

Recent Advancements in the Gyrokinetic Moment-Based Approach for Boundary Plasmas

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- Motivations and The GyroKinetic (GK) Moment-Based Approach
- Numerical studies and benchmarks with GENE code
- Conclusions

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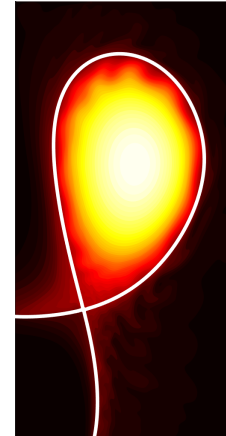
- **Boundary region** (edge and scrape-off-layer) controls the performance of fusion devices (H-L transition, pedestal, ELMs, ...)
- **Boundary region** is characterised by **different plasma collisionality regimes**:

ITER: $T_e \sim 10 - 10^4$ eV and $n \sim 10^{18} - 10^{20}$ m⁻³ $\Rightarrow \lambda_{mpf}/R_0 \sim \frac{T_e^2}{n_e R_0} \sim 10^{-1} - 10^3$

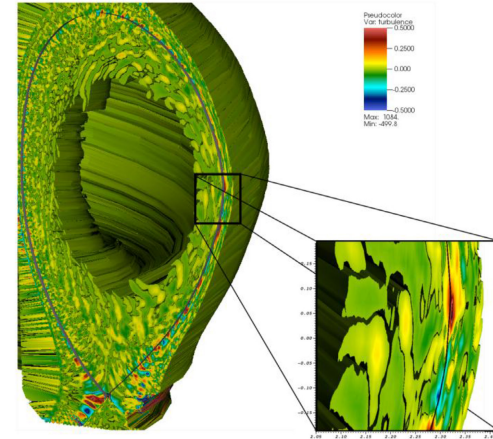
- Two currently-used turbulent modelling approaches:
 - **Drift-Kinetic (DK) fluid modelling** (lowest-order moments, less expensive, $k_{\perp} \rho_i \ll 1$ and $\lambda_{mfp} k_{\parallel} \ll 1$, no kinetic effects, Full-F)
 - **GyroKinetic (GK) modelling** (expensive, $k_{\perp} \rho_i \sim 1$, kinetic effects, also suitable for $\lambda_{mfp} k_{\parallel} \gtrsim 1$)

GK Moment Approach

$$k_{\parallel} \lambda_{mfp} \ll 1 \quad \& \quad k_{\parallel} \lambda_{mfp} \gtrsim 1$$



Giacomin M. *et al.* JPP, 2020

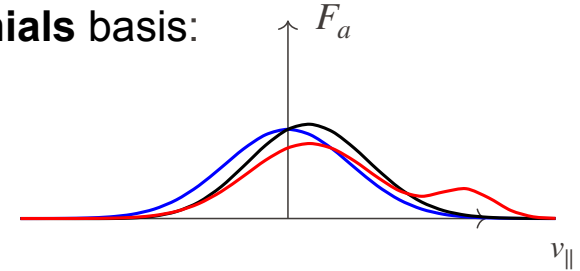


Chang C. S. *et al.* Nucl. Fusion, 2017

- Expansion of F_a (Full-F) on a **Hermite-Laguerre polynomials** basis:

Gyro-Moment of F_a

$$F_a = \sum_p \sum_j \boxed{N_a^{pj}(\mathbf{R}, t)} \frac{H_p(v_{\parallel}) L_j(\mu)}{\sqrt{2^p p!}} F_{Ma}$$



- Equation for N_a^{pj}** from the projection onto the Hermite-Laguerre basis of GK Equ.

$$\|\dots\|^{pj} = \int d\theta \int d\mu dv_{\parallel} H_p(v_{\parallel}) L_j(\mu) \dots \quad \|(\text{GK Equ.})\|^{pj}$$

- Gyro-Moment Hierarchy Equation and GK Field Equations**

$$\frac{\partial}{\partial t} N_a^{pj} + \nabla \cdot \|\dot{\mathbf{R}}\|^{pj} - \frac{\sqrt{2l}}{v_{th\|a}} \|\dot{v}_{\parallel}\|^{p-1j} + \mathcal{F}_a^{pj} = \sum_b C_{ab}^{pj} \quad C_{ab}^{pj} = \|\langle C_{ab} \rangle\|^{pj}$$

Projection of GK collision operator

EPFL Nonlinear GK Coulomb in the Gyro-Moment Approach 6

- In (\mathbf{x}, \mathbf{v}) phase-space, **Coulomb** (“Landau”) collision operator is defined by

$$C_{ab}[f_a, f_b] = C_{ab}(\mathbf{x}, \mathbf{v}) = \frac{\partial}{\partial \mathbf{v}} \cdot \left[\mathbf{A}_{ab} f_a + \frac{\partial}{\partial \mathbf{v}} \cdot (\mathbf{D}_{ab} f_b) \right]$$

- Gyroaveraged** of C_{ab} performed in $(\mathbf{R}, \mu, v_{\parallel}, t)$ phase-space and project on Hermite-Laguerre, i.e.

$$C_{ab}^{pj} = \int d\mathbf{v} H_p(v_{\parallel}) L_j(\mu) \langle C_{ab} \rangle$$

Full-F Nonlinear GK Coulomb Collision Operator

$$C_{ab}^{pj} = \nu_{ab} \sum_{\mathbf{k}, \mathbf{k}'} \sum_{l=0}^{\infty} \sum_{m=-l}^{m=l} \sum_{r,s,p,q,\dots} \left[\nu_{ab\dots}^{lpj\dots} (T^{-1})_{pq}^{rsm} \right] \cdot \left[K_s(b_a) K_q(b'_b) \right] \left[N_a^{rs}(\mathbf{k}) N_b^{pq}(\mathbf{k}') \right],$$

Numerical coefficients

FLR Kernels

Gyro-Moments

- Arbitrary k_{\perp} , and m_a/m_b , T_a/T_b ratios

Jorge R., B. J. Frei and P. Ricci, JPP **85**, 905850604 (2019)

- Assuming $F_a = F_{aM} + f_a$ ($f_a \ll F_{aM}$)

$$C_{ab} \simeq C_{ab}^T + C_{ab}^F \quad C_{ab}^T \stackrel{\text{Test}}{=} C_{ab}[f_a, f_{bM}] \quad C_{ab}^F \stackrel{\text{Field}}{=} C_{ab}[f_{aM}, f_b]$$

- Linearized **GK Coulomb Collision Operator** in the gyro-moment approach

$$C_{ab}^{Tpj} = \sum_{l=0}^{\infty} \sum_{m=-l}^{m=l} \sum_{r,s,p,q,\dots} \underbrace{\nu_{ab\dots}^{Tlpj} \dots (T^{-1})_{pq}^{rsm}}_{\text{Numerical coefficients}} \underbrace{K_s(b_a)}_{\text{FLR Kernels}} \underbrace{N_a^{rs}}_{\text{Gyro-Moments}}$$

- Gyro-moment approach applied to different **simplified/ad hoc** linearized collision operators (e.g. GK/DK Sugama, GK/DK pitch-angle, GK/DK Dougherty)

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- Assuming $F_a = F_{aM} + f_a$ ($f_a \ll F_{aM}$)

$$\begin{aligned}
 \frac{\partial}{\partial t} \mathbf{N}_a &+ \overset{\text{Landau Damping}}{\mathbf{H}_a \cdot \nabla_{\parallel}} \mathbf{N}_a + \overset{\text{Particle trapping}}{(\mathbf{M}_a \cdot \mathbf{N}_a) \nabla_{\parallel}} \ln B + \overset{\text{Magnetic Drifts}}{\mathbf{D}_a \cdot \mathbf{N}_a} \\
 &= \underset{\text{Background Gradient drive}}{\mathbf{S}_\phi \phi} + \underset{\text{Linearized GK Collision Operators}}{\mathbf{S}_\psi \psi} + \sum_b \left(\mathbf{C}_{ab}^T \cdot \mathbf{N}_a + \mathbf{C}_{ab}^F \cdot \mathbf{N}_b \right)
 \end{aligned}$$

With $\delta \mathbf{E}_{es} = -\nabla \phi$, $\delta \mathbf{B}_{\perp} \simeq -\mathbf{b} \times \nabla \psi$ and $\mathbf{N}_a = [N_a^{00}, N_a^{01}, \dots, N_a^{10}, N_a^{11}, \dots, N_a^{PJ}]^T$

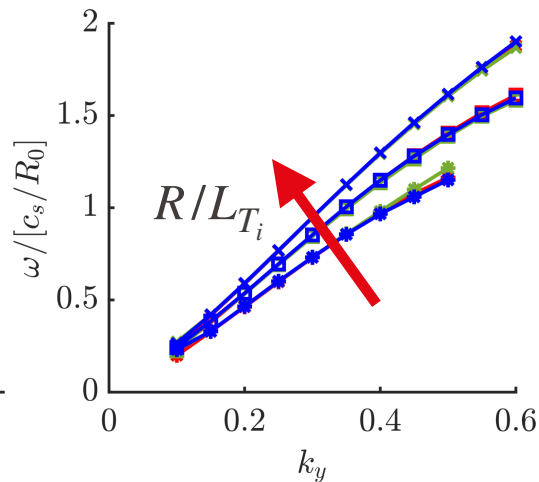
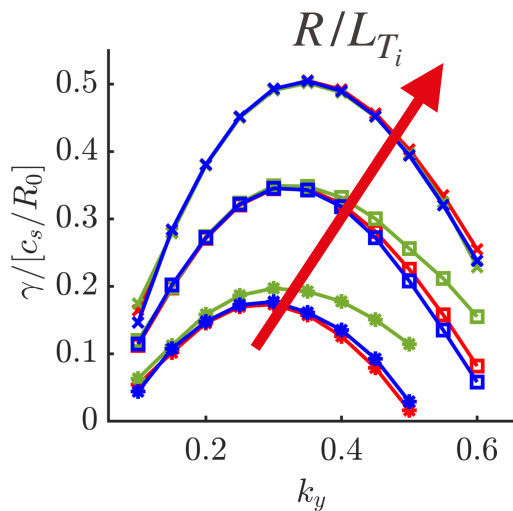
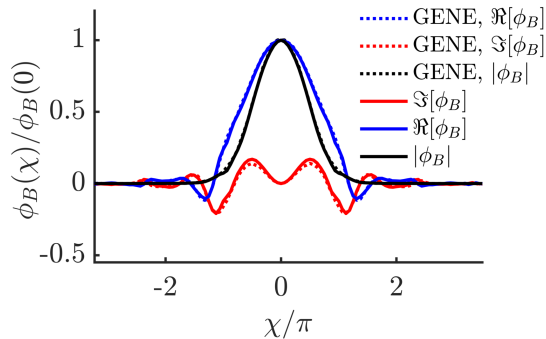
- Closure by truncation (arbitrary collisionality closure available): $P = P(\nu, k_{\parallel})$, $J = J(\nu, k_{\perp})$
- Implemented in a flux-tube code with concentric, circular flux surface: $\mathbf{B} = B_0 \nabla x \times \nabla y$, $\mathbf{k}_{\perp} = k_x \nabla x + k_y \nabla y$

Benchmark ($\nu \ll 1$) with GENE: Cyclone Base Case

Collisionless ITG with adiabatic electron:

- $\mathbf{B} = B_0 \nabla x \times \nabla y$ and $\mathbf{k}_\perp = k_x \nabla x + k_y \nabla y$

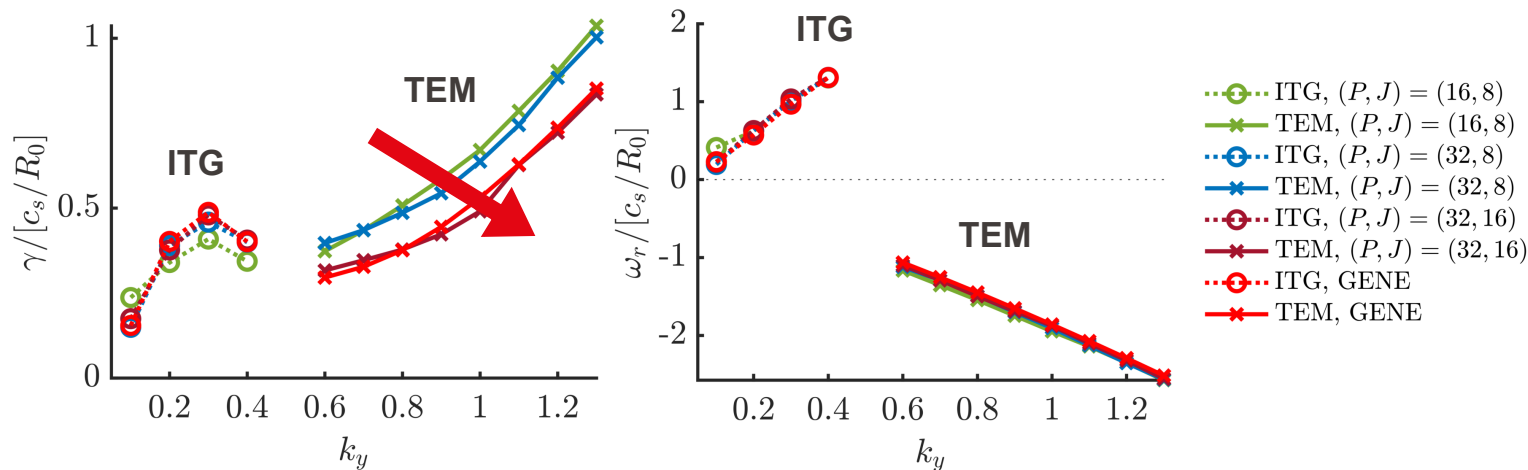
$$R_{T_i} = \frac{R}{L_{T_i}} \simeq -R \partial_x \ln T_i$$



- \square $R_{T_i} = 8$, GENE
- \times $R_{T_i} = 10$, GENE
- \star $R_{T_i} = 6$, GENE
- \star $(P, J) = (16, 8)$
- \square $(P, J) = (16, 8)$
- \times $(P, J) = (16, 8)$
- \star $(P, J) = (32, 16)$
- \square $(P, J) = (32, 16)$
- \times $(P, J) = (32, 16)$

Benchmark ($\nu \ll 1$) with GENE: ITG to TEM mode

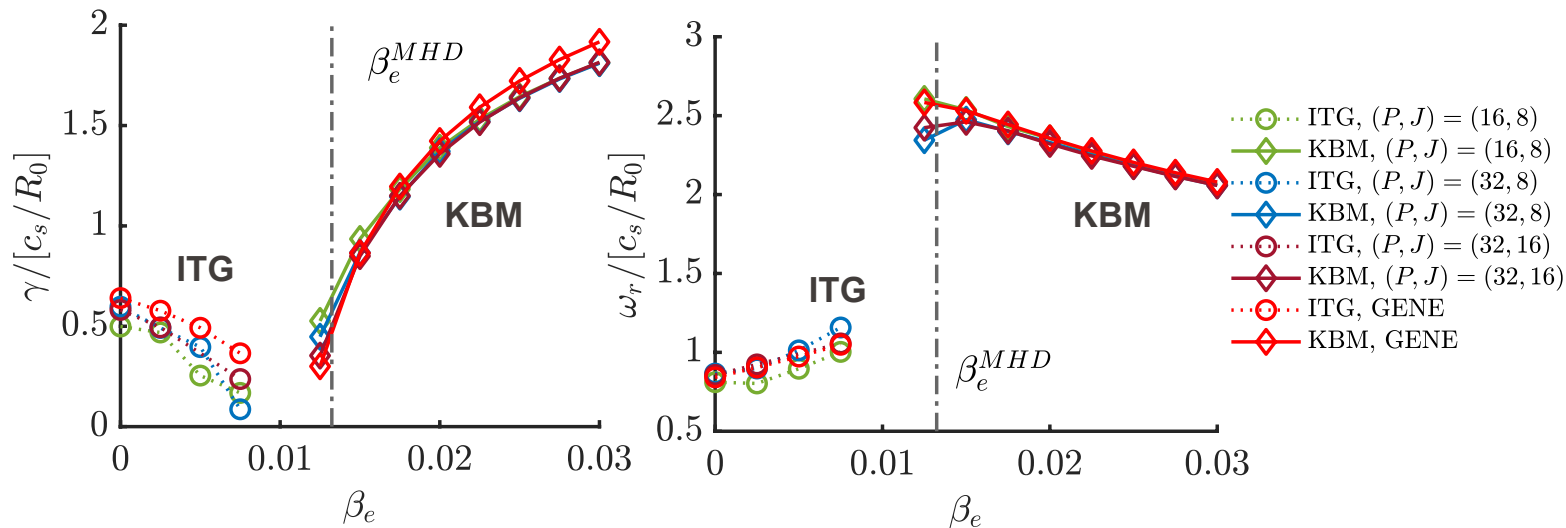
- Retrieve **collisionless ITG to TEM** transition (cyclone base case parameters) with $(P, J) \simeq (32, 16)$ (GENE velocity-space resolution $(N_{v\parallel}, N_{\mu}) \simeq (64, 24)$)



- Good agreement with ballooning mode structures
- Strong kinetic physics (e.g. trapped/passing boundary) **require larger** (P, J)

Benchmark ($\nu \ll 1$) with GENE: ITG to KBM mode

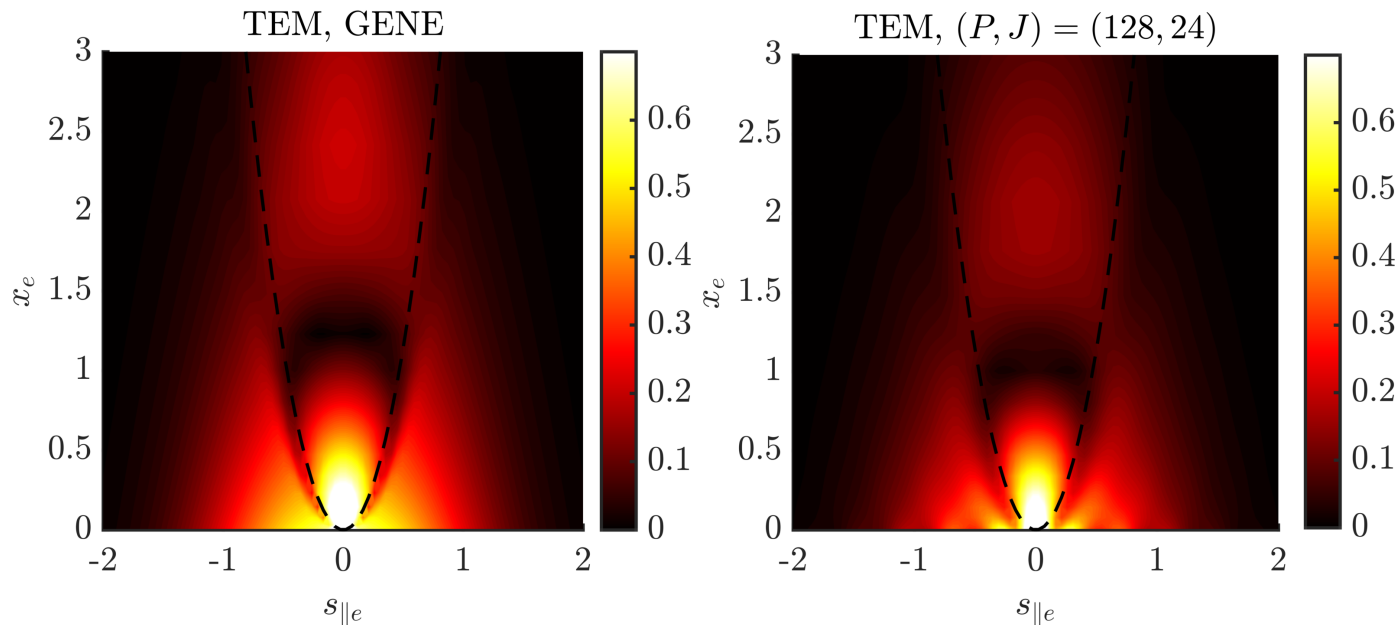
- Retrieve **KBM** transition (cyclone base case parameters) at finite β_e , with $(P, J) \simeq (16, 8)$



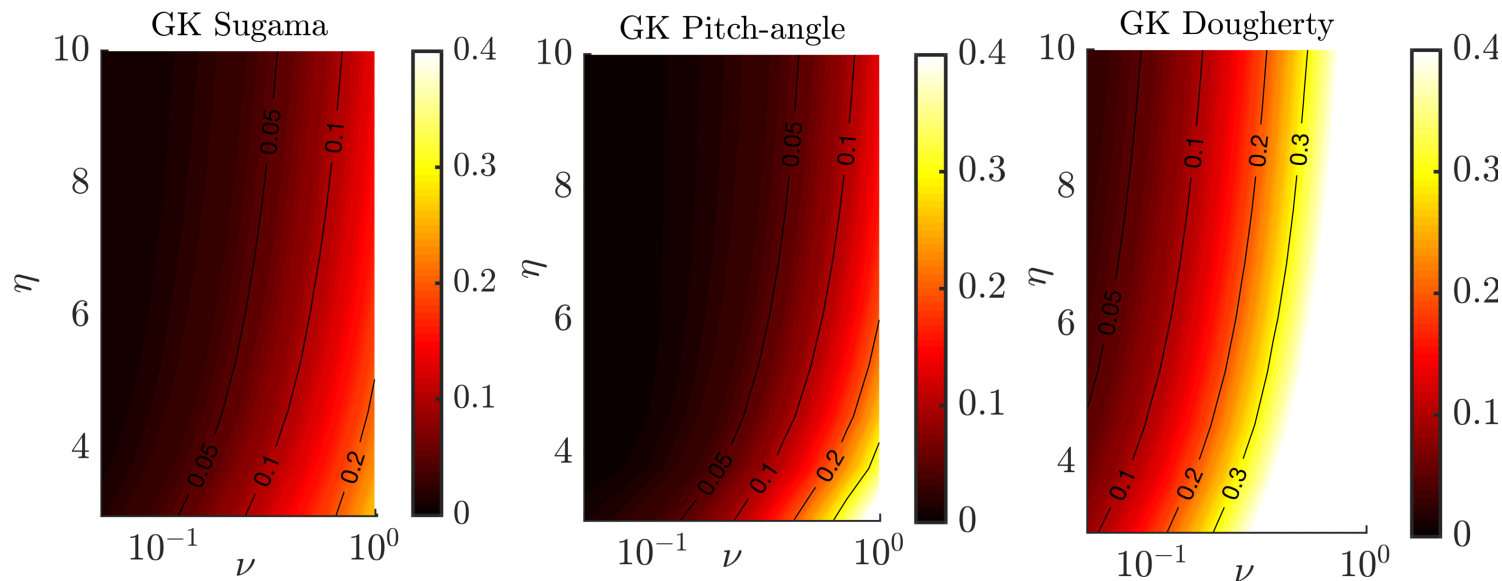
- Low number of gyro-moments (w.r.t to ITG) underlies the “fluid-like” properties of KBM

Benchmark ($\nu \ll 1$) with GENE: ITG to TEM modes

- Fine scale structures near **passing/trapped boundary**
- With sufficient number of gyro-moments, good qualitative agreement with the modulus of f_e at the outboard mid plane



- Relative deviation of ITG growth rate peak w.r.t GK Coulomb as a function of collisionality, $\nu = \nu_{ii}/[c_s/L_N]$, and normalized temperature gradient, $\eta = L_N/L_T$



- Deviation up to 20% with GK Sugama (benchmarked with GENE)

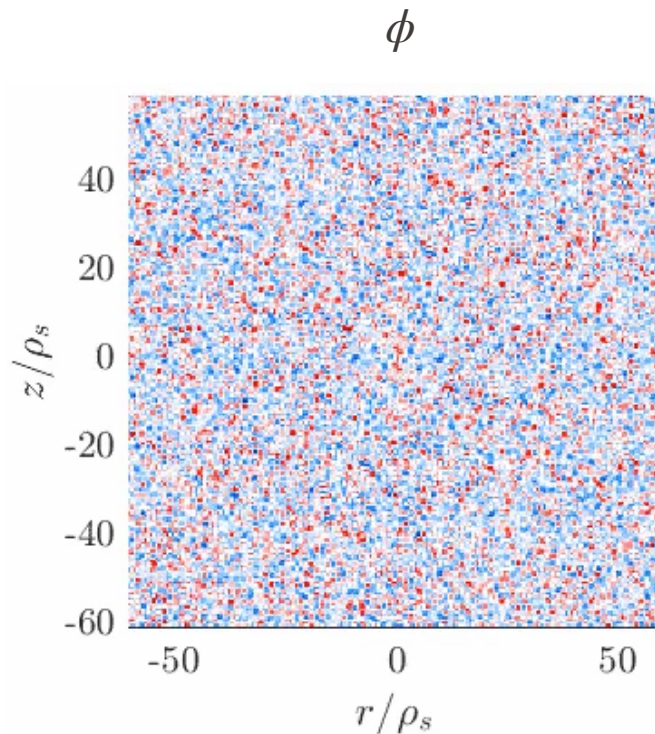
- Z-Pinch geometry ($k_{\parallel} = 0$), kinetic electrons

- Retain $\mathbf{E} \times \mathbf{B}$ nonlinearity

$$\sim \sum_{\mathbf{k}, \mathbf{k}'} \frac{i\mathbf{E}(\mathbf{k}) \times \mathbf{B}}{B^2} \cdot \mathbf{k}' N_a^{pj}(\mathbf{k}') e^{i(\mathbf{k} + \mathbf{k}') \cdot \mathbf{R}}$$

- Collisional effects with **linearized GK collision operators** (GK Coulomb, GK Sugama, GK Dougherty)

- Cyclic, bursting behaviour of fluctuations, transport, and zonal flows



- **A new GK model** for the tokamak boundary based on a **Hermite-Laguerre polynomials basis** (gyro-moment expansion)
- **A nonlinear GK Coulomb collision operator** (arbitrary mass and temperature ratios) is derived; a **linearized version derived and numerically implemented**; gyro-moment expansion of different **simplified/ad hoc** collision operator models (e.g., Sugama operator)
- Recover relevant **microinstability linear properties** (e.g., TEM, ITG and KBM); in agreement with GENE
- Explore **deviations** between **GK Coulomb** and other GK collision operator models (e.g. ITG mode)
- **Number of gyro-moment decreases with collisionality (SOL)** \Rightarrow optimal for edge and SOL modelling.
- Extend to flux-tube code to nonlinear and realistic 3D equilibrium and Full-F conditions

Thank you for your attention