



Towards first-principles simulations of the L- to H-mode transition with the global gyrokinetic turbulence code GENE-X



TSVV4/1 Workshop 2024 / H-mode Workshop 2024

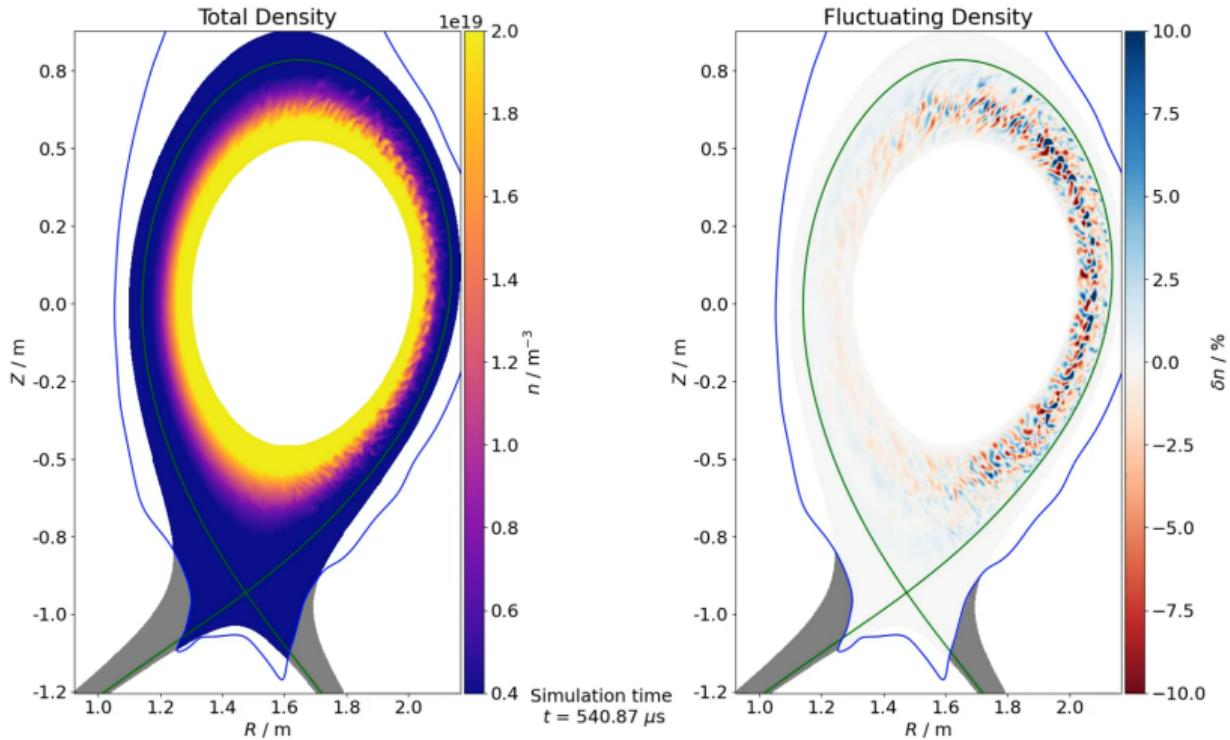
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³See author list of H. Zohm et al, 2024 Nucl. Fusion <https://doi.org/10.1088/1741-4326/ad249d>

Example: Simulation of turbulence in AUG L-mode



In this talk we investigate the transition between turbulent states in AUG by performing a power ramp on an L-mode

Part I

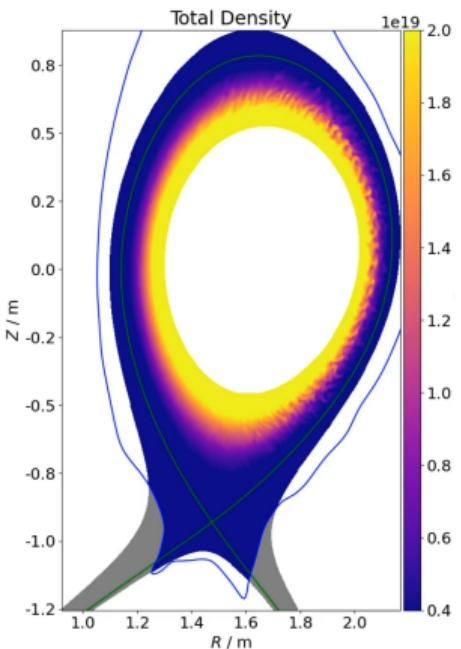
- Overview of the GENE-X code

Part II

- Previous validation studies in AUG and TCV

Part III

- Turbulence simulations in AUG with a power ramp

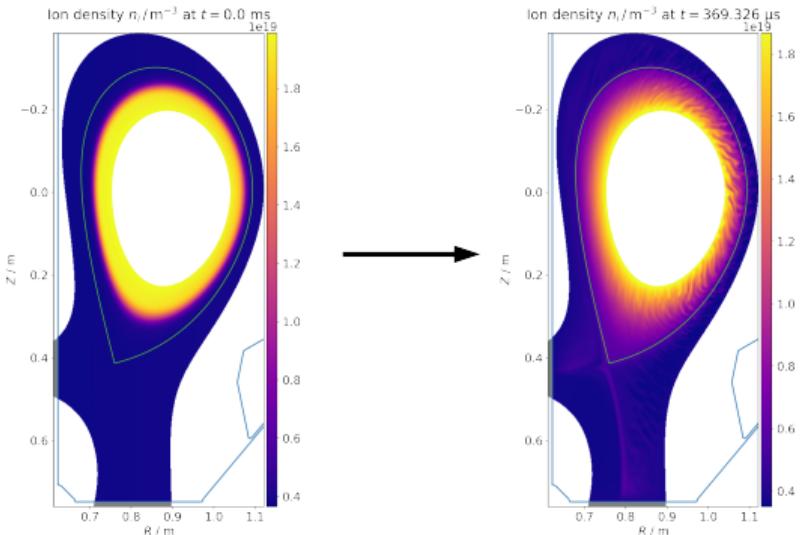


GENE-X Overview

GENE-X enables gyrokinetic turbulence simulations in X-point geometries

Features:

- grid-based (Eulerian)
- global
- non-linear
- full-f
- electromagnetic (EM)
- collisional



GENE-X can simulate from the core to the wall.

Efficiently designed for massively parallelized conventional (CPU-based) supercomputers. Strong scaling with 93% efficiency up to 512 nodes ($\approx 20k$ cores).

[D. Michels, A. Stegmeir, P. Ulbl et al. CPC 264 (2021)] [D. Michels, P. Ulbl et al. PoP 29 (2022)] [P. Ulbl, T. Body et al. PoP 30 (2023)]

GENE-X solves a full-f, collisional, EM gyrokinetic model

evolution

B* advection

perpendicular drifts

$$\frac{\partial f_\alpha}{\partial t} + v_{||} \frac{\mathbf{B}^*}{B_{||}^*} \cdot \nabla f_\alpha + \frac{c}{q_\alpha B_{||}^*} \mathbf{b} \times (\mu \nabla B + q_\alpha \nabla \phi_1 + \nabla H_2) \cdot \nabla f_\alpha$$

$$- \frac{\mathbf{B}^*}{m_\alpha B_{||}^*} \cdot (\mu \nabla B + q_\alpha \nabla \phi_1 + \nabla H_2) \frac{\partial f_\alpha}{\partial v_{||}} - \frac{q_\alpha}{m_\alpha c} \frac{\partial A_{1,||}}{\partial t} \frac{\partial f_\alpha}{\partial v_{||}} = \sum_\beta C_{\alpha\beta}(f_\alpha, f_\beta)$$

vspace-advection

magnetic induction

collisions

with

$$-\nabla \cdot \left(\sum_\alpha \frac{m_\alpha c^2 n_\alpha}{B^2} \nabla_\perp \phi_1 \right) = \sum_\alpha q_\alpha \int f_\alpha \, dW, \quad -\Delta_\perp A_{1,||} = 4\pi \sum_\alpha \frac{q_\alpha}{c} \int v_{||} f_\alpha \, dW.$$

$$\mathbf{B}^* = \mathbf{B} + \frac{m_\alpha c}{q_\alpha} v_{||} \nabla \times \mathbf{b} + \nabla A_{1,||} \times \mathbf{b}, \quad H_2 = -\frac{m_\alpha c^2}{2B^2} |\nabla_\perp \phi_1|^2, \quad dW = 2\pi B_{||}^* / m_\alpha dv_{||} d\mu.$$

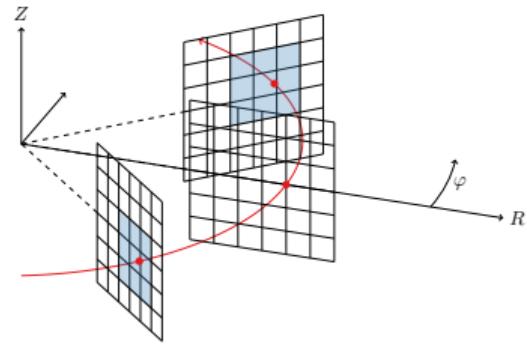
Based on [D. Michels, P. Ulbl, W. Zholobenko et al., PoP 29 (2022) 032307].

The Flux-Coordinate Independent Approach (FCI) allows for simulations in X-point geometries

Plasma turbulence is field aligned, **but** conventional field aligned coordinates break down at the X-point.

Solution: FCI approach

- Collection of Cartesian poloidal planes.
 - Connected with magnetic field lines.
- **locally field aligned coordinates.**



Discretization:

- 4th order sym. FD for $x, z, v_{||}$ derivatives
- 2nd order Arakawa for non-linear terms
- 2nd order elliptic solvers (GMRES + multigrid))



- Field line tracing
- Bicubic interpolation
- 4th order sym. FD for y



Bhatnagar-Gross-Krook (BGK)

$$C_{\alpha\beta} = \nu_{\alpha\beta} \left(\frac{B}{B_{||}^*} \mathcal{M}_{\alpha\beta} - f_\alpha \right).$$

Lenard-Bernstein/Dougherty (LBD)

$$C_{\alpha\beta} = \frac{\nu_{\alpha\beta}}{B_{||}^*} \left\{ \frac{\partial}{\partial v_{||}} \left[(v_{||} - u_{||,\alpha\beta}) B_{||}^* f_\alpha + \frac{1}{2} v_{\text{th},\alpha\beta}^2 \frac{\partial B_{||}^* f_\alpha}{\partial v_{||}} \right] + \frac{\partial}{\partial \mu} \left[2\mu B_{||}^* f_\alpha + \frac{m_\alpha v_{\text{th},\alpha\beta}^2}{B} \mu \frac{\partial B_{||}^* f_\alpha}{\partial \mu} \right] \right\}.$$

- Conservative finite volume discretization (2nd order)

Fokker-Planck/Landau (FPL)

$$C_{\alpha\beta} = -\frac{1}{B_{||}^*} \frac{\partial}{\partial \mathbf{v}} \cdot \left[B_{||}^* \left(\mathbf{E}_{\alpha\beta} f_\alpha + \mathbf{D}_{\alpha\beta} \cdot \frac{\partial}{\partial \mathbf{v}} f_\alpha \right) \right],$$

$$\mathbf{E}_{\alpha\beta} = \frac{\Gamma_{\alpha\beta}}{m_\beta} \int d\mathbf{v}' \mathbf{U}_{\alpha\beta}^E \cdot \frac{\partial}{\partial \mathbf{v}'} f'_\beta,$$

$$\mathbf{D}_{\alpha\beta} = -\frac{\Gamma_{\alpha\beta}}{m_\alpha} \int d\mathbf{v}' \mathbf{U}_{\alpha\beta}^D f'_\beta,$$

$$\mathbf{U}_{\alpha\beta}^E = \begin{pmatrix} U_{\perp,\perp} & U_{\perp,||} \\ U_{||,\perp} & U_{||,||} \end{pmatrix}, \quad \mathbf{U}_{\alpha\beta}^D = \begin{pmatrix} U_{\perp,\perp} & U_{\perp,||} \\ U_{||,\perp} & U_{||,||} \end{pmatrix}.$$

- Components of \mathbf{U} are given by linear combinations of elliptic integrals $E(m)$, $K(m)$
- 2nd order FV discretization

See [R. Hager et. al, JCP315 (2016)]

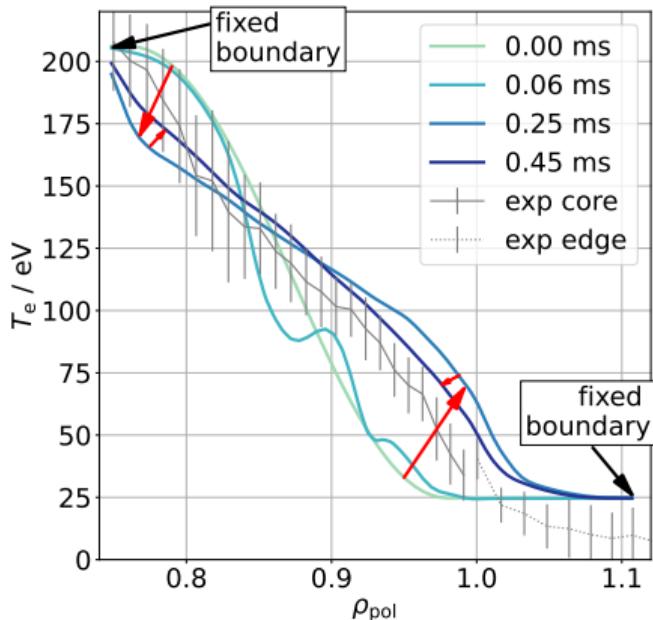
Boundary conditions provide heat and particle fluxes in our simulations

"First-principles" modelling: we start from an arbitrary initial state, not imposing the experimental profiles as a whole.

GENE-X BCs

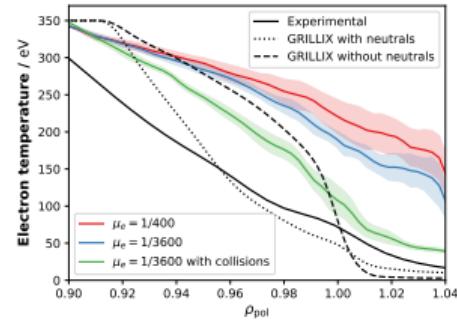
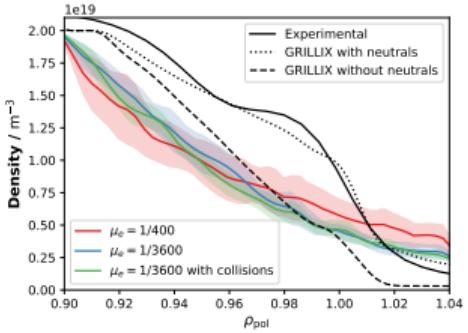
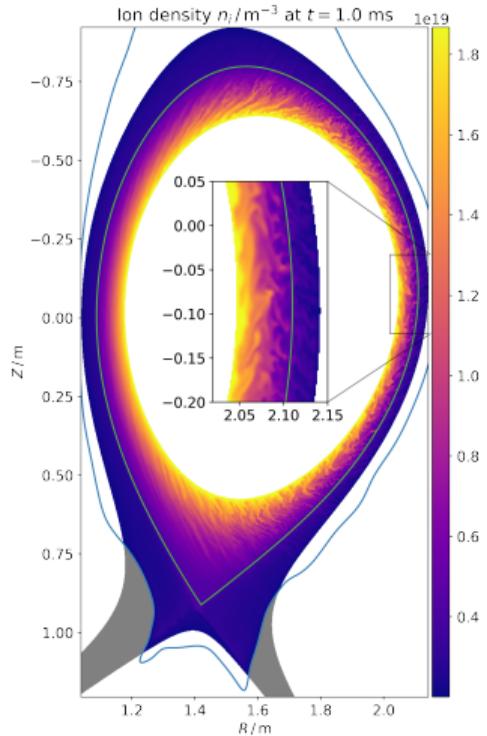
Dirichlet with

- **Distribution function:** Maxwellian with experimental profile values for n and T (no flow)
- **Potentials:** zero



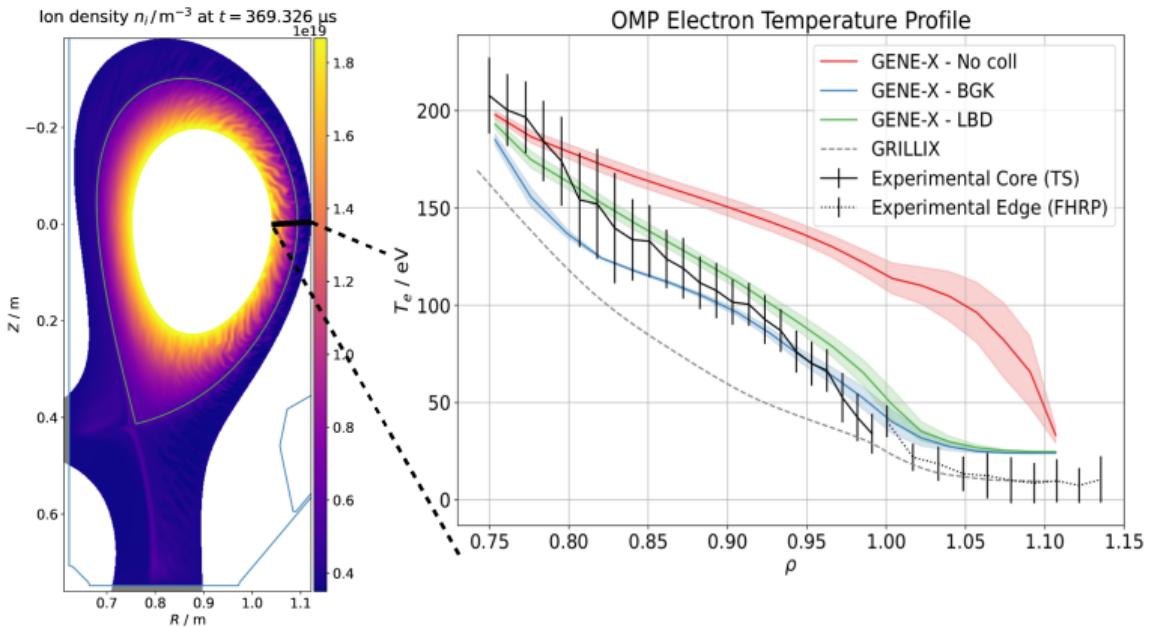
Previous Validation Studies with GENE-X

ASDEX Upgrade Simulations show that realistic mass ratio and collisions are required



- Reduced/realistic mass ratio and BGK coll.
- Electron cooling by collisions
- SOL fall-off length λ_q broadened by collisions
- Case where neutrals are important, electron temperature set too high

Simulation with LBD collisions reproduces experimental electron temperature profile in TCV-X21

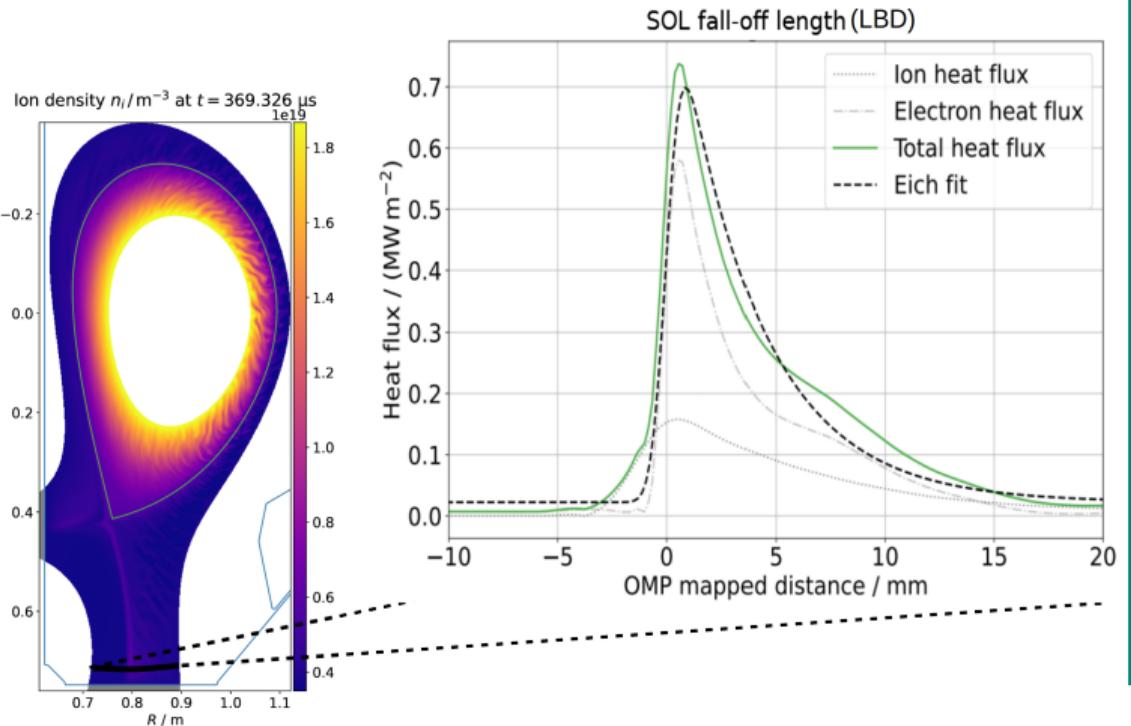


- Realistic T_e profile reproduced by global *collisional* gyrokinetic simulations
- Collisional de-trapping of trapped electrons in the SOL is essential
- Results and data published open access

[P. Ulbl et al. PoP 30 (2023)]

[Data: 10.5281/zenodo.7894731]

Collisional simulations develop "lambda-q" close to experiment



- Divertor heat flux fall-off follows Eich-fit function
- SOL fall-off length λ_q
Experimental: **5.5 mm***

Fluid Models*

GRILLIX	1.1 mm
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*[D. S. Olivera, T. Body et al. NF 62 (2022)]

GENE-X (Gyrokinetic)

No Coll	1.34 mm
Coll BGK	4.68 mm
Coll LBD	3.75 mm

[P. Ulbl et al. PoP 30 (2023)]

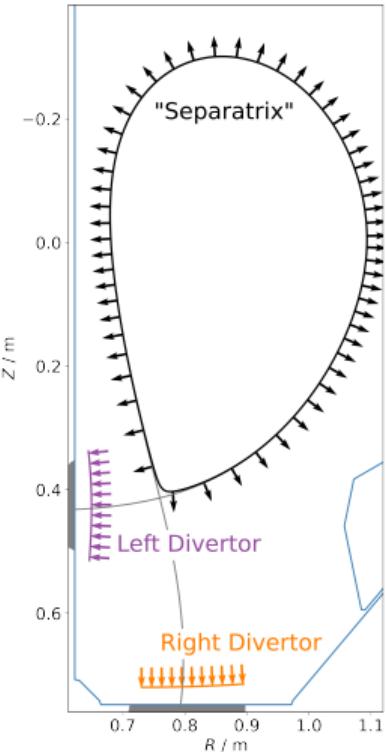
[Data: 10.5281/zenodo.7894731]

A consistent power balance is only achieved using LBD

Q_r / kW	Exp	No Coll	BGK	LBD
Separatrix	120	394	35.6	131.5
Both Divertors	-	136.5	48.8	135.6
Right Divertor	38.1	101.3	51.6	68.6
Left Divertor	-	35.2	-2.8	67.0

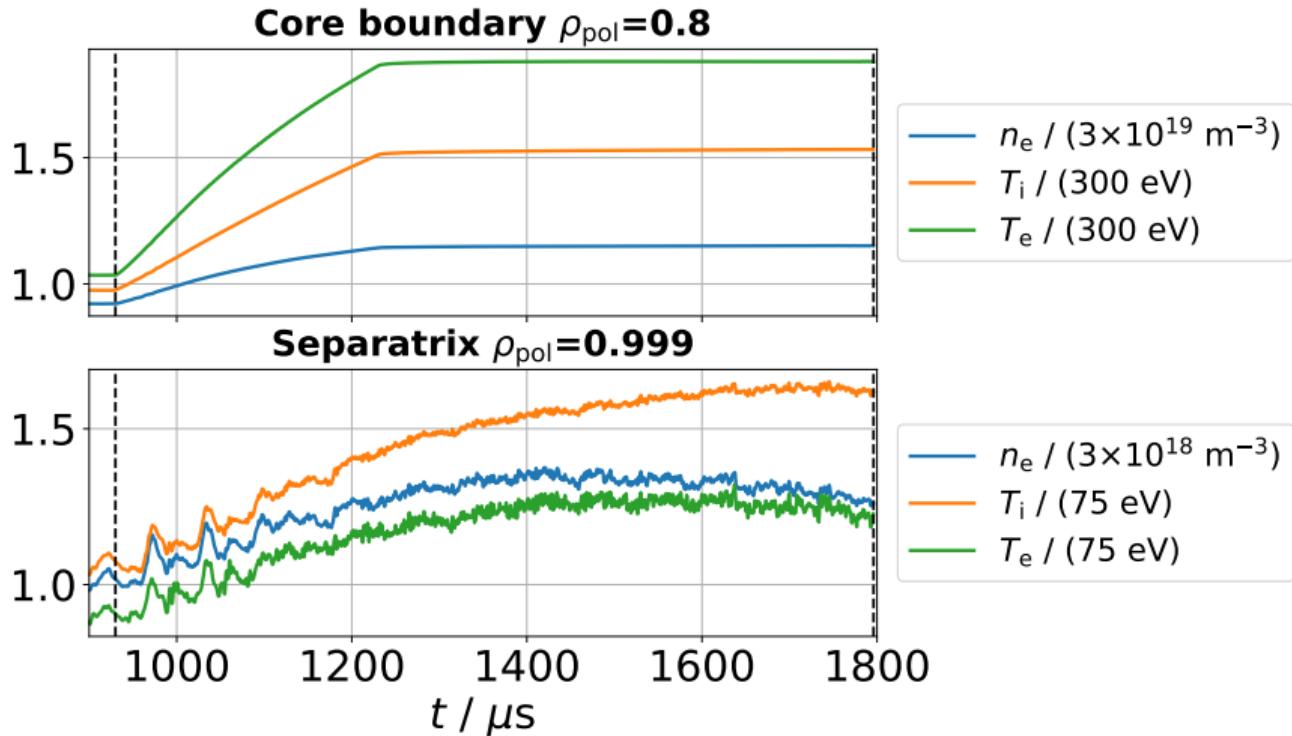
Takeaways

- Separatrix power matches experiment within 10%
- LBD simulation has consistent power balance
- No neutrals and radiation yet

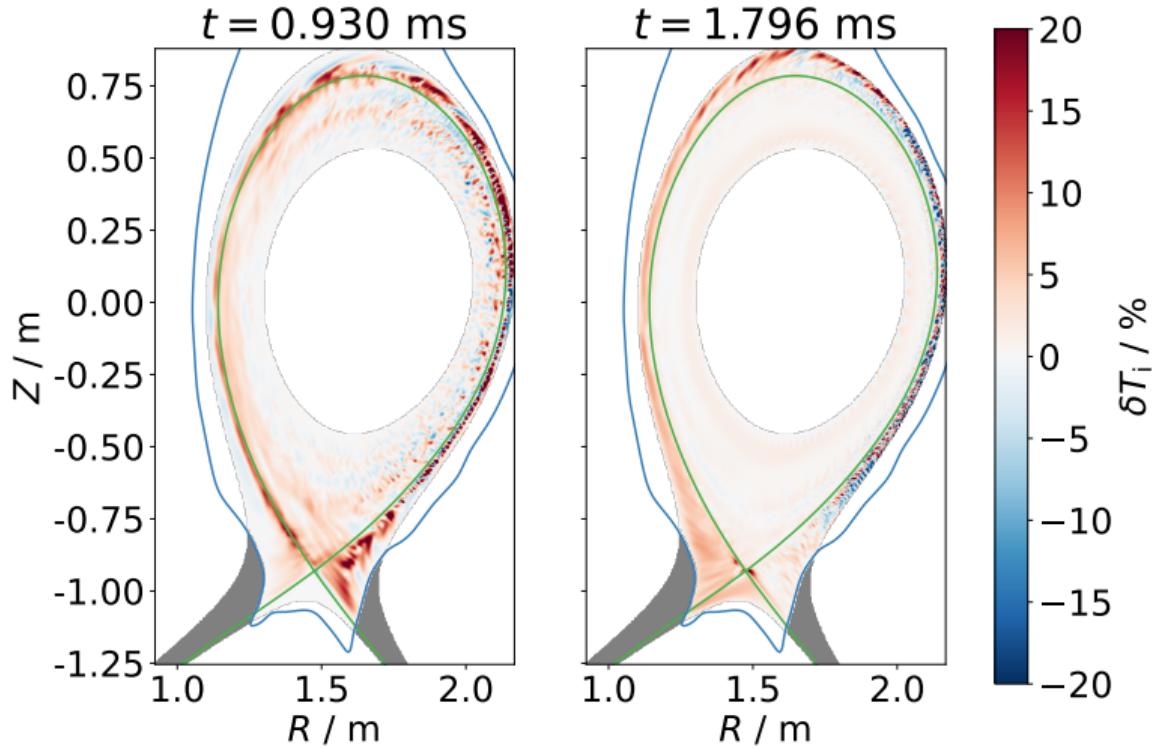


Turbulence simulations with a power ramp

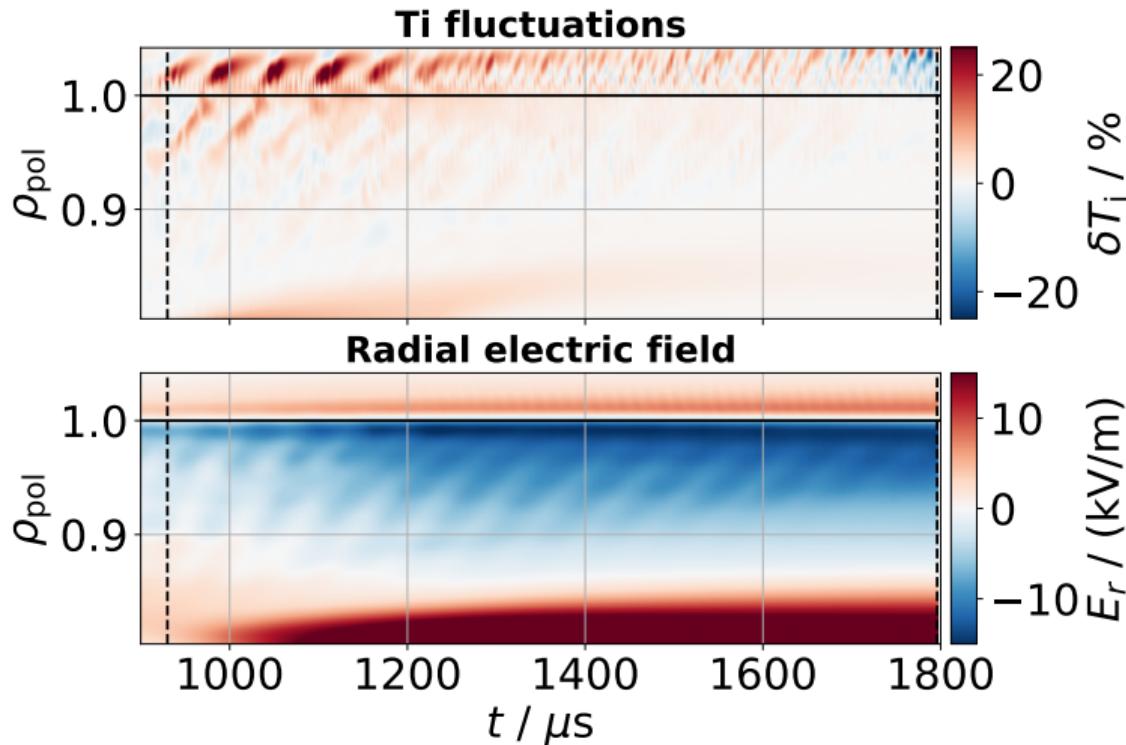
We implement a power ramp by adjusting the inner BCs



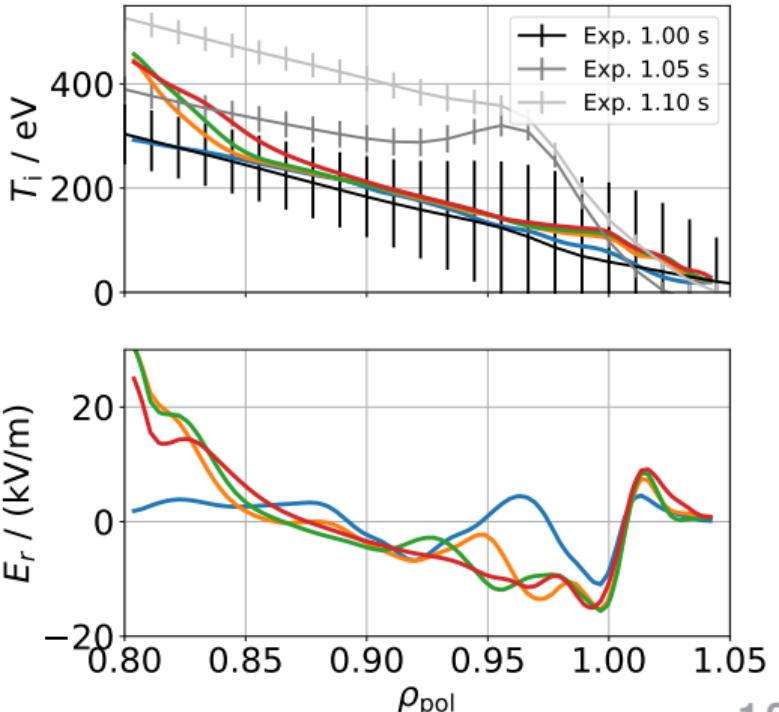
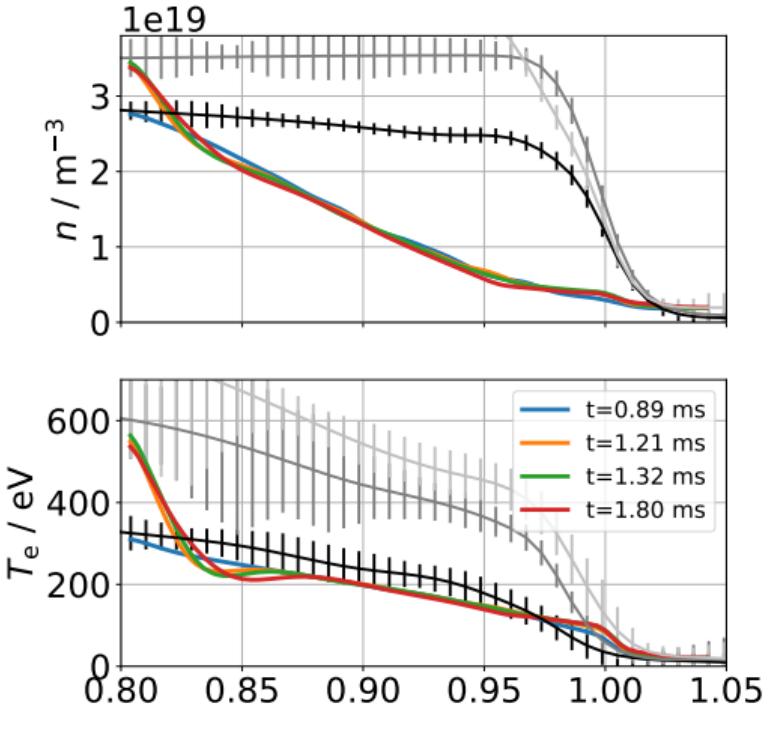
We observe a transition into a turbulence-suppressed state



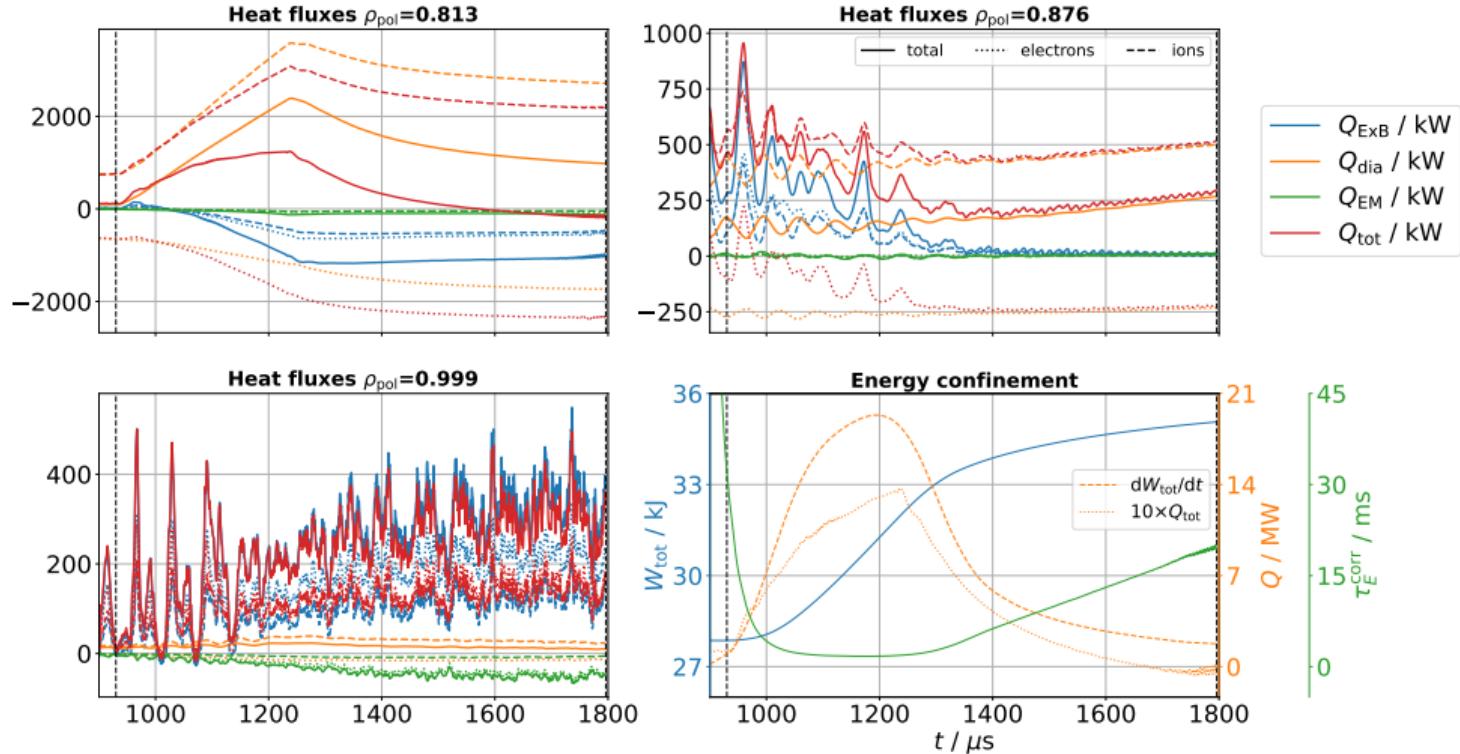
Fluctuations within $\rho < 1$ are suppressed, E_r well builds up



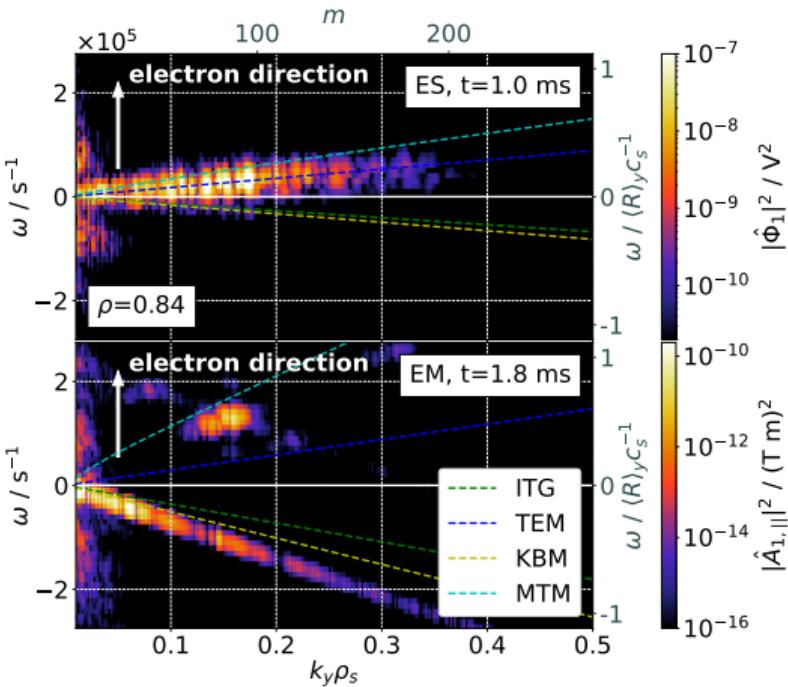
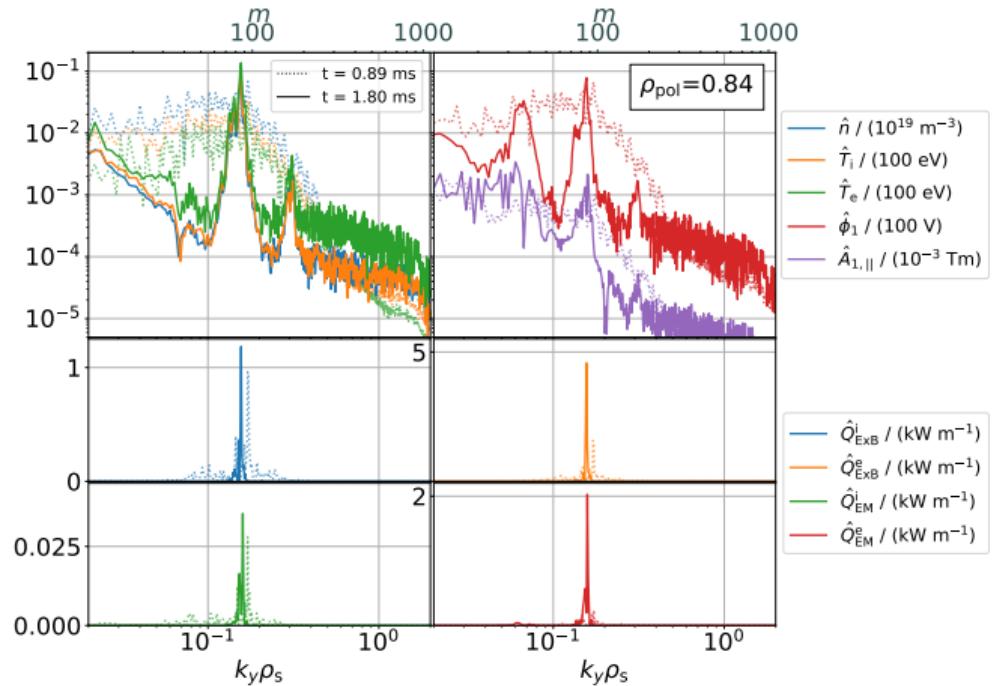
Low density regime due to missing neutral gas particle source close to the separatrix → H-mode highly unlikely



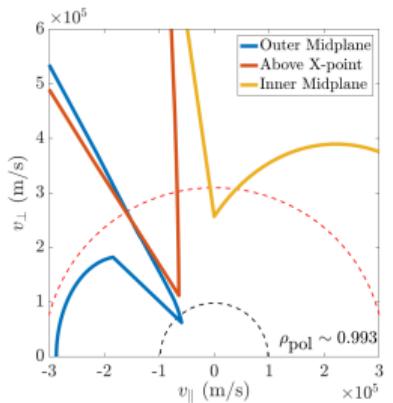
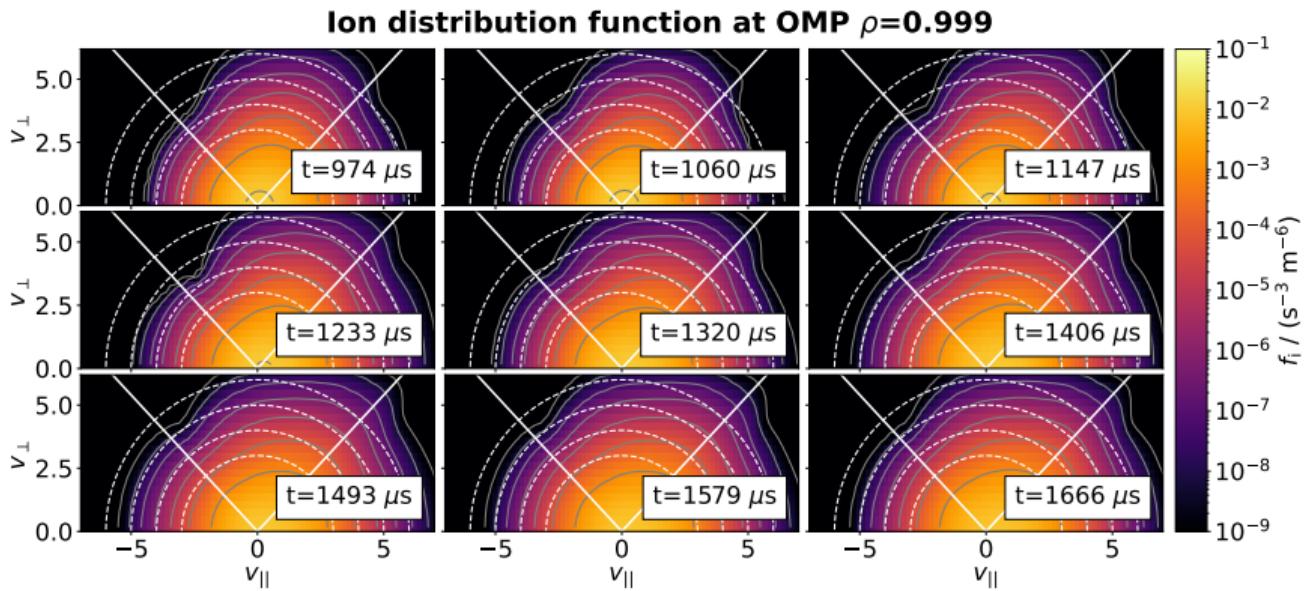
Transport changes locally - NEO and EM become important



Turbulence characteristics change locally - TEM to MTM/KBM?

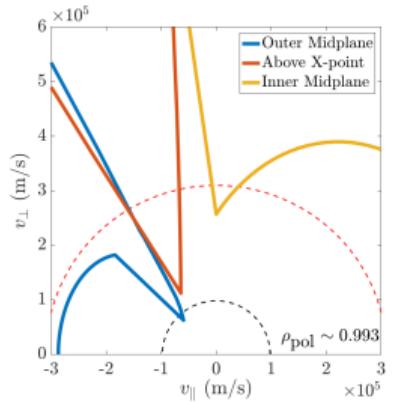
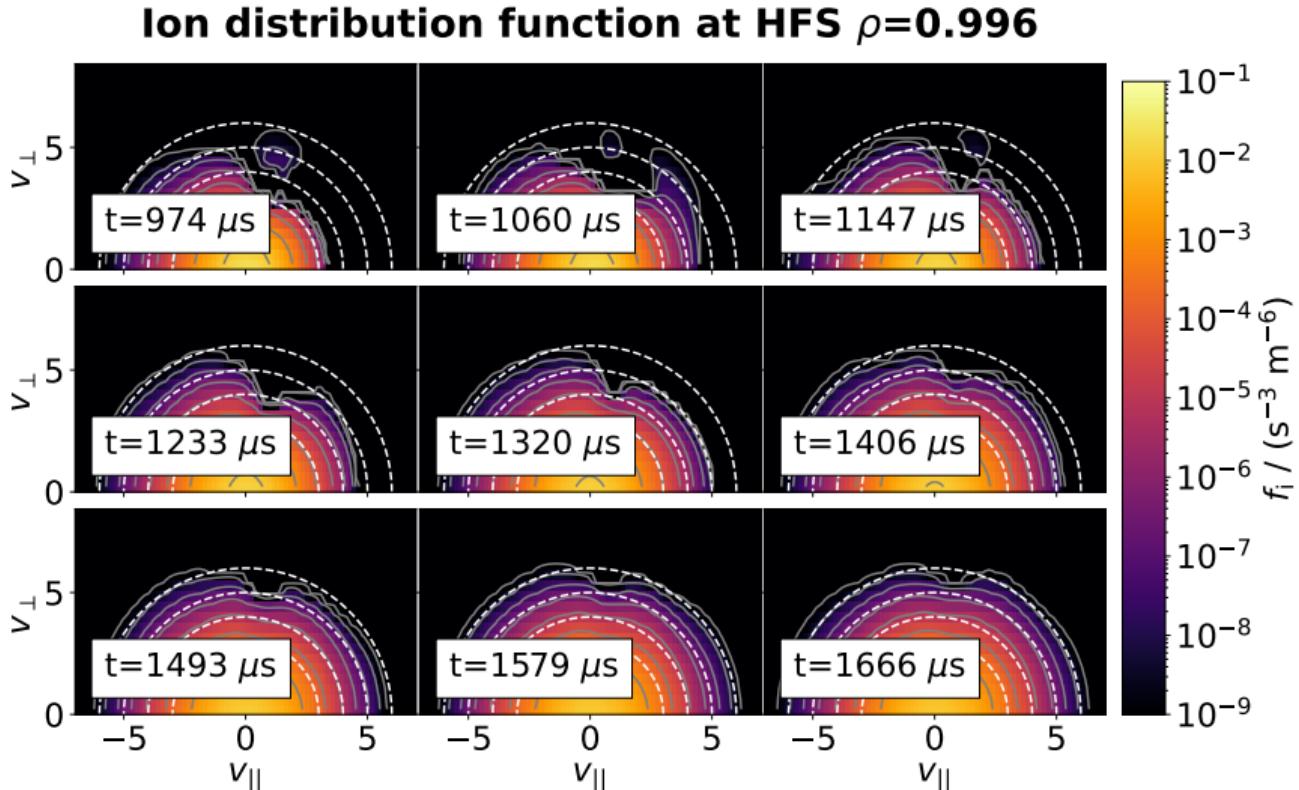


Indications for ion orbit losses (IOL) close to separatrix → diagnostic capabilities are required



[R. Brzozowski, PhD Thesis]

Shape of loss region at HFS differs to LFS in line with theory



[R. Brzozowski, PhD Thesis]

Summary & Outlook

Low-density turbulence simulations with power ramp show transition into turbulence-suppressed state

Key Takeaways

- Fluctuations in confined region mostly suppressed, while SOL scales become smaller
- E_r -well build-up, outer shear increases
- Transport and turbulence change locally, EM and NEO effects become important
- IOL dynamics observed
(further analysis required)

Next steps

- Implement heat and particle sources (partially completed)
- Repeat simulations with core heat source and separatrix particle source
- Diagnostics for IOL
(who does the work?)
- **Why is turbulence suppressed in the current case?**

Contact

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