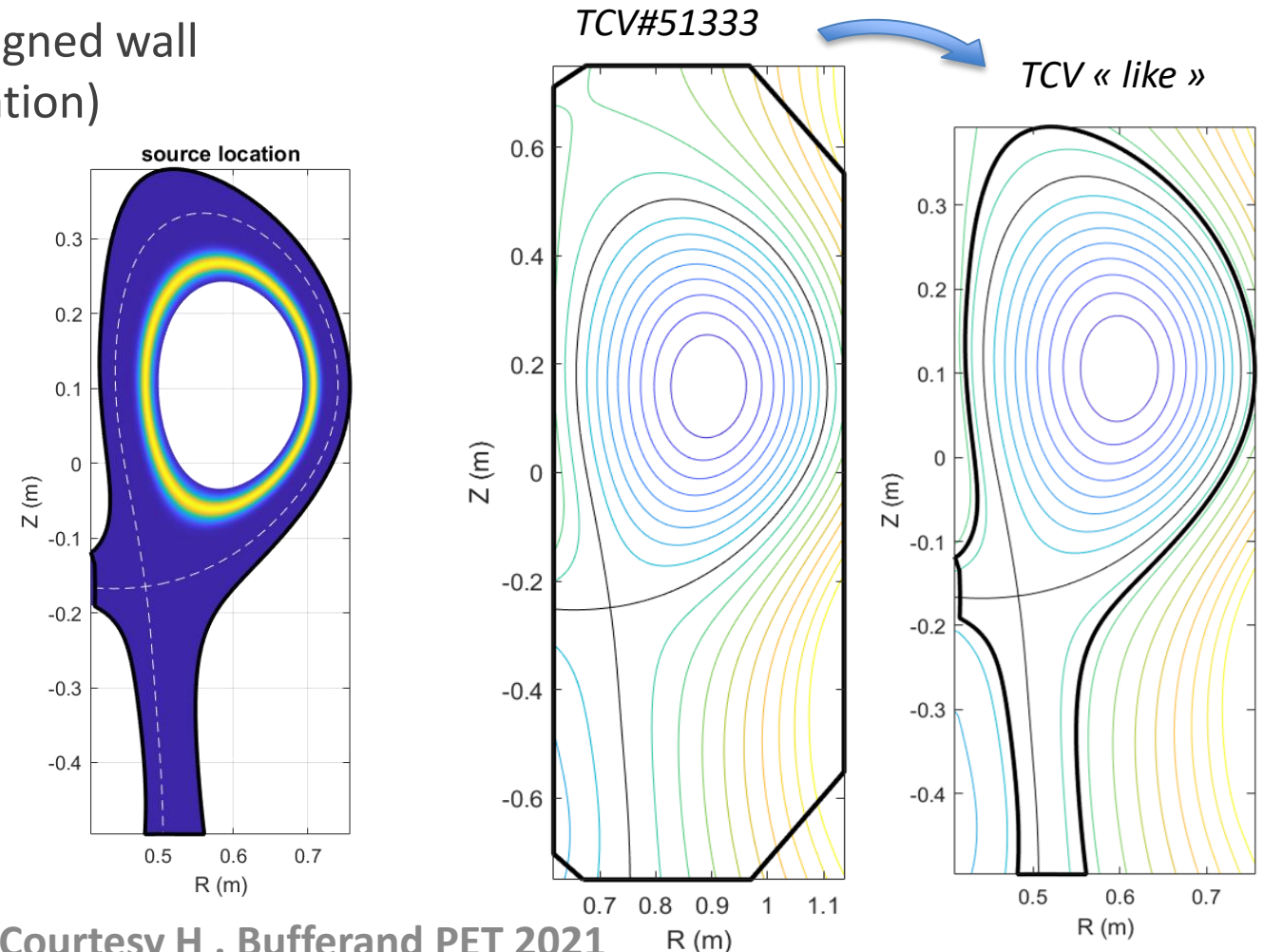
- ✓ New SOLEDGE3X code enables 2D/3D transport and 3D turbulent simulations for multi-species plasma.
- **Status of implementation:**
 - ✓ Code runs routinely in 2D transport mode (including drifts) and 3D turbulent mode in divertor geometry
 - ✓ **Neutrals:** coupling to fluid model tested in 2D and 3D, coupling to EIRENE tested for 2D transport only + advanced A&M model (from [Kotov PPCF 2008](#) w/o n-n collisions) [[Bufferand PET 2021](#), [Rivals PET 2021](#)]
 - ✓ First application to **power scan on TCV case including neutrals** shows onset of turbulence at low power and recovers generation of E_r well accompanied by turbulence reduction around separatrix at high power
 - ✓ Drift and current associated to the stress tensor divergence, including parallel ion viscosity – implemented
under test before merging into SOLEDGE3X git master branch and release
- ✗ **Missing terms:** polarization velocity not included in particle, parallel momentum and energy transport
- ✗ **Electromagnetic** version under development

Purpose: test code capability to simulate a turbulent plasma in **divertor geometry including fluid neutrals**

- Equilibrium based on TCV#51333 (magnetic field rescaled by $\frac{1}{2}$)
- Wall geometry modified to fit a flux surface (to avoid potential artefacts due to a non-aligned wall – consequence: closer divertor in the simulation)

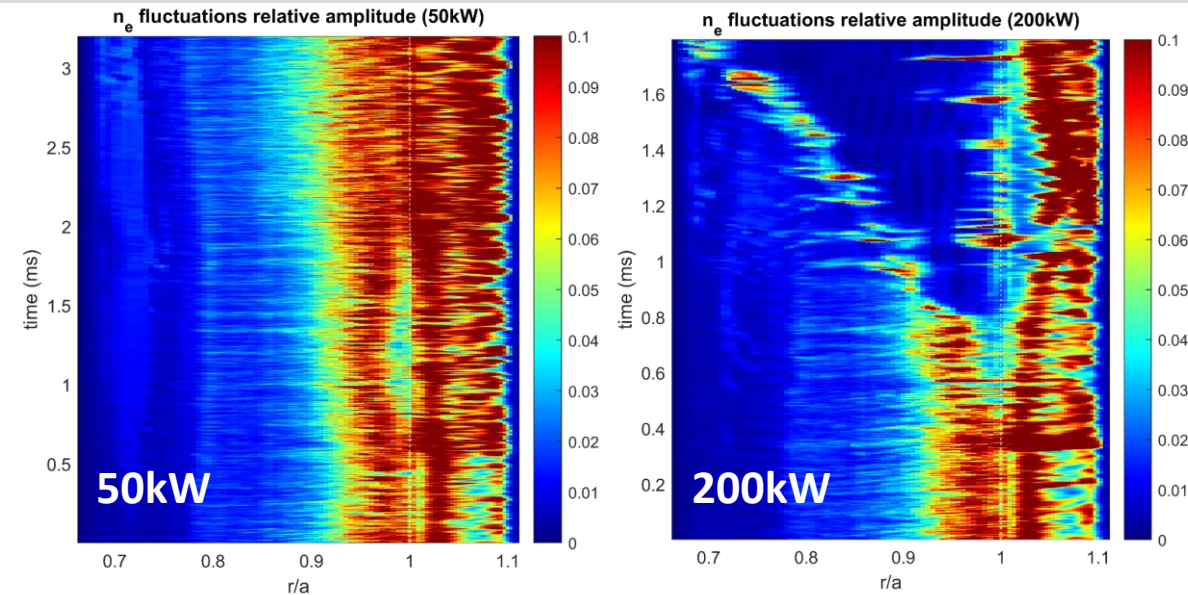
Simulation setup:

- Pure D plasma with **fluid neutrals**
- Anomalous diffusivities set to $10^{-2} m^2 s^{-1}$ (classical collision level)
- Energy source at $\frac{r}{a} \approx 0.75$
Power scan 50 100 200 kW
- Recycling on the wall set to 80%
Gas puff at outer midplane adjusted to get $n_e \approx 10^{19} m^{-3}$ at the separatrix
- Parallel resistivity artificially $\times 10$
 $N_\psi \sim 92$; $N_\theta \sim 736$; $N_\varphi = 64$ (1/4 torus)
grid points $\sim 4 \cdot 10^6$

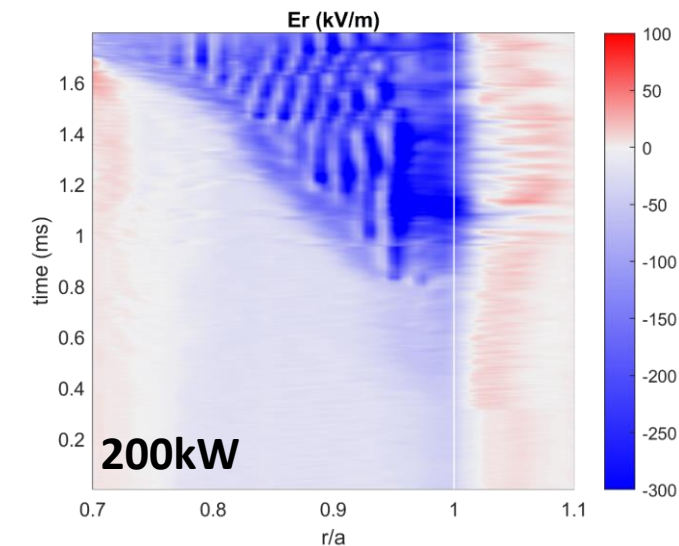
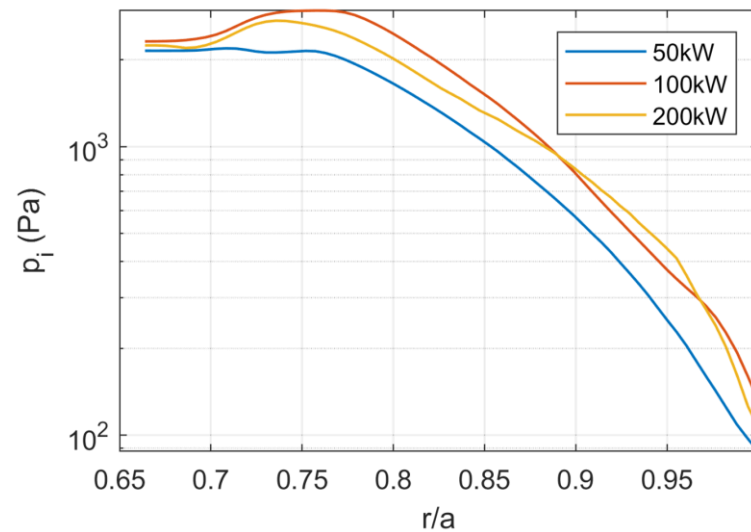


Courtesy H. Bufferand PET 2021

- Well developed turbulence at low power ($50kW$)
Avalanches cross the separatrix
- Reduction of fluctuation level around the separatrix at higher power ($100kW$ and $200kW$)
“gap” propagates from separatrix inward
Associated with higher $E \times B$ shear
Stationary zonal flows observed in the low turbulence region
Similar to [Giacomin, J. Plasma Phys., 2020]



- 1ms after turbulence reduction, no clear steepening in pressure profile
local steepening near separatrix
- E_r well recovered though **very high value**
missing a term to control plasma rotation?
Ion viscosity effects
[Sigmar & Helander, Zholobenko et al., PPCF 2021]



Courtesy H . Bufferand PET 2021

Neoclassical viscosity impact study: plan towards M2.7 and deliverable 2022

- 2D transport simulations – verification of P-S neoclassical transport
- 3D turbulence reference case in circular bottom limiter geometry with $\eta=0$
(same parameters as GYSELA : Tore Supra rescaled with $\frac{1}{2}B$ $\frac{1}{4}$ of torus)
- 3D turbulence simulations in circular bottom limiter case [with parallel ion viscosity](#) - power scan

Magnetic drift impact: plan towards M4.2 and deliverable 2023

- WEST #54903 LSN @t = 4.8s
- Compare to reversed B field case: USN
- Higher q_{95} to verify the effect observed experimentally
- Increase power in USN case

WEST #54903

2D SOLEDGE3X-EIRENE simulation setup

- LSN
Pure D
- Upstream density = $1.6e^{19}$ part/m³
- Input power $P_{in} = P_{heat} - P_{rad}^{core} \simeq 450$ kW
- $D = \nu = 0.3$ m²s⁻¹, $\chi_e = \chi_i = 1$ m²s⁻¹
- Recycling coefficient $R_{wall} = 1$, $R_{pump} = 0.95$

Courtesy M. Peret, H. Yang

