

The role of neutral gas in validated global edge turbulence simulations

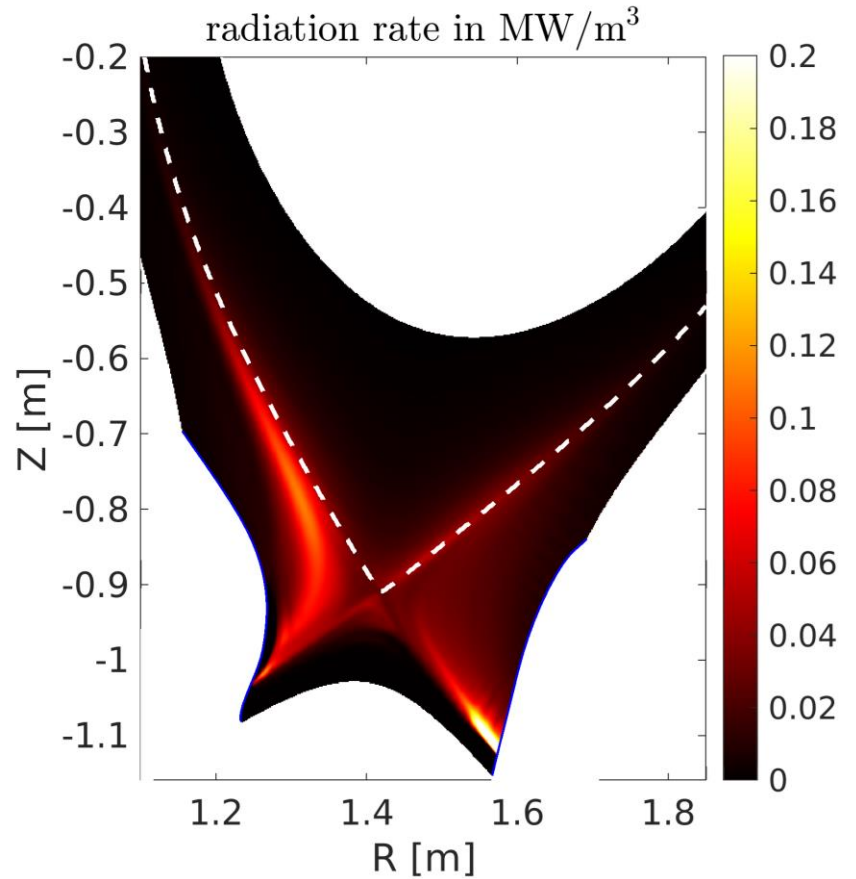
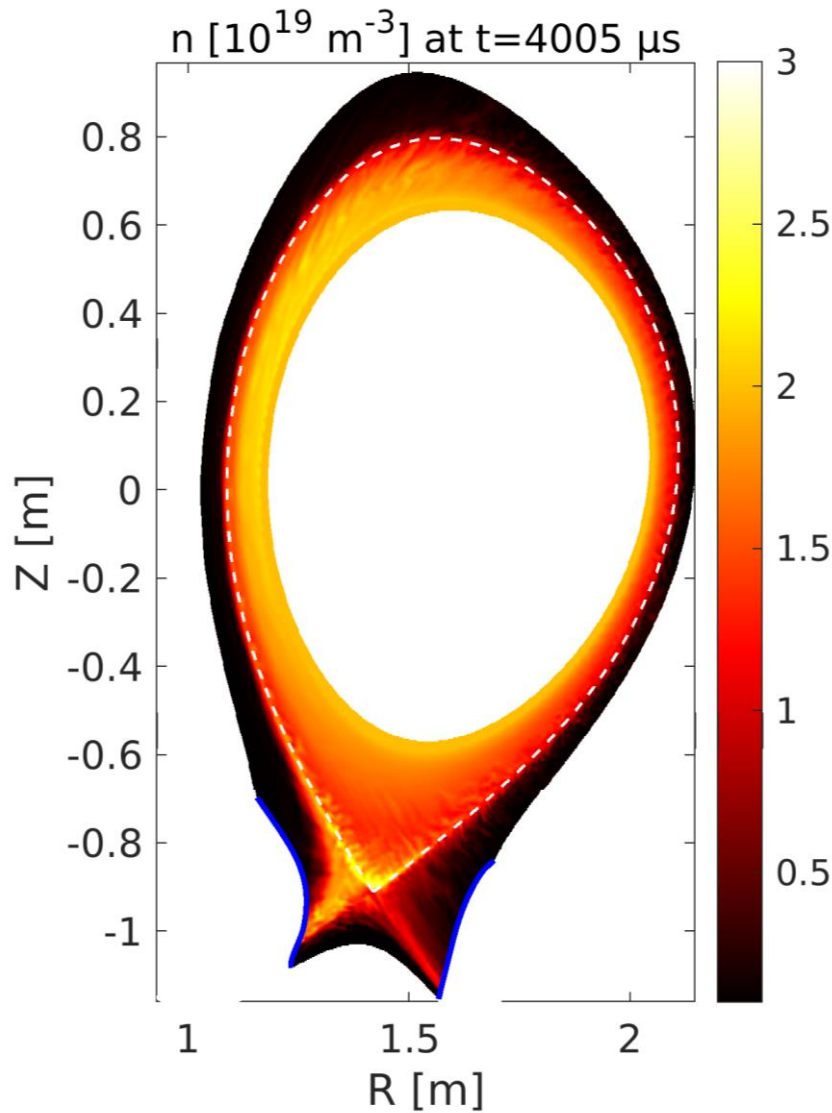
W. Zholobenko, A. Stegmeir, M. Griener, G. D. Conway, T. Body, D. Coster, F. Jenko and the ASDEX Upgrade Team

W. Zholobenko *et al* 2021 *Nucl. Fusion* **61** 116015,
<https://doi.org/10.1088/1741-4326/ac1e61>

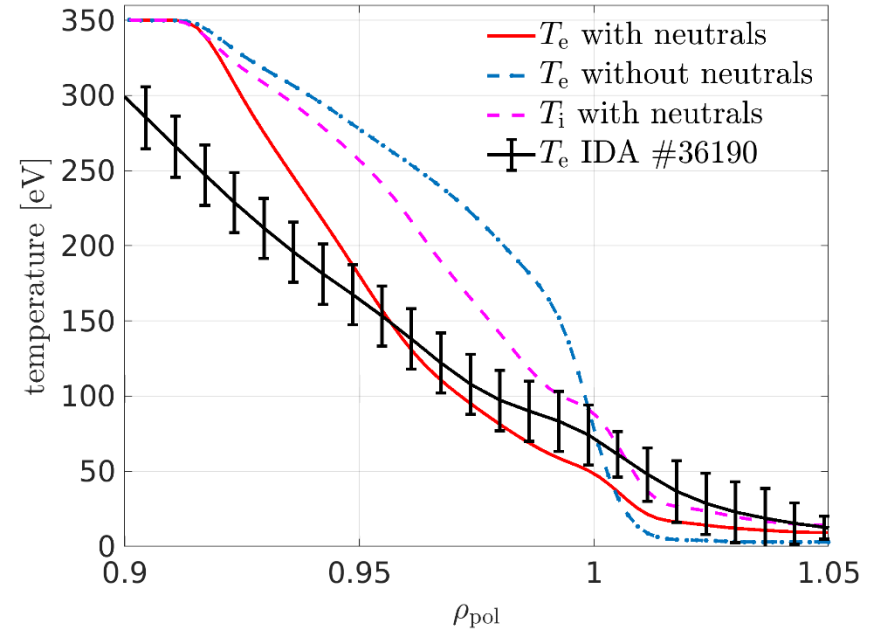
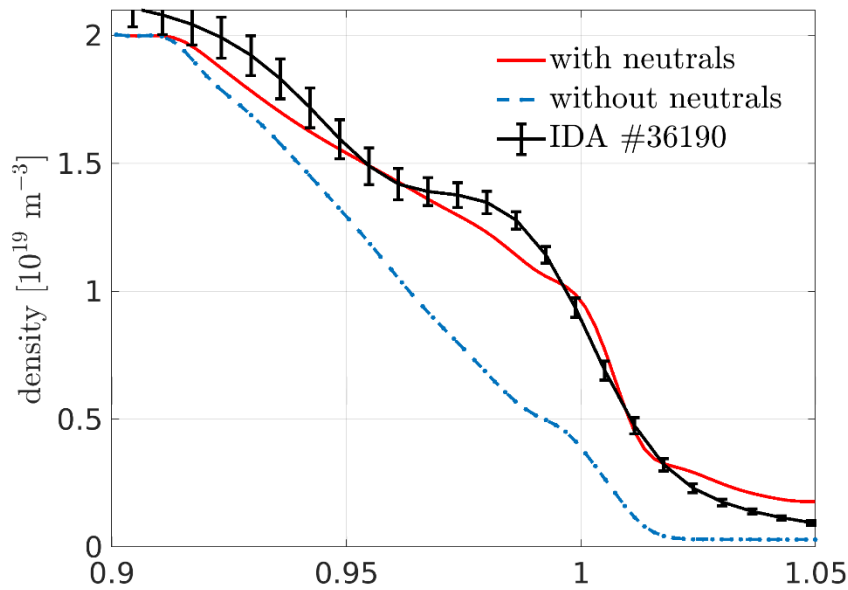
Global drift-reduced Braginskii with diffusive neutrals

plasma density	$\frac{d}{dt}n = nC(\varphi) - C(p_e) + \nabla \cdot [(j_{\parallel} - nu_{\parallel}) \mathbf{b}] + \mathcal{D}_n(n) + k_{iz}nN + S_n,$
quasineutrality / vorticity	$\nabla \cdot \left[\frac{n}{B^2} \left(\frac{\partial}{\partial t} + \delta_0 \left(\frac{\mathbf{B}}{B^2} \times \nabla_{\perp} \varphi \right) \cdot \nabla_{\perp} + Nk_{iz} + Nk_{cx} + u_{\parallel} \nabla_{\parallel} \right) \left(\nabla_{\perp} \varphi + \zeta \frac{\nabla_{\perp} p_i}{n} \right) \right] =$ $-C(p_e + \zeta p_i) + \nabla \cdot (j_{\parallel} \mathbf{b}) - \frac{\zeta}{6} C(G) + \mathcal{D}_{\Omega}(\Omega),$
momentum	$\left(\frac{d}{dt} + u_{\parallel} \nabla_{\parallel} \right) u_{\parallel} = -\frac{\nabla_{\parallel} (p_e + \zeta p_i)}{n} + \zeta T_i C(u_{\parallel}) - \frac{2}{3} \zeta \frac{B^{3/2}}{n} \nabla_{\parallel} \frac{G}{B^{3/2}} + \mathcal{D}_u(u_{\parallel}),$
Ohm & Ampère	$\beta_0 \frac{\partial}{\partial t} A_{\parallel} + \mu \left(\frac{d}{dt} + v_{\parallel} \nabla_{\parallel} \right) \frac{j_{\parallel}}{n} = - \left(\frac{\eta_{\parallel 0}}{T_e^{3/2}} \right) j_{\parallel} - \nabla_{\parallel} \varphi + \frac{\nabla_{\parallel} p_e}{n} + 0.71 \nabla_{\parallel} T_e + \mathcal{D}_{\Psi}(\Psi_m), \quad \nabla_{\perp}^2 A_{\parallel} = -j_{\parallel},$
electron heat	$\frac{3}{2} \left(\frac{d}{dt} + v_{\parallel} \nabla_{\parallel} \right) T_e = T_e C(\varphi) - \frac{T_e}{n} C(p_e) - \frac{5}{2} T_e C(T_e) - T_e \nabla \cdot (v_{\parallel} \mathbf{b}) + 0.71 \frac{T_e}{n} \nabla \cdot (j_{\parallel} \mathbf{b})$ $+ \frac{1}{n} \nabla \cdot \left[(\chi_{\parallel e 0} T_e^{5/2}) \mathbf{b} \nabla_{\parallel} T_e \right] - 2\nu_{e0} \mu \left(\frac{n}{T_e^{3/2}} \right) (T_e - \zeta T_i) + \left(\frac{\eta_{\parallel 0}}{T_e^{3/2}} \right) \frac{j_{\parallel}^2}{n} + \frac{3}{2} \left(\mathcal{D}_{T_e}(T_e) + S_{T_e} \right),$ $S_{T_e} = -\frac{2}{3} (W_{iz} N + W_{rec} n) - (k_{iz} N - k_{rec} n) T_e$
ion heat	$\frac{3}{2} \left(\frac{d}{dt} + u_{\parallel} \nabla_{\parallel} \right) T_i = T_i C(\varphi) - \frac{T_i}{n} C(p_e) + \frac{5}{2} \zeta T_i C(T_i) - T_i \nabla \cdot (u_{\parallel} \mathbf{b}) + \frac{T_i}{n} \nabla \cdot (j_{\parallel} \mathbf{b})$ $+ \frac{1}{n} \nabla \cdot \left[(\chi_{\parallel i 0} T_i^{5/2}) \mathbf{b} \nabla_{\parallel} T_i \right] + 2\nu_{e0} \mu \left(\frac{n}{T_e^{3/2}} \right) \left(\frac{1}{\zeta} T_e - T_i \right) + \frac{2w_{GTi}}{9\eta_{i0}} \frac{G^2}{n T_i^{5/2}} + \frac{3}{2} \left(\mathcal{D}_{T_i}(T_i) + S_{T_i} \right),$
neutral particles	$\frac{\partial}{\partial t} N = \nabla \cdot \frac{\zeta \nabla N T_i}{n k_{cx}} - k_{iz} n N, \quad N \text{ fixed at the divertor.}$

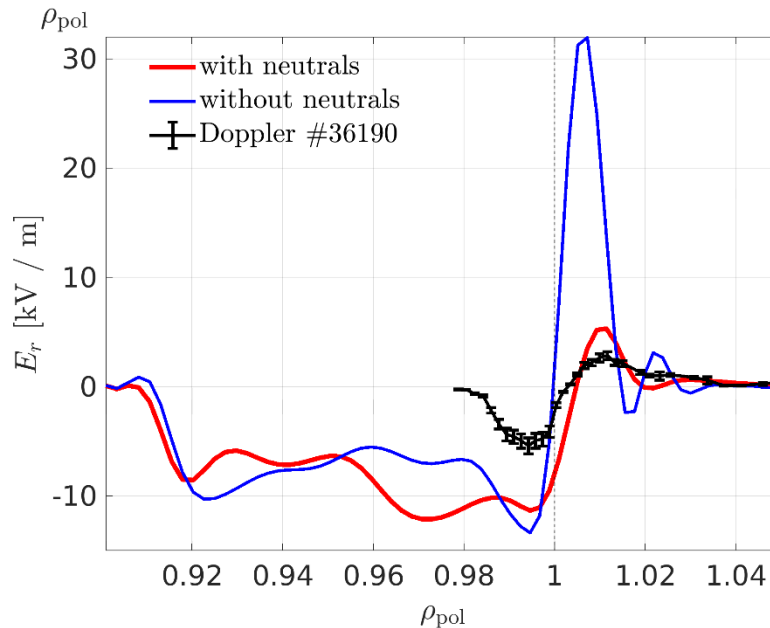
Ionization source and poloidal asymmetry



Validation against attached AUG L-mode (#36190)



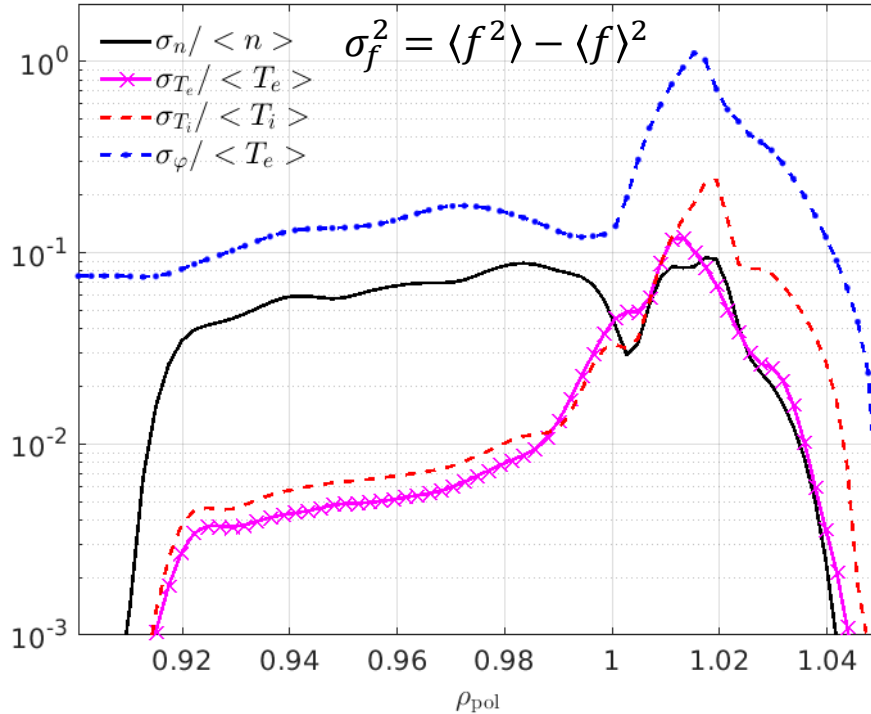
outboard mid-plane profiles



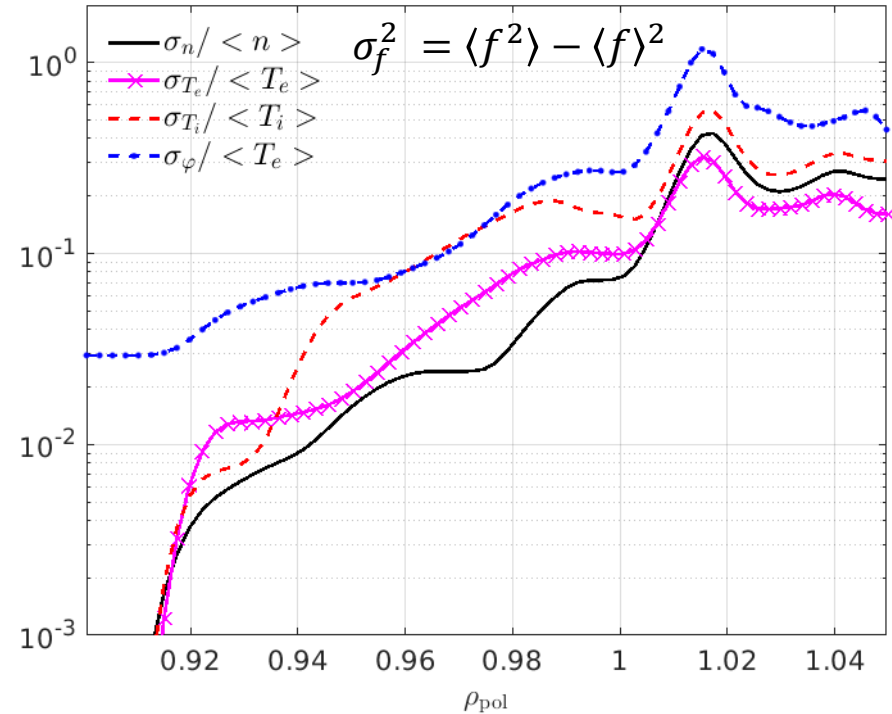
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Wladimir Zholobenko, TSVV 3 - neutrals

without neutrals (core source)



with neutrals (edge source)



$$\frac{\sigma_\varphi}{\langle T_e \rangle} > 2 \frac{\sigma_n}{\langle n \rangle} \gg \frac{\sigma_T}{\langle T \rangle} \Rightarrow \text{ballooning}$$

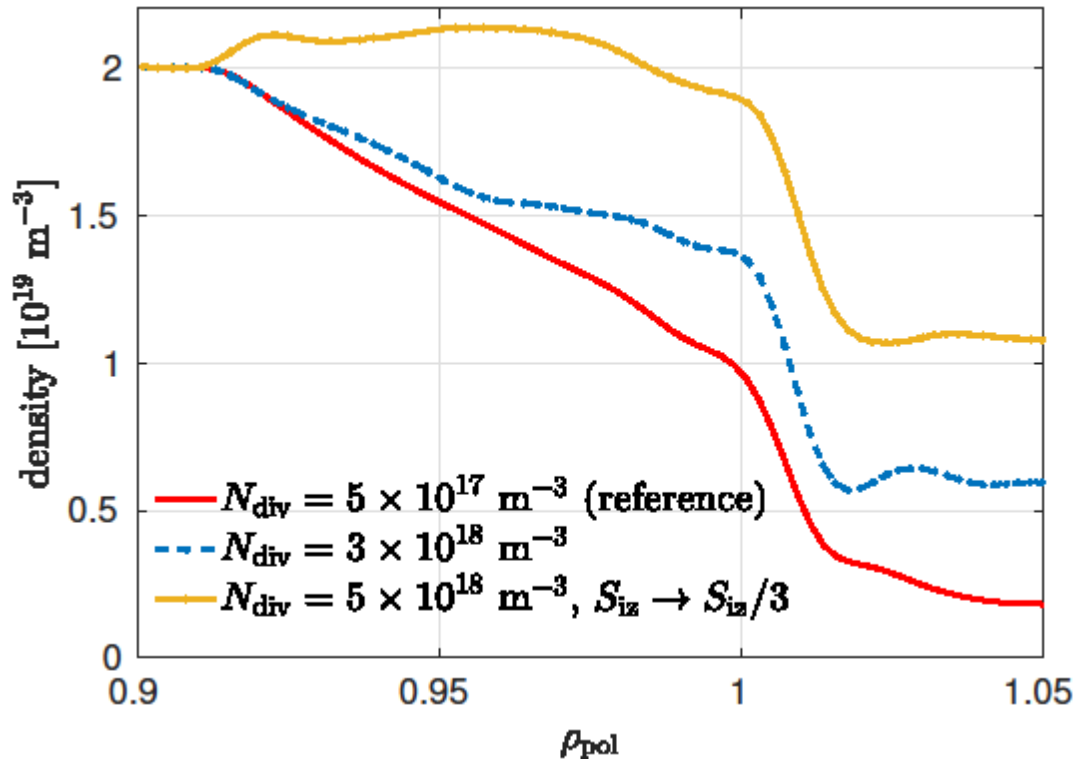
$$\chi_1^e = 2/3D_\perp$$

$$\frac{\sigma_\varphi}{\langle T_e \rangle} \geq \frac{\sigma_{T_i}}{\langle T_i \rangle} > \frac{\sigma_{T_e}}{\langle T_e \rangle} > \frac{\sigma_n}{\langle n \rangle} \Rightarrow \text{ITG}$$

$$\langle D_\perp \rangle \approx 0$$

Towards high-recycling conditions

outboard mid-plane plasma density as a function of divertor neutrals density



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- Even the simple diffusive neutral gas model makes simulations much more realistic!
- Will take advantage of in upcoming modelling efforts
- Extensions naturally remain of high interest!