

# Mitigation of Alfvén Eigenmodes in Negative Triangularity plasmas at TCV

**P. Oyola,** M. García-Muñoz, M. Vallar, E. Viezzer, J. Rueda-Rueda, J. Domínguez-Palacios, J. Gonzalez-Martin, Y. Todo, S. Sharapov, A. Fasoli, B. Duval, A. Karpushov, S. Coda, O. Sauter and the TCV team.





This work has been carried out within the framework of the EUROfusion Consortium and has received funding from the Euratom research and training programme 2014-2018 and 2019-2020 under grant agreement No 633053. The views and opinions expressed herein do not necessarily reflect those of the European Commission.

- Experimental observations of TAEs in NT
- MEGA: 3D nonlinear hybrid kinetic-MHD
- TAEs in NT and PT

Outline

- Wave-particle resonances in the FI phase-space
- Fast-ion losses induced by TAE in NT and PT





#### Outline

- Experimental observations of TAEs in NT
- MEGA: 3D nonlinear hybrid kinetic-MHD
- TAEs in NT and PT
- Wave-particle resonances in the FI phase-space
- Fast-ion losses induced by TAE in NT and PT





## AEs in NT firstly observed in DIII-D



 Experiments in DIII-D<sup>4</sup> to obtain AEs, shows TAEs excited in NT and PT.



<sup>4</sup> M. A. Van Zeeland *et al.*, NF **59** 086028 (2019)

# Gyrofluid simulations indicate negligible impact on AE activity

 Experiments in DIII-D<sup>4</sup> to obtain AEs, shows TAEs excited in NT and PT.

- Numerical studies<sup>5</sup> with FAR3d<sup>6</sup>:
  - Linear EP-driven AE.
  - 2-moments gyrofluid model for FI
  - Negligible impact of triangularity on AE growth rate



<sup>6</sup>L. A. Charlton et al., J. Comp. Phys 86 270 (1990)

1.0 **(d)** 

 $\phi_{ES}^{n=3}(R,z)$ 

<sub>0.6</sub>(c)

0.4

0.2

0.0

-0.2

-0.6



## Strong NT impact on AEs at TCV



- Strong impact of triangularity on Alfvénic modes:
  - Amplitude reduction
  - Frequency drops
- Uncontrolled changes in many variables:
  - Density rise during NT phase (better confinement)
  - Direct comparison between triangularities is difficult.
- Nonlinear hybrid simulations help unreveal the impact of  $\delta$  in the Alfvén Eigenmodes and induced fast-ion transport.





- Experimental observations of TAEs in NT
- MEGA: 3D nonlinear hybrid kinetic-MHD
- TAEs in NT and PT

Outline

- Wave-particle resonances in the FI phase-space
- Fast-ion losses induced by TAE in NT and PT





9/32

## MEGA<sup>7</sup>: Nonlinear 3D hybrid kinetic-MHD code



#### Bulk plasma

• Full resistive-MHD model.

$$\begin{aligned} \frac{\partial \rho}{\partial t} + \vec{\nabla} \cdot (\rho \vec{v}) &= \vec{\nabla} \cdot \left(\nu_n \vec{\nabla} \rho\right) \\ \frac{\partial \vec{U}}{\partial t} + \left(\vec{v} \cdot \vec{\nabla}\right) \vec{U} &= -\vec{\nabla} p + \left(\vec{J} - \vec{J}_{FI}\right) \times \vec{B} \\ &+ \frac{4}{3} \left(\nu \rho \vec{\nabla} \cdot \vec{v}\right) - \vec{\nabla} \times (\nu \rho \vec{\omega}) \\ \frac{\partial p}{\partial t} + \vec{\nabla} \cdot (p \vec{v}) + (\gamma - 1) p \vec{\nabla} \cdot \vec{v} &= \\ &\vec{\nabla} \cdot \left(\chi \vec{\nabla} (p - p_{eq})\right) \\ &(\gamma - 1) \left[\nu \rho \left(\vec{\nabla} \times \vec{v}\right)^2 + \frac{4}{3} \left(\vec{\nabla} \cdot \vec{v}\right)^2\right] \\ &(\gamma - 1) \eta \left(\vec{J} - \vec{J}_{FI}\right) \cdot \left(\vec{J} - \vec{J}_{eq}\right) \\ \vec{E} &= -\vec{v} \times B + \eta \vec{J} \end{aligned}$$

<sup>7</sup>Y. Todo et al., PoP 5 1321 (1998)

2023 Annual TSVV 2 Workshop - P. Oyola



#### Bulk plasma

• Full resistive-MHD model.

#### **Fast-ions**

- *Kinetic description*: markers sampling distribution function.
- Gyrokinetic equation ( $\delta f$  or *full*-f).

$$\begin{aligned} \frac{\partial \rho}{\partial t} + \vec{\nabla} \cdot (\rho \vec{v}) &= \vec{\nabla} \cdot \left(\nu_n \vec{\nabla} \rho\right) \\ \frac{\partial \vec{U}}{\partial t} + \left(\vec{v} \cdot \vec{\nabla}\right) \vec{U} &= -\vec{\nabla} p + \left(\vec{J} - \vec{J}_{FI}\right) \times \vec{B} \\ &+ \frac{4}{3} \left(\nu \rho \vec{\nabla} \cdot \vec{v}\right) - \vec{\nabla} \times (\nu \rho \vec{\omega}) \end{aligned}$$
$$\begin{aligned} \frac{\partial p}{\partial t} + \vec{\nabla} \cdot (p \vec{v}) + (\gamma - 1) p \vec{\nabla} \cdot \vec{v} &= \\ \vec{\nabla} \cdot \left(\chi \vec{\nabla} (p - p_{eq})\right) \\ \left(\gamma - 1\right) \left[\nu \rho \left(\vec{\nabla} \times \vec{v}\right)^2 + \frac{4}{3} \left(\vec{\nabla} \cdot \vec{v}\right)^2\right] \\ \left(\gamma - 1\right) \eta \left(\vec{J} - \vec{J}_{FI}\right) \cdot \left(\vec{J} - \vec{J}_{eq}\right) \end{aligned}$$

<sup>7</sup>Y. Todo et al., PoP **5** 1321 (1998)



#### Bulk plasma

• Full resistive-MHD model.



#### **Fast-ions**

- *Kinetic description*: markers sampling distribution function.
- Gyrokinetic equation (δf or *full-*f).

$$\begin{aligned} \frac{\partial \rho}{\partial t} + \vec{\nabla} \cdot (\rho \vec{v}) &= \vec{\nabla} \cdot \left(\nu_n \vec{\nabla} \rho\right) \\ \frac{\partial \vec{U}}{\partial t} + \left(\vec{v} \cdot \vec{\nabla}\right) \vec{U} &= -\vec{\nabla} p + \left(\vec{J} - \vec{J}_{FI}\right) \times \vec{B} \\ &+ \frac{4}{3} \left(\nu \rho \vec{\nabla} \cdot \vec{v}\right) - \vec{\nabla} \times (\nu \rho \vec{\omega}) \end{aligned}$$
$$\begin{aligned} \frac{\partial p}{\partial t} + \vec{\nabla} \cdot (p \vec{v}) + (\gamma - 1) p \vec{\nabla} \cdot \vec{v} = \end{aligned}$$

$$\begin{aligned} \frac{\partial p}{\partial t} + \vec{\nabla} \cdot (p\vec{v}) + (\gamma - 1) p\vec{\nabla} \cdot \vec{v} &= \\ \vec{\nabla} \cdot \left(\chi \vec{\nabla} (p - p_{eq})\right) \\ (\gamma - 1) \left[\nu \rho \left(\vec{\nabla} \times \vec{v}\right)^2 + \frac{4}{3} \left(\vec{\nabla} \cdot \vec{v}\right)^2\right] \\ (\gamma - 1) \eta \left(\vec{J} - \vec{J}_{FI}\right) \cdot \left(\vec{J} - \vec{J}_{eq}\right) \\ \vec{E} &= -\vec{v} \times B + \eta \vec{J} \end{aligned}$$

<sup>7</sup>Y. Todo et al., PoP 5 1321 (1998)



#### Bulk plasma

• Full resistive-MHD model.



#### **Fast-ions**

- *Kinetic description*: markers sampling distribution function.
- Gyrokinetic equation ( $\delta f$  or *full*-f).
- 4<sup>th</sup> order finite differences in cylindrical coordinates (R, φ, z).

$$\frac{\partial \rho}{\partial t} + \vec{\nabla} \cdot (\rho \vec{v}) = \vec{\nabla} \cdot \left(\nu_n \vec{\nabla} \rho\right)$$
$$\frac{\partial \vec{U}}{\partial t} + \left(\vec{v} \cdot \vec{\nabla}\right) \vec{U} = -\vec{\nabla} p + \left(\vec{J} - \vec{J}_{FI}\right) \times \vec{B}$$
$$+ \frac{4}{3} \left(\nu \rho \vec{\nabla} \cdot \vec{v}\right) - \vec{\nabla} \times (\nu \rho \vec{\omega})$$

$$\begin{aligned} \frac{\partial p}{\partial t} + \vec{\nabla} \cdot (p\vec{v}) + (\gamma - 1) p\vec{\nabla} \cdot \vec{v} &= \\ \vec{\nabla} \cdot \left(\chi \vec{\nabla} (p - p_{eq})\right) \\ (\gamma - 1) \left[\nu \rho \left(\vec{\nabla} \times \vec{v}\right)^2 + \frac{4}{3} \left(\vec{\nabla} \cdot \vec{v}\right)^2\right] \\ (\gamma - 1) \eta \left(\vec{J} - \vec{J}_{FI}\right) \cdot \left(\vec{J} - \vec{J}_{eq}\right) \\ \vec{E} &= -\vec{v} \times B + \eta \vec{J} \end{aligned}$$

<sup>7</sup>Y. Todo et al., PoP 5 1321 (1998)



#### Bulk plasma

• Full resistive-MHD model.



#### **Fast-ions**

- *Kinetic description*: markers sampling distribution function.
- Gyrokinetic equation (δf or *full-*f).
- 4<sup>th</sup> order finite differences in cylindrical coordinates (R, φ, z).
- Explicit 4<sup>th</sup> Runge-Kutta for time-integration.

$$\frac{\partial \rho}{\partial t} + \vec{\nabla} \cdot (\rho \vec{v}) = \vec{\nabla} \cdot \left(\nu_n \vec{\nabla} \rho\right)$$

$$\frac{\partial \vec{U}}{\partial t} + \left(\vec{v} \cdot \vec{\nabla}\right) \vec{U} = -\vec{\nabla} p + \left(\vec{J} - \vec{J}_{FI}\right) \times \vec{B}$$

$$+ \frac{4}{3} \left(\nu \rho \vec{\nabla} \cdot \vec{v}\right) - \vec{\nabla} \times (\nu \rho \vec{\omega})$$

$$\begin{aligned} \frac{\partial p}{\partial t} + \vec{\nabla} \cdot (p\vec{v}) + (\gamma - 1) p\vec{\nabla} \cdot \vec{v} &= \\ \vec{\nabla} \cdot \left(\chi \vec{\nabla} (p - p_{eq})\right) \\ (\gamma - 1) \left[\nu \rho \left(\vec{\nabla} \times \vec{v}\right)^2 + \frac{4}{3} \left(\vec{\nabla} \cdot \vec{v}\right)^2\right] \\ (\gamma - 1) \eta \left(\vec{J} - \vec{J}_{FI}\right) \cdot \left(\vec{J} - \vec{J}_{eq}\right) \\ \vec{E} &= -\vec{v} \times B + \eta \vec{J} \end{aligned}$$

<sup>7</sup> Y. Todo *et al.*, PoP **5** 1321 (1998)

## Simulation setup for the $\delta$ comparison



• Flipped equilibrium to isolate the  $+\delta$  /  $-\delta$  effects on AE activity<sup>8</sup>.

#### **Simulation parameters**

- $\delta$ f-method for kinetic species.
- #markers = 23M particles
- Multi-*n* simulation (n < 5)



<sup>&</sup>lt;sup>8</sup> P. Oyola et al., in preparation

18 July 2023

2023 Annual TSVV 2 Workshop - P. Oyola



Analytical anisotropic slowing-down distribution

$$f_0 \propto e^{-\frac{(\rho-\rho_0)^2}{2(\Delta\rho_0)^2}} \frac{1}{v^3 + v_{crit}^3} erfc\left(\frac{v - v_{birth}}{\Delta v}\right) e^{-\frac{(\Lambda-\Lambda_0)^2}{2(\Delta\Lambda)^2}}$$



Analytical anisotropic slowing-down distribution

$$f_0 \propto e^{-\frac{(\rho-\rho_0)^2}{2(\Delta\rho_0)^2}} \frac{1}{v^3 + v_{crit}^3} erfc\left(\frac{v - v_{birth}}{\Delta v}\right) e^{-\frac{(\Lambda-\Lambda_0)^2}{2(\Delta\Lambda)^2}}$$



Analytical anisotropic slowing-down distribution

$$f_0 \propto e^{-\frac{(\rho-\rho_0)^2}{2(\Delta\rho_0)^2}} \frac{1}{v^3 + v_{crit}^3} erfc\left(\frac{v-v_{birth}}{\Delta v}\right) e^{-\frac{(\Lambda-\Lambda_0)^2}{2(\Delta\Lambda)^2}}$$







18 July 2023

2023 Annual TSVV 2 Workshop - P. Oyola

14 / 32



Analytical anisotropic slowing-down distribution

$$f_0 \propto e^{-\frac{(\rho - \rho_0)^2}{2(\Delta \rho_0)^2}} \frac{1}{v^3 + v_{crit}^3} erfc\left(\frac{v - v_{birth}}{\Delta v}\right) e^{-\frac{(\Lambda - \Lambda_0)^2}{2(\Delta \Lambda)^2}}$$

- Scan in different pitch-angle injections  $\Lambda_0 \equiv \frac{\mu B_{\rm axis}}{E}$
- Scan in different fast-ion gradient location



 $\rho_0$ 

## Initial FI drive is the same for NT and PT

Analytical slowing-down distribution:

$$f_0 \propto e^{-\frac{(\rho-\rho_0)^2}{2(\Delta\rho_0)^2}} \frac{1}{v^3 + v_{crit}^3} erfc\left(\frac{v-v_{birth}}{\Delta v}\right) e^{-\frac{(\Lambda-\Lambda_0)^2}{2(\Delta\Lambda)^2}}$$

- Scan in different pitch-angle injections  $\Lambda_0 \equiv \frac{\mu B_{axis}}{E}$
- Scan in different fast-ion gradient location

$$\gamma_{TAE} \propto \beta_{FI} \left( \frac{\partial f_0}{\partial E} + \frac{n}{\omega} \frac{\partial f_0}{\partial P_{\phi}} \right)$$





2023 Annual TSVV 2 Workshop - P. Oyola

 $\rho_0$ 

#### Outline

- Why Negative Triangularity ?
- MEGA: 3D nonlinear hybrid kinetic-MHD
- TAEs in NT and PT
- Wave-particle resonances in the FI phase-space
- Fast-ion losses induced by TAE in NT and PT





## TAEs is mitigated in NT vs PT



TAEs appear both in PT and NT:

• PT reaches an energy ~40% higher.



- PT reaches an energy ~40% higher.
- SAW is similar in PT & NT.





## Linear growth rate for NT and PT



- PT reaches an energy ~40% higher.
- SAW is similar in PT & NT.
- NT shows a smaller growth rate.



## Linear growth rate for NT and PT



- PT reaches an energy ~40% higher.
- SAW is similar in PT & NT.
- NT shows a smaller growth rate.
- This trend is kept independent on the pair of parameters  $\Lambda_0$  and  $\rho_0$ .



## Linear growth rate for NT and PT



- PT reaches an energy ~40% higher.
- SAW is similar in PT & NT.
- NT shows a smaller growth rate.
- This trend is kept independent on the pair of parameters  $\Lambda_0$  and  $\rho_0$ .



- Experimental observations of TAEs in NT
- MEGA: 3D nonlinear hybrid kinetic-MHD
- TAEs in NT and PT

Outline

- Wave-particle resonances in the FI phase-space
- Fast-ion losses induced by TAE in NT and PT





21/32

Resonant energy exchange in FI phase-space

• Power exchange in FI phase-space shows particlewave resonances.

 $\Delta E > 0 \longrightarrow$  Energy to the FI

 $\Delta E < 0 \longrightarrow$  Energy to the wave

• Two main regions of the phase-space providing energy to TAE.





2023 Annual TSVV 2 Workshop - P. Oyola

Resonant energy exchange in FI phase-space

 Power exchange in FI phase-space shows particlewave resonances.

 $\Delta E > 0 \longrightarrow$  Energy to the FI

 $\Delta E < 0 \longrightarrow$  Energy to the wave

- Two main regions of the phase-space providing energy to TAE:
  - Wave-particle resonances<sup>9</sup>.





<sup>9</sup>Y. Todo, Rev. Mod. Plasma Phys **3**, 1 (2019)

#### NT damps the lower bounce harmonic

- Alignment of analytical resonances with structures in FI phase-space.
- In PT, lower transit harmonic is most excited.
- In NT, damps lower transit harmonics.





#### NT damps the lower bounce harmonic.

- Alignment of analytical resonances with structures in FI phase-space.
- In PT, lower transit harmonic is most excited.
- In NT, damps lower transit harmonics.
- Overall energy transfer is larger in PT.





- Experimental observations of TAEs in NT
- MEGA: 3D nonlinear hybrid kinetic-MHD
- TAEs in NT and PT

Outline

- Wave-particle resonances in the FI phase-space
- Fast-ion losses induced by TAE in NT and PT





## Synthetic wall in MEGA<sup>10</sup>



• 2D wall implemented in MEGA for TCV tokamak.



<sup>10</sup> P. Oyola *et al.*, RSI **92** (2021)

#### TAE-induced FIL are 3x lower in NT



- 2D wall implemented in MEGA for TCV tokamak.
- Correlated FIL bursts with TAE saturation.



### TAE-induced FIL are 3x lower in NT



- 2D wall implemented in MEGA for TCV tokamak.
- Correlated FIL bursts with TAE saturation.

- Fast-ion losses in NT is **smaller** than its counterpart in PT.
  - 3x times lower at the peak.
  - 3x times lower integrated FIL.



### TAE-induced FIL are 3x lower in NT



- 2D wall implemented in MEGA for TCV tokamak.
- Correlated FIL bursts with TAE saturation.

- Fast-ion losses in NT is **smaller** than its counterpart in PT.
  - 3x times lower at the peak.
  - 3x times lower integrated FIL.



#### Conclusions & Outlook

- In experiments, TAEs appear weaker in NT than in its counterpart PT.
- MEGA sims used to isolate the  $\delta$  effects.
- ~40% lower energy in NT with respect to PT.
- Lower transit harmonics are damped in NT.
- Fast-ion losses are 3x lower in NT.

2023 Annual TSVV 2 Workshop - P. Oyola





**Conclusions & Outlook** 

Work-in-progress

- Extend the scan to several  $\delta \in (-0.6, 0.6)$ .
- Studying the particle transport by the TAE in both PT and NT.





Height (m)



#### **Conclusions & Outlook**

- In experiments, TAEs appear weaker in NT than in its counterpart PT.
- MEGA sims used to isolate the  $\delta$  effects. •
- ~40% lower energy in NT with respect to PT.
- Lower transit harmonics are damped in NT. •
- Fast-ion losses are 3x lower in NT.









## **Backup slides**





This work has been carried out within the framework of the EUROfusion Consortium and has received funding from the Euratom research and training programme 2014-2018 and 2019-2020 under grant agreement No 633053. The views and opinions expressed herein do not necessarily reflect those of the European Commission.