

BOUNDARY CONDITIONS IN THE GBS SIMULATION OF THE COMPASS TOKAMAK

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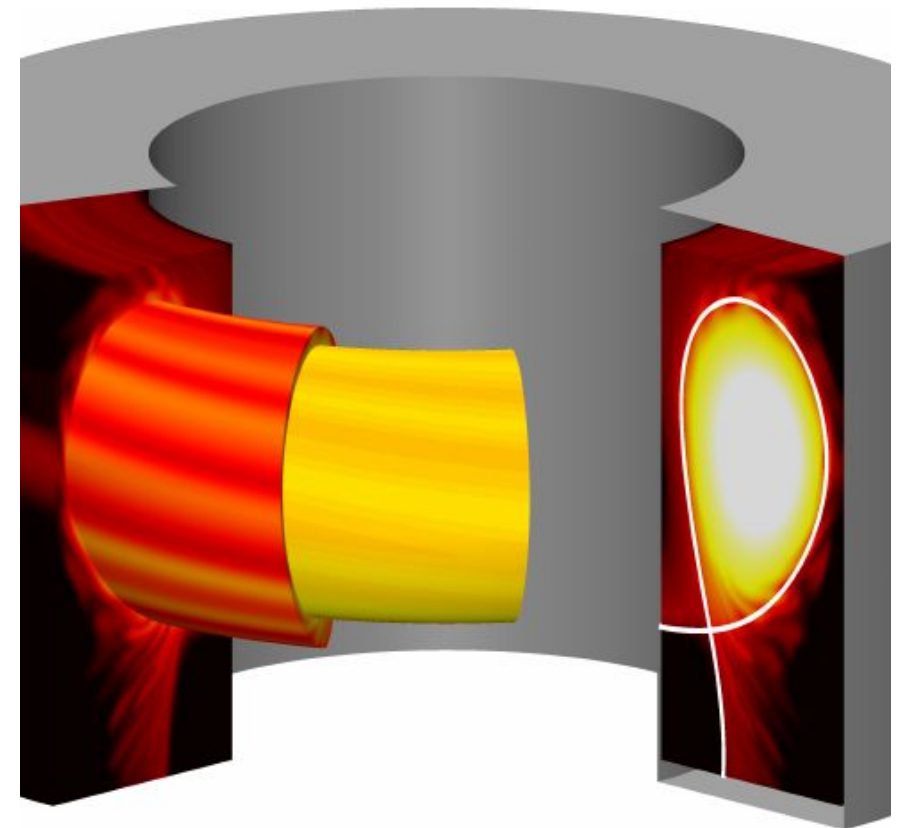
List of affiliations:

- 1) Institute of the plasma physics of the CAS
- 2) Faculty of nuclear science and physical engineering
- 3) Ecole Polytechnique Fédérale de Lausanne (EPFL), Swiss Plasma Center (SPC)

- Boundary condition types
- BCs used in COMPASS simulation and simulation domain
- An impact of BCs on turbulence
- Impact of BCs on electron velocities and vertical electric field
- Conclusion

GLOBAL BRAGINSKII SOLVER

- First principle, 3D, flux-driven, global, turbulence code for plasma edge simulations based on Braginskii equations [1].
- Full plasma volume, Divertor geometry, electromagnetic effects, kinetic neutrals, ion temperature dynamics, self-consistent turbulence evolution.
- High computational requirements (~2000 cores, ~5-10 M CPU hours).
- Validation on COMPASS tokamak – first validation of full-size simulation after TCV.
- Validation on COMPASS will include electron temperature and plasma potential fluctuations.



EQUATIONS

- Braginskii equations are solved, Boussinesq approximation is not used.
- 7 fields are evolved during each step:
 - Density, electron and ion parallel velocity, vorticity, electron and ion temperature, and psi (if electromagnetic effects are enabled).
- If kinetic neutrals are included:
 - Neutral density, and neutral parallel velocity.

$$\nabla \cdot (n \nabla_{\perp} \phi) = \Omega - \frac{\nabla_{\perp}^2 p_i}{e},$$

$$\left(\nabla_{\perp}^2 - \frac{e^2 \mu_0}{m_e} n \right) v_{\parallel e} = \nabla_{\perp}^2 U_{\parallel e} - \frac{e^2 \mu_0}{m_e} n v_{\parallel i} + \frac{e^2 \mu_0}{m_e} \bar{j}_{\parallel},$$

$$U_{\parallel e} = v_{\parallel e} + e\psi/m_e$$

Particle confinement $\frac{\partial n}{\partial t} = -\frac{1}{B}[\phi, n] + \frac{2}{eB} [C(p_e) - nC(\phi)] - \nabla_{\parallel}(nv_{\parallel e}) + D_n \nabla_{\perp}^2 n + s_n + v_{iz}n_n - v_{rec}n,$ (1)

Vorticity $\frac{\partial \Omega}{\partial t} = -\frac{1}{B} \nabla \cdot [\phi, \omega] - \nabla \cdot (v_{\parallel i} \nabla_{\parallel} \omega) + \frac{B\Omega_{ci}}{e} \nabla_{\parallel} j_{\parallel} + \frac{2\Omega_{ci}}{e} C(p_e + p_i) + \frac{\Omega_{ci}}{3e} C(G_i) + D_{\Omega} \nabla_{\perp}^2 \Omega - \frac{n_n}{n} v_{cx} \Omega,$ (2)

Electron inertia $\frac{\partial U_{\parallel e}}{\partial t} = -\frac{1}{B}[\phi, v_{\parallel e}] - v_{\parallel e} \nabla_{\parallel} v_{\parallel e} + \frac{e}{m_e} \left(\frac{j_{\parallel}}{\sigma_{\parallel}} + \nabla_{\parallel} \phi - \frac{1}{en} \nabla_{\parallel} p_e - \frac{0.71}{e} \nabla_{\parallel} T_e - \frac{2}{3en} \nabla_{\parallel} G_e \right) + D_{v_{\parallel e}} \nabla_{\perp}^2 v_{\parallel e} + \frac{n_n}{n} (v_{en} + 2v_{iz})(v_{\parallel n} - v_{\parallel e}),$ (3)

ion inertia $\frac{\partial v_{\parallel i}}{\partial t} = -\frac{1}{B}[\phi, v_{\parallel i}] - v_{\parallel i} \nabla_{\parallel} v_{\parallel i} - \frac{1}{m_i n} \nabla_{\parallel} (p_e + p_i) - \frac{2}{3m_i n} \nabla_{\parallel} G_i + D_{v_{\parallel i}} \nabla_{\perp}^2 v_{\parallel i} + \frac{n_n}{n} (v_{iz} + v_{cx})(v_{\parallel n} - v_{\parallel i}),$ (4)

electron energy confinement $\frac{\partial T_e}{\partial t} = -\frac{1}{B}[\phi, T_e] - v_{\parallel e} \nabla_{\parallel} T_e + \frac{2}{3} T_e \left[0.71 \frac{\nabla_{\parallel} j_{\parallel}}{en} - \nabla_{\parallel} v_{\parallel e} \right] + \frac{4}{3} \frac{T_e}{eB} \left[\frac{7}{2} C(T_e) + \frac{T_e}{n} C(n) - eC(\phi) \right] + \nabla_{\parallel} (\chi_{\parallel e} \nabla_{\parallel} T_e) + D_{T_e} \nabla_{\perp}^2 T_e + s_{T_e} - \frac{n_n}{n} v_{en} m_e \frac{2}{3} v_{\parallel e} (v_{\parallel n} - v_{\parallel e}) - 2 \frac{m_e}{m_i} \frac{1}{\tau_e} (T_e - T_i) + \frac{n_n}{n} v_{iz} \left[-\frac{2}{3} E_{iz} - T_e + m_e v_{\parallel e} \left(v_{\parallel e} - \frac{4}{3} v_{\parallel n} \right) \right],$ (5)

ion energy confinement $\frac{\partial T_i}{\partial t} = -\frac{1}{B}[\phi, T_i] - v_{\parallel i} \nabla_{\parallel} T_i + \frac{4}{3} \frac{T_i}{eB} \left[C(T_e) + \frac{T_e}{n} C(n) - eC(\phi) \right] - \frac{10}{3} \frac{T_i}{eB} C(T_i) + \frac{2}{3} T_i \left[(v_{\parallel i} - v_{\parallel e}) \frac{\nabla_{\parallel} n}{n} - \nabla_{\parallel} v_{\parallel e} \right] + \nabla_{\parallel} (\chi_{\parallel i} \nabla_{\parallel} T_i) + D_{T_i} \nabla_{\perp}^2 T_i + s_{T_i} + 2 \frac{m_e}{m_i} \frac{1}{\tau_e} (T_e - T_i) + \frac{n_n}{n} (v_{iz} + v_{cx}) \left[T_n - T_i + \frac{1}{3} (v_{\parallel n} - v_{\parallel i})^2 \right],$ (6)

BOUNDARY CONDITIONS

- BCs play an important role in the simulation.
- Applied at the magnetic presheath.
- There are specific sets of BCs for plasma potential and for other fields.
- Plasma potential uses multiple conditions:
 - **Man** - Dirichlet, Neumann – fixing all fields.
 - **pAT** – fixing potential to ΛT_e , other fields Man.
 - **Tar** - fixing potential to ΛT_e , others Mag
 - **Robin** – allowing potential to vary from ΛT_e .
 - **Magnetic** – full conductive condition.

Man

$$\phi(x, z) = A$$

pAT

$$\phi(x, z) = \Lambda(x, z)T_e(x, z),$$

$$\Lambda(x, z) = \lambda - \sqrt{1 + \tau \frac{T_i(x, z)}{T_e(x, z)}}$$

Rob

$$\phi(x, z) = \Lambda(x, z)T_e(x, z) - R_{\text{fac}}(x)\partial_x \Psi(x)B_{\text{sign}}(x)T_e(x, z) (\partial_y n(x, z) + 1.71\partial_y T_e(x, z)),$$

where

$$R_{\text{fac}}(x) = \sum_z \frac{T_e(x, z)}{c_s(x)\sqrt{1 + \tau \frac{T_i(x, z)}{T_e(x, z)}}} \mu\mu_{\text{spitzer}}(x)F_{\text{fac}},$$

Mag

$$\partial_y \phi(x, z) = -\frac{\frac{\partial_x \Psi}{|\partial_x \Psi|} c_s(x, z)}{\sqrt{1 + \tau \frac{T_i(x, z)}{T_e(x, z)}}} \partial_y v_{\parallel i}(x, z),$$

$$v_{\parallel i} = \pm c_s \sqrt{1 + \frac{T_i}{T_e}},$$

$$v_{\parallel e} = \pm c_s \sqrt{1 + \frac{T_i}{T_e}} \exp\left(\Lambda - \frac{e\phi}{T_e}\right),$$

$$\partial_s n = \mp \frac{n}{c_s \sqrt{1 + \frac{T_i}{T_e}}} \partial_s v_{\parallel i},$$

$$\partial_s T_e = \partial_s T_i = 0,$$

$$\Omega = \mp \frac{m_i n}{e} c_s \sqrt{1 + \frac{T_i}{T_e}} \partial_{ss}^2 v_{\parallel i},$$

$$\partial_s \phi = \mp \frac{m_i c_s}{e \sqrt{1 + \frac{T_i}{T_e}}} \partial_s v_{\parallel i},$$

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$$\begin{aligned} v_{\parallel i} &= \pm c_s \sqrt{1 + \frac{T_i}{T_e}}, \\ v_{\parallel e} &= \pm c_s \sqrt{1 + \frac{T_i}{T_e}} \exp\left(\Lambda - \frac{e\phi}{T_e}\right), \\ \partial_s n &= \mp \frac{n}{c_s \sqrt{1 + \frac{T_i}{T_e}}} \partial_s v_{\parallel i}, \\ \partial_s T_e &= \partial_s T_i = 0, \\ \Omega &= \mp \frac{m_i n}{e} c_s \sqrt{1 + \frac{T_i}{T_e}} \partial_{ss}^2 v_{\parallel i}, \\ \partial_s \phi &= \mp \frac{m_i c_s}{e \sqrt{1 + \frac{T_i}{T_e}}} \partial_s v_{\parallel i}, \end{aligned}$$

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$$\phi(x, z) = \Lambda(x, z) T_e(x, z) - R_{\text{fac}}(x) \partial_x \Psi(x) B_{\text{sign}}(x) T_c(x, z) (\partial_y n(x, z) + 1.71 \partial_y T_c(x, z)),$$

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$$\partial_y \phi(x, z) = - \frac{\frac{\partial_x \Psi}{|\partial_x \Psi|} c_s(x, z)}{\sqrt{1 + \tau \frac{T_i(x, z)}{T_e(x, z)}}} \partial_y v_{\parallel i}(x, z),$$

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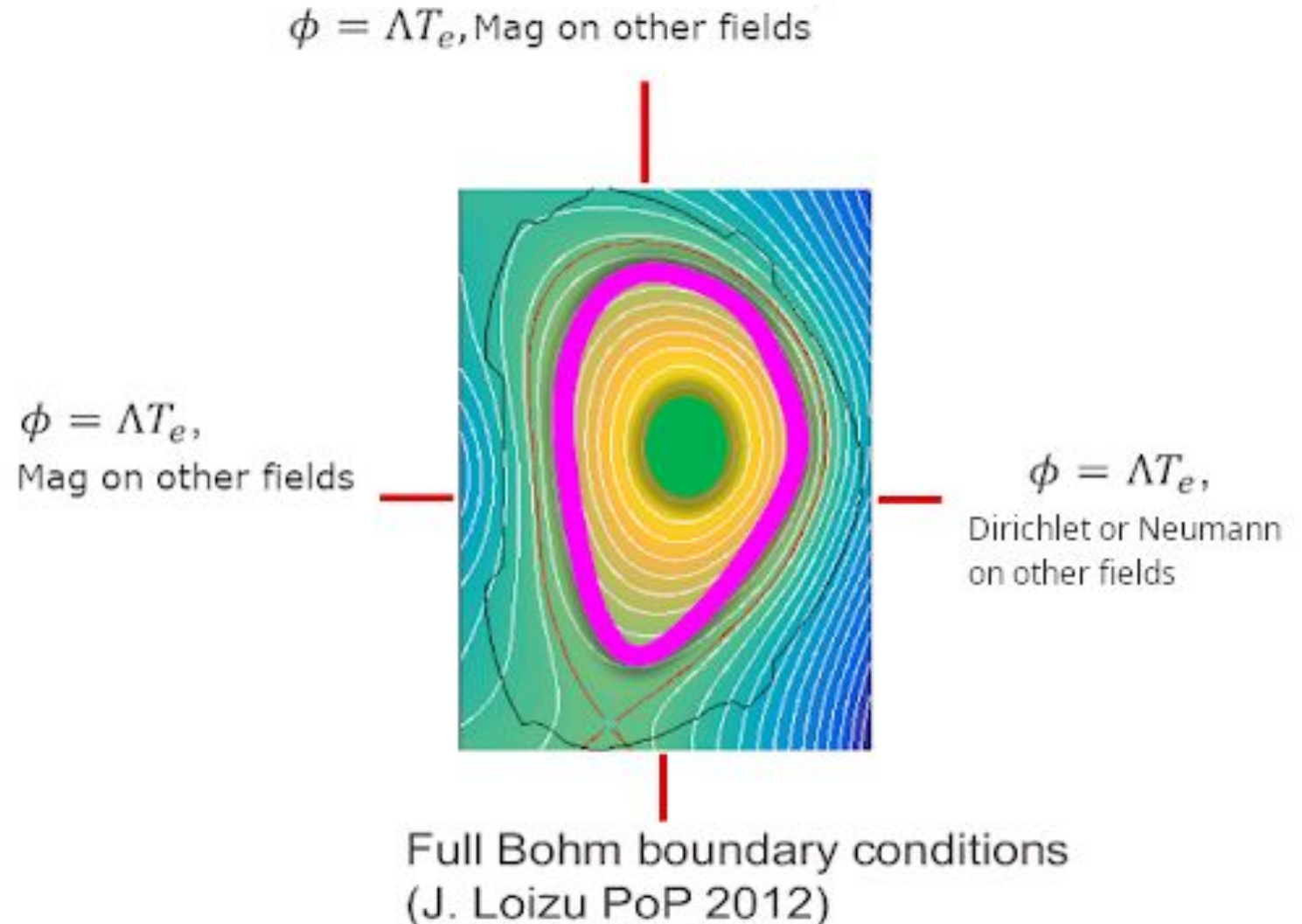
$$R_{\text{fac}}(x) = \sum_z \frac{T_e(x, z)}{c_s(x) \sqrt{1 + \tau \frac{T_i(x, z)}{T_e(x, z)}}} \mu \mu_{\text{spitzer}}(x) F_{\text{fac}},$$

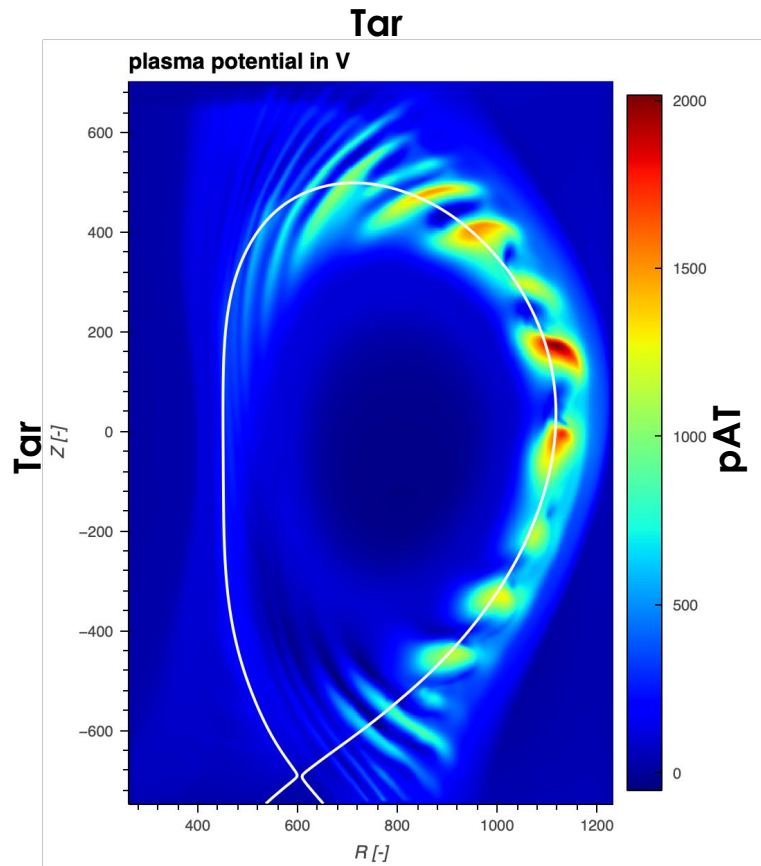
- **Tar** - fixing potential to ΛT_e , others Mag
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Mag

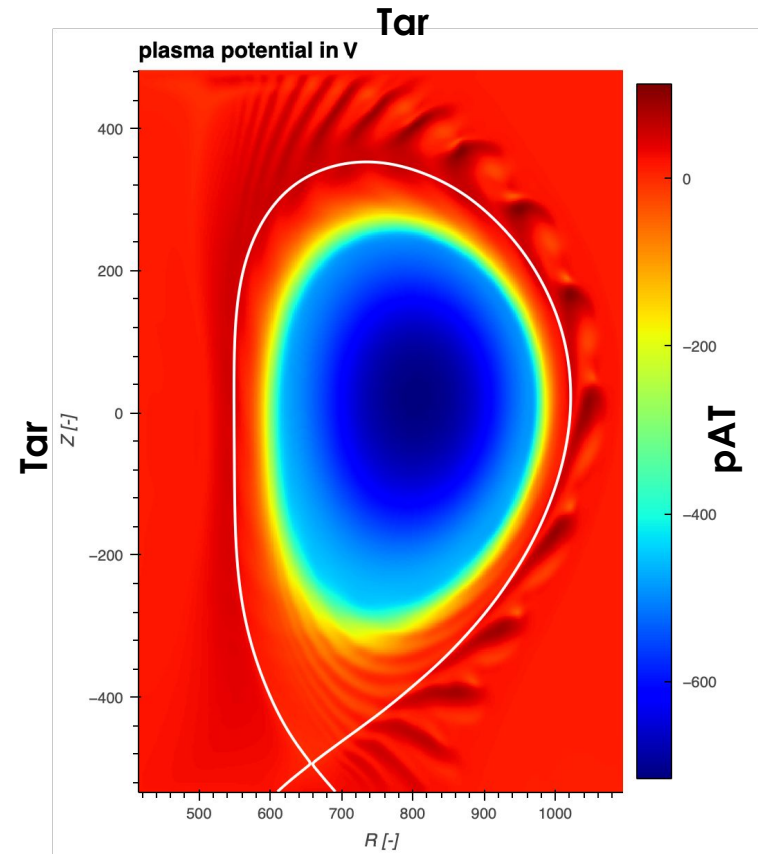
$$\partial_y \phi(x, z) = - \frac{\frac{\partial_x \Psi}{|\partial_x \Psi|} c_s(x, z)}{\sqrt{1 + \tau \frac{T_i(x, z)}{T_e(x, z)}}} \partial_y v_{\parallel i}(x, z),$$

- In the case of COMPASS simulation, all the **Tar**, **Rob**, and **Magnetic** BC were applied on the bottom boundary - the divertor position.
- **Tar** condition (potential fixed to λT_e) used at left and top boundary.
- **Dirichlet** and **Neumann** conditions are set on right boundary for all fields.
- It was shown, turbulent structures disappear before touching the right boundary.
- Influence of the Dirichlet boundary is therefore not propagating inside plasma.

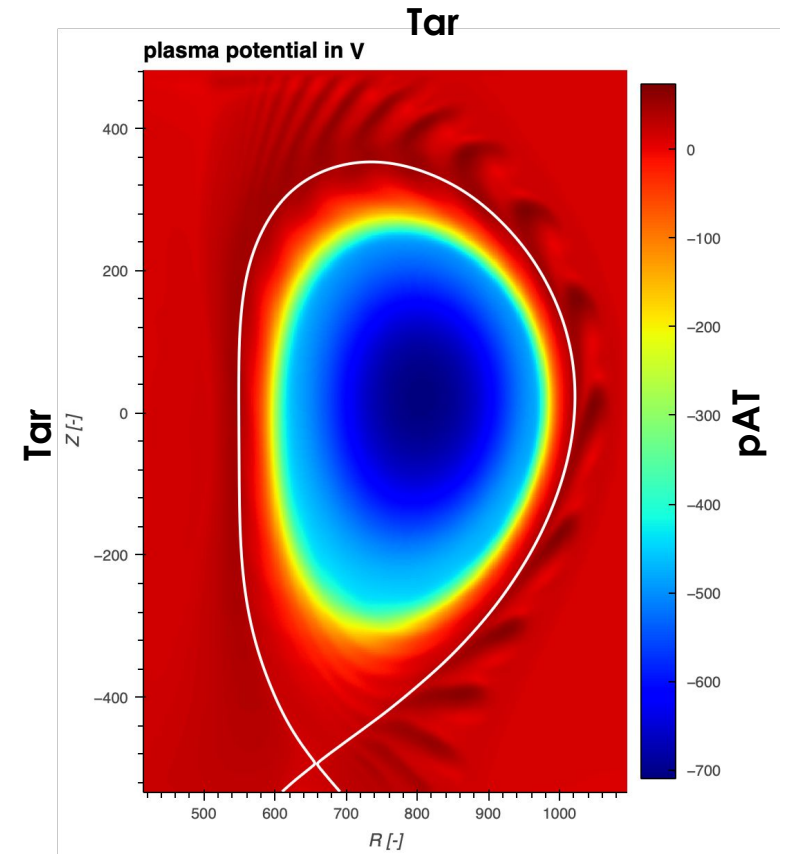




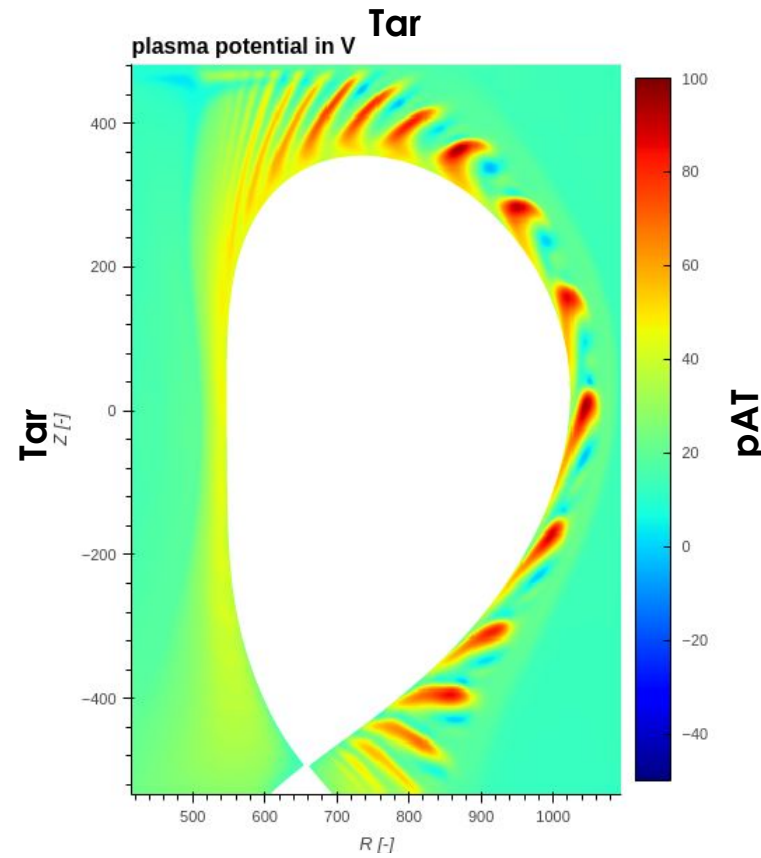
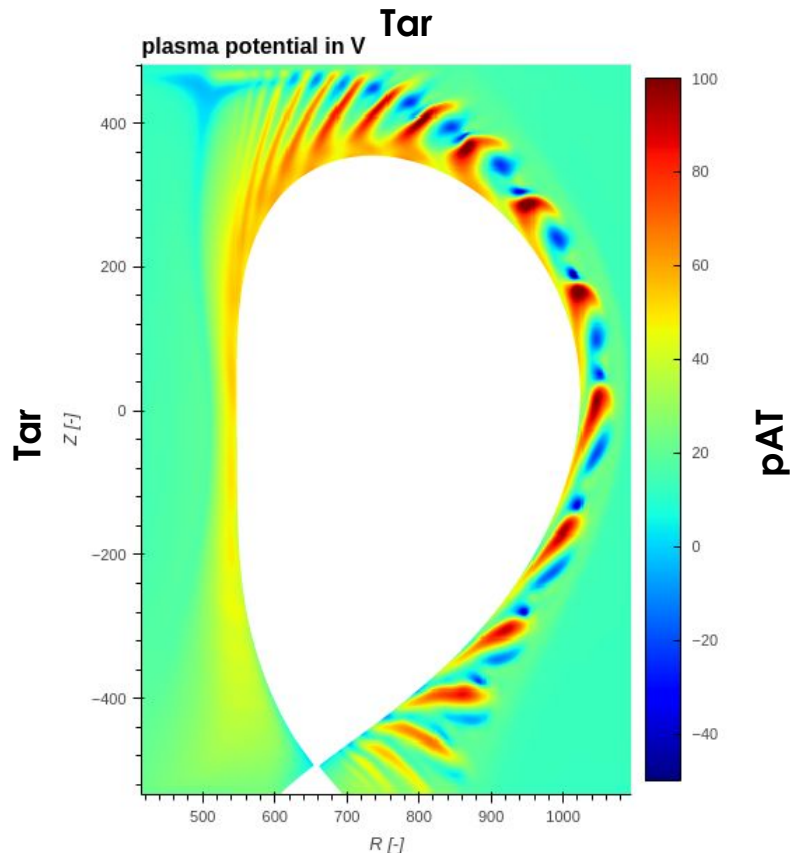
Tar
Potential fixed to ΔT_e ,
Other fields Mag



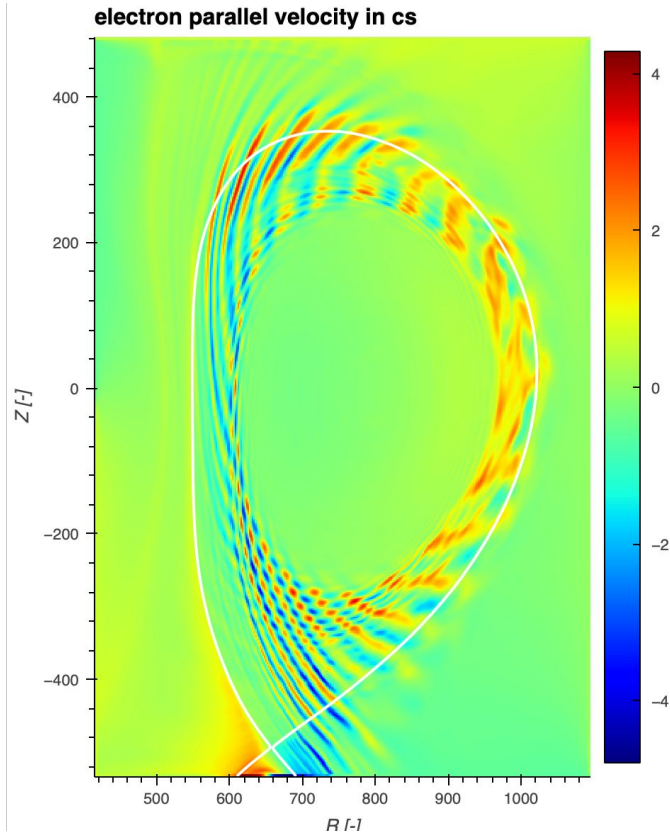
Rob
Potential can vary ΔT_e ,
Other fields Mag



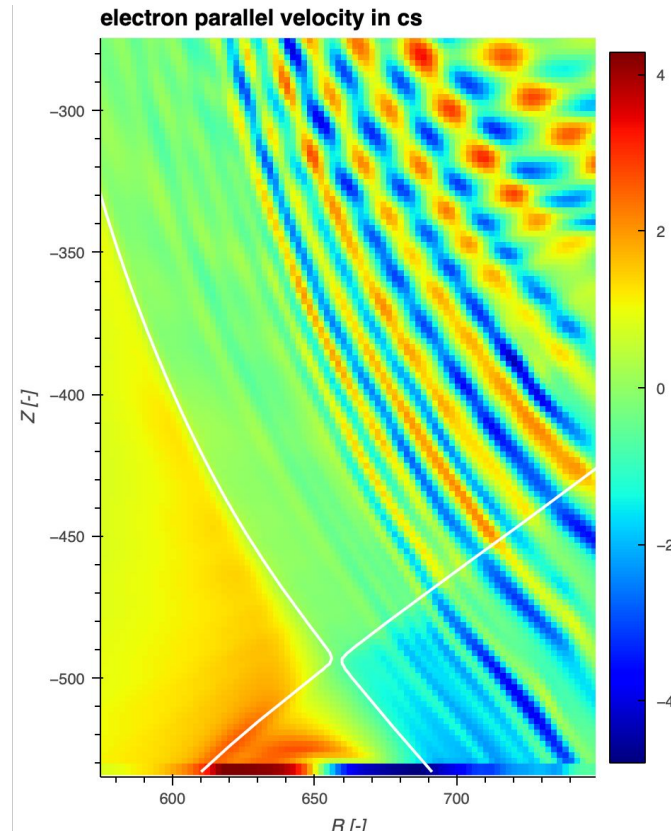
Mag
Full conductive



- Similar blob shapes and amplitudes are observed for both the **Rob** and **Mag**.
- A bit lower amplitudes of negative structures for the **Rob** BC.
- Both **Rob** and **Mag** should be equivalent to each other.
- Usage of **Mag** is however preferred, **Rob** used for transition from **Tar**.
- Furthermore, problem with Poisson solver and **Rob** BC in COMPASS simulation leading to code slowing.

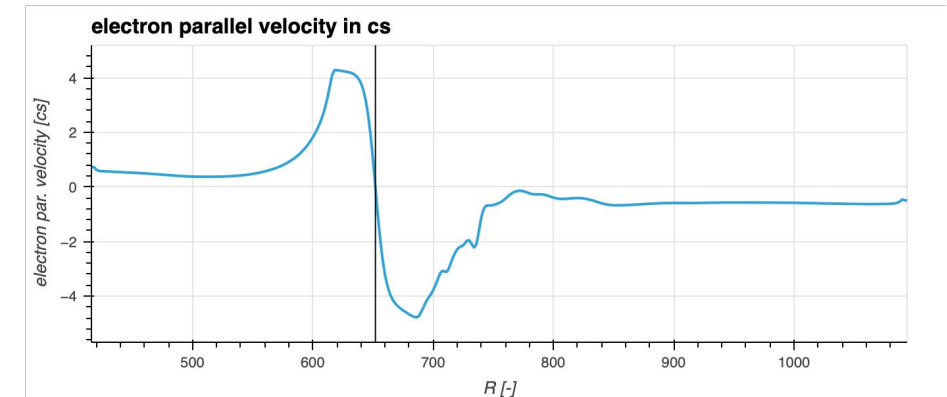


Electron parallel velocity



Electron parallel velocity
Zoomed on FPR

- Significantly higher values ($\sim 3x$) of v_{pare} for **Mag BC**.
- Sign represents direction with respect to magnetic field lines alignment.
- The 1D profile is smoothed by Tanh function in ghost cells where mag. Field is tangent to surface.
- Electron and ion parallel velocity set to 0 between the two regions (plus zero derivative of v_{pari}).



1D profile of electron parallel velocity at the bottom boundary

$$v_{\parallel i} = \pm c_s \sqrt{1 + \frac{T_i}{T_e}},$$

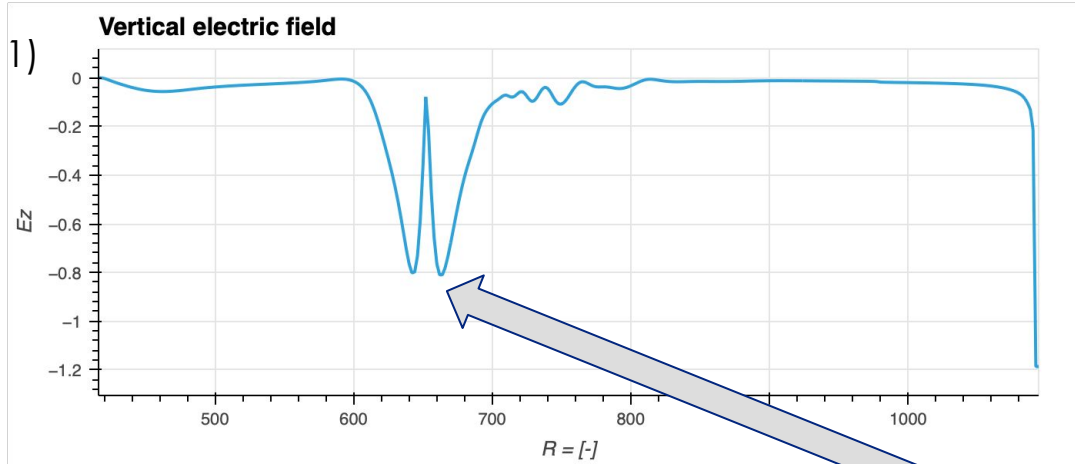
$$v_{\parallel e} = \pm c_s \sqrt{1 + \frac{T_i}{T_e}} \exp\left(\Lambda - \frac{e\phi}{T_e}\right),$$

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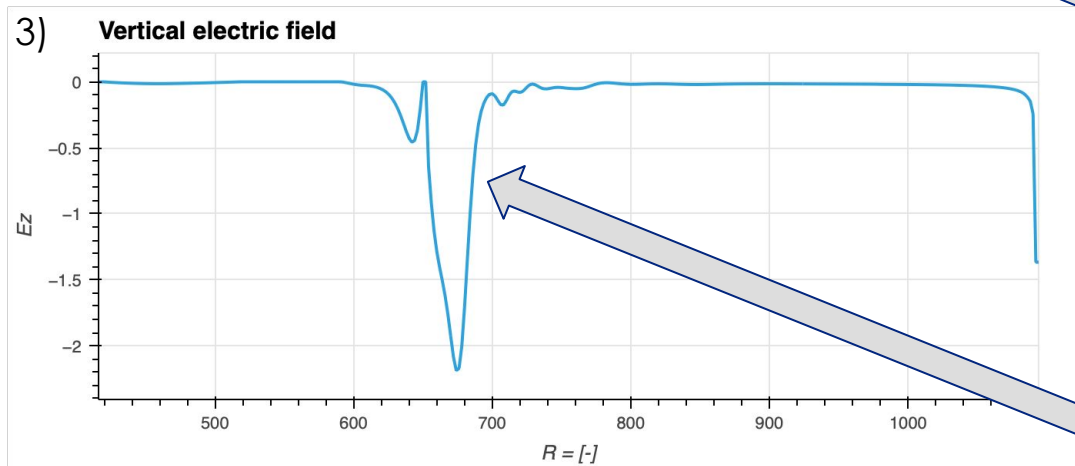
$$\Omega = \mp \frac{m_i n}{e} c_s \sqrt{1 + \frac{T_i}{T_e}} \partial_{ss}^2 v_{\parallel i},$$

$$\partial_s \phi = \mp \frac{m_i c_s}{e \sqrt{1 + \frac{T_i}{T_e}}} \partial_s v_{\parallel i},$$

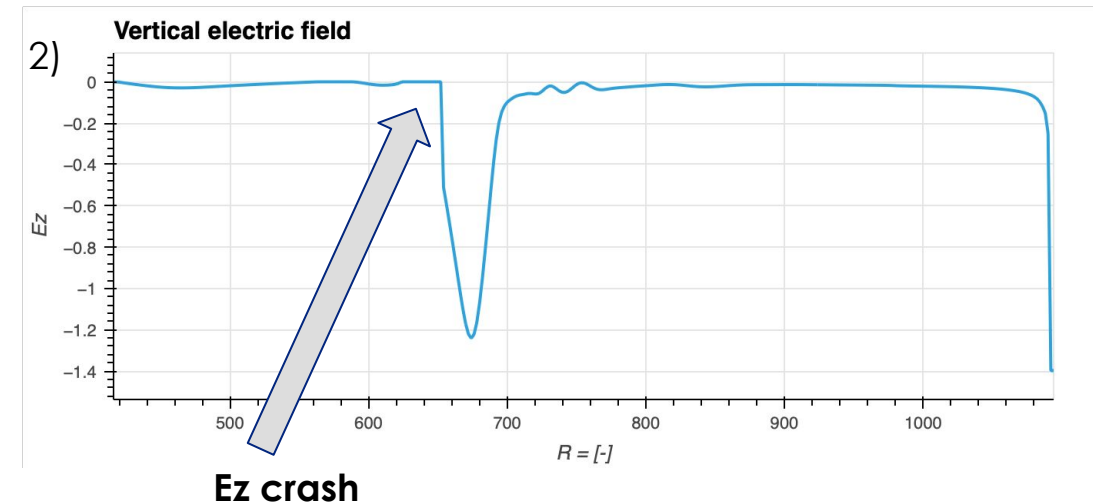


- **E_z calculated between ghost cells!**
- Crash of vertical electric field at inner strike point.
- Hard limit in code cutting $E_z > 0$ to 0
- The point in middle forced to zero by setting derivative of ion parallel velocity to zero.
- Causing parallel electron velocity to increase.
- Leading to crash in consequence.
- Self-refresh after several time units.

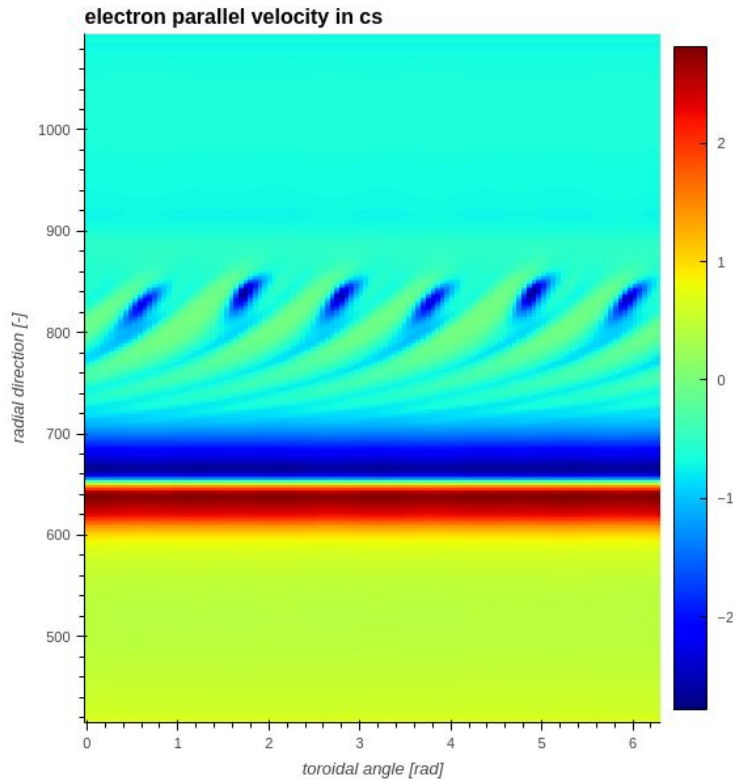
Initial E_z



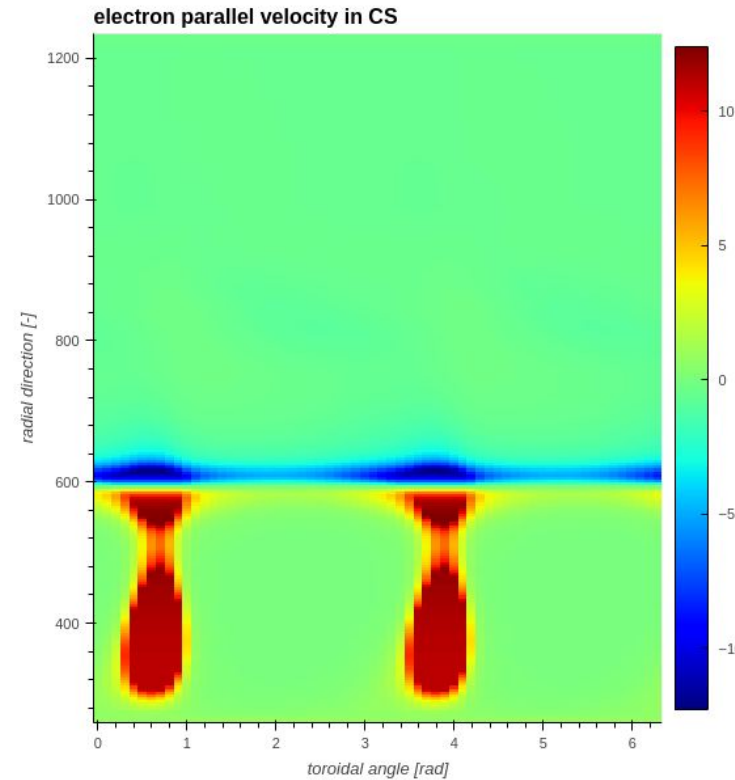
E_z self-refreshed



E_z crash



Normal values



Crash in electron velocity

- The crash in electron parallel velocities, caused by vertical electric field, was threatened by adjusting the mass ratio and the sheath drop Δ .
- The mechanism of positive vertical electric field formation is however still not well understood.
- Already appeared several times using **Mag** BC in past.
- Since now, the problem did not appear again.

SUMMARY

- COMPASS simulation showed significant impact of BCs on simulation:
 - **Tar** condition caused huge amplitudes in potential, exceeding 2 000 V leading to crash.
 - Both **Robin** and **Magnetic** BCs showed similar turbulence properties.
- Problems with Poisson solver combined with **Robin** boundary condition were observed.
- **Magnetic** boundary condition performed best, however, problems with vertical electric field were observed:
 - Increasing E_z caused acceleration of electrons and simulation crash.
 - The problem however fixed itself after enough simulation time and mass ratio adjustment.
 - This mechanism must be further investigated.
 - Since now, the problem with the bottom boundary did not happen again.

1. M. Giacomini et al J. Comput. Phys. 463 (2022) 111294 (The GBS code for the self-consistent simulation of plasma turbulence and kinetic neutral dynamics in the tokamak boundary)
2. M. Giacomini *et al* 2021 *Nucl. Fusion* **61** 076002 (Theory-based scaling laws of near and far scrape-off layer widths in single-null L-mode discharges)



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