



# ISOTOPE MIX EFFECT ON EDGE TURBULENT TRANSPORT IN PRE-L-H TRANSITION CONDITIONS IN JET-ILW

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- $\mathbf{E} \times \mathbf{B}$  shear introduction
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# CONTEXT – JET-ILW EXPERIMENTAL DATA

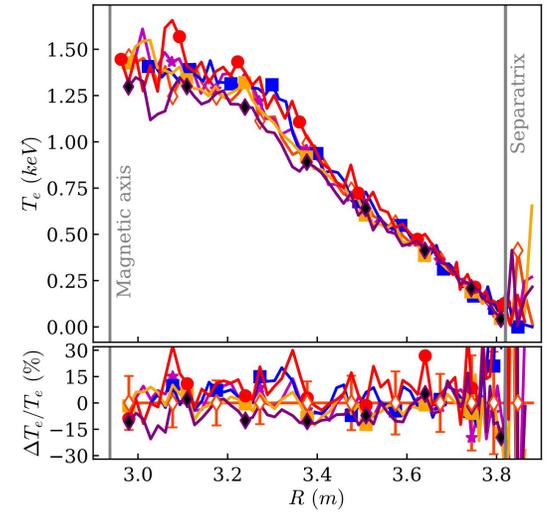
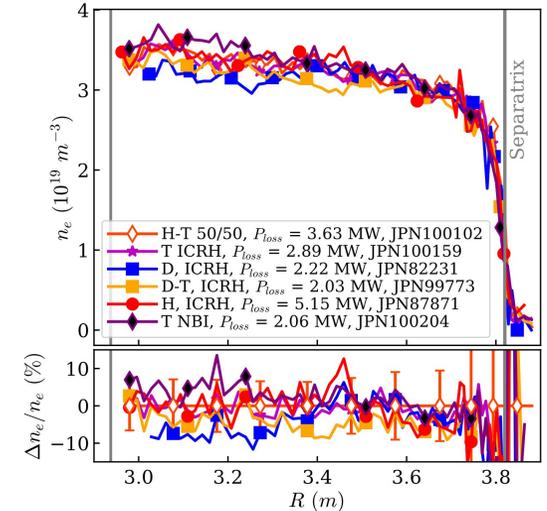
- L-mode pre-LH transition with **ICRH heating only**
- Low impurity content: **no ion/rotation measurements**
- 6 Pulses with similar profiles and magnetic configuration (Favourable x-point)
- $B = 1.8 \text{ T}$  |  $I_p = 1.65 \text{ MA}$  |  $q_{95} = 3.65$

$$P_{\text{loss}} - P_{\text{rad}} = P_{\text{sep}} \xrightarrow{\text{at } t_{\text{LH}}} P_{\text{LH}}$$

Pulse #	Species	$P_{\text{LH}}$ (MW)	$P_{\text{rad}}$ (MW)	$A_{\text{eff}}$
87871	H	4.55	0.75	1.01
82231	D	1.81	0.52	2
100102	53% H + 47% T	2.97	0.85	1.96
100102	26% H + 74% T	2.42	0.69	2.5
99773	49% D + 51% T	1.66	0.57	2.5
100159	T	1.68	1.22	2.92

**Isotope effect**

$$P_{\text{LH}} \propto 1/A_{\text{eff}}$$

$$P_{\text{LH}}(\text{H}) > P_{\text{LH}}(\text{D}) > P_{\text{LH}}(\text{T})$$


[Birkenmeier et al. *PPCF*, 2023]

# CONTEXT – ISOTOPE EFFECT

- **Isotope effect in pure plasma explained by passing electron dynamics**  
[Bonanomi et al. *PoP*, 2021]

- In JET-ILW is observed

$$P_{LH} \propto \chi_{eff}$$

- Transport is observed to be higher in lower isotope mass plasmas !
- Identical profiles but higher input power leads to **stronger heat fluxes (turbulent)**

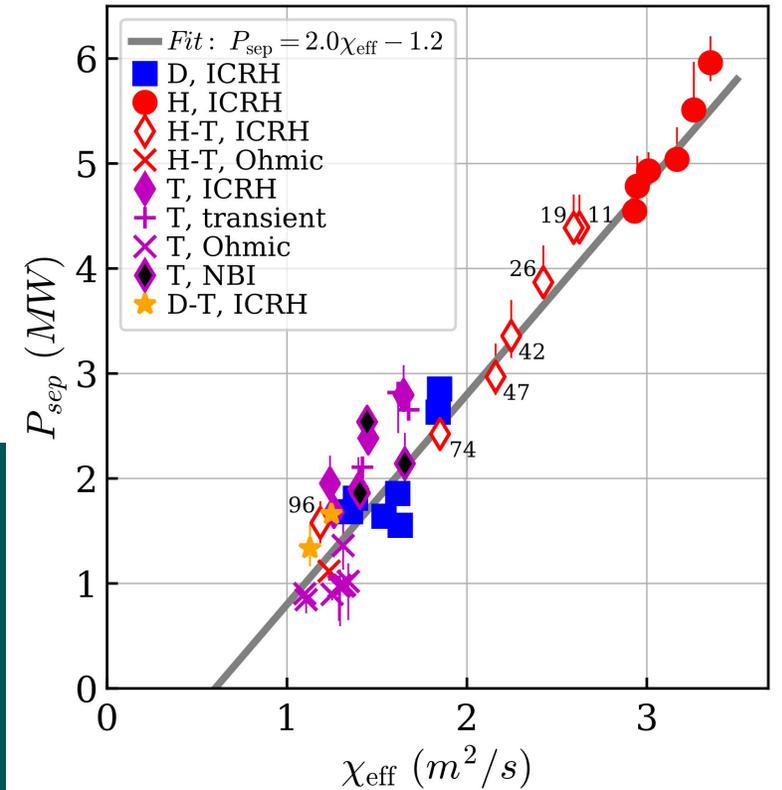
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**Isotope effect**

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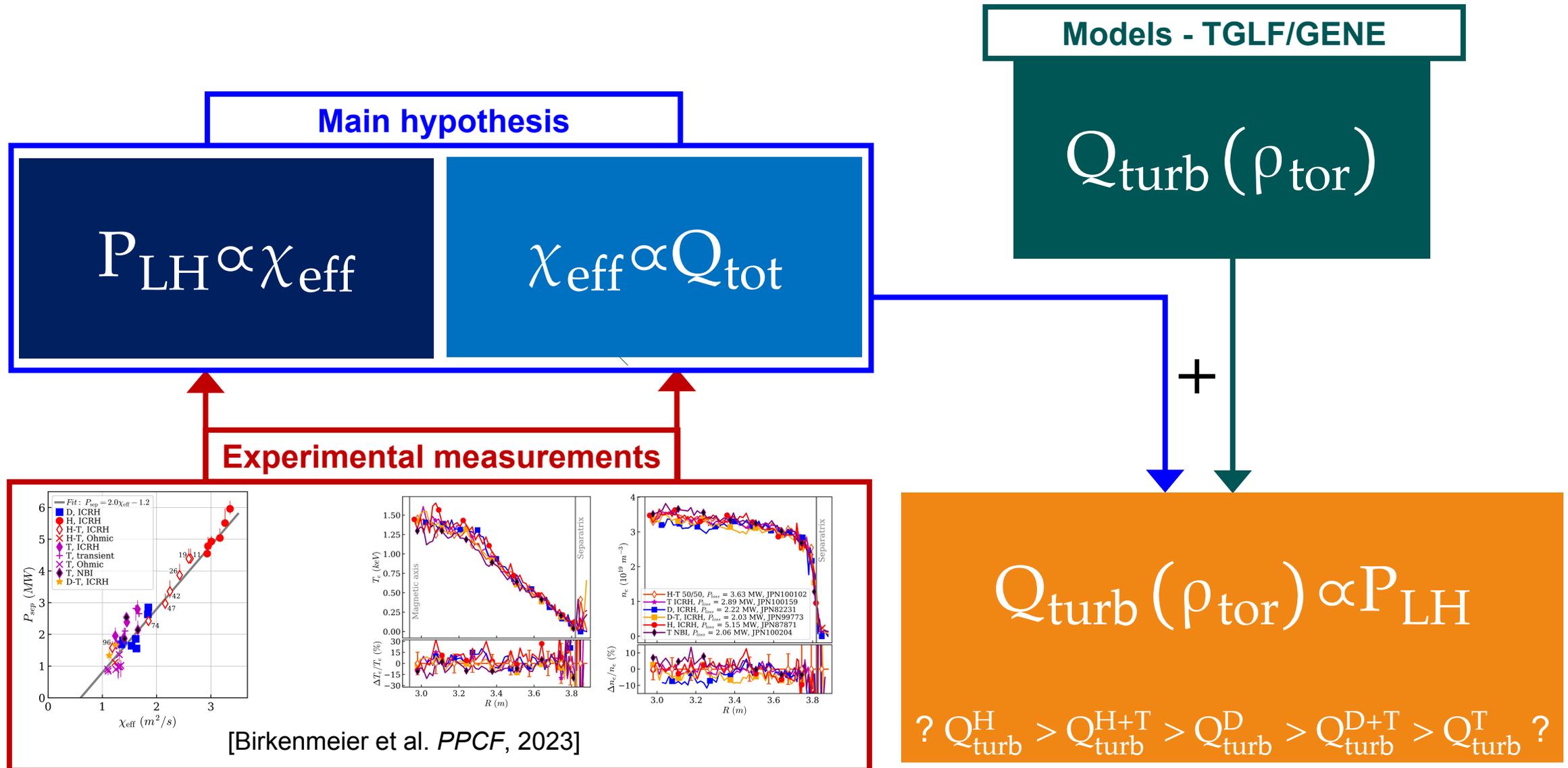
$P_{LH}(H) > P_{LH}(D) > P_{LH}(T)$

But at same  $A_{eff}$  ...

$$P_{LH}(H + T) > P_{LH}(D)$$


[Birkenmeier et al. *PPCF*, 2023]

# APPROACH OF THIS STUDY

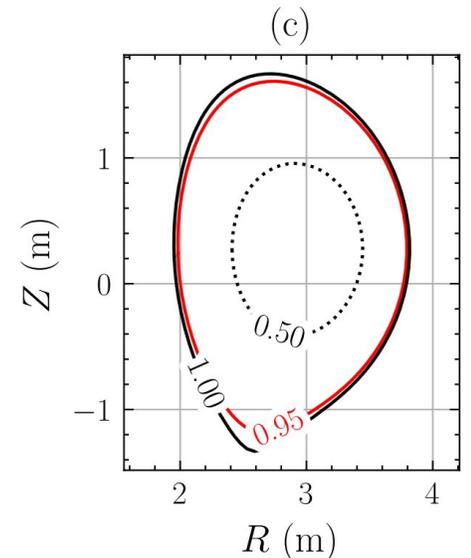
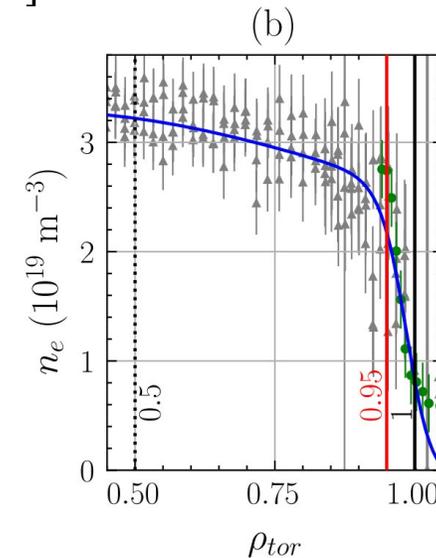
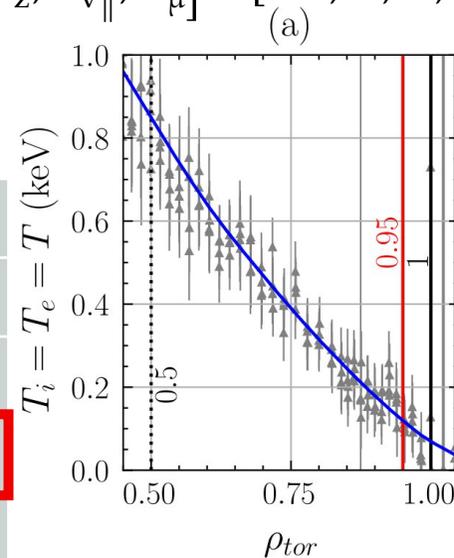


# SIMULATIONS SETUP

- **GENE flux-tube** used at  $\rho_{tor} = 0.95$
- Field-aligned coordinates  $[x, y, z, v_{\parallel}, \mu]$  with **Sugama collision operator w/ FLR effects**
- All simulations ran with the same reference values (see below) with  $\mathbf{m}_{ref} = \mathbf{m}_D = 2$
- Non-linear simulation resolution set as  $[N_x, N_y, N_z, N_{v_{\parallel}}, N_{\mu}] = [512, 72, 72, 48, 12]$
- $[L_x, L_y] = [248, 314] \rho_{c,D}$

(c) Flux surfaces taken from EQDSK file – equilibrium reconstruction

$v_e^*$	6.12	$R/L_n$	36.29
$\hat{v}_{ei}$	12.89	$R/L_T$	40.6
$\beta_e$	$3.29 \times 10^{-4}$	$n_e$	$2.07 \times 10^{19} \text{ m}^{-3}$
$s/\hat{q}$	3.63/3.65	$T_i \equiv T_e = T$	<b>0.119 keV</b>
$T_i/T_e$	1	$B_{tor}$	1.91 T
$\rho_{c,D}$	0.827 mm	$\rho^* \equiv \rho_{c,D}/L_T$	$\sim 1/87$

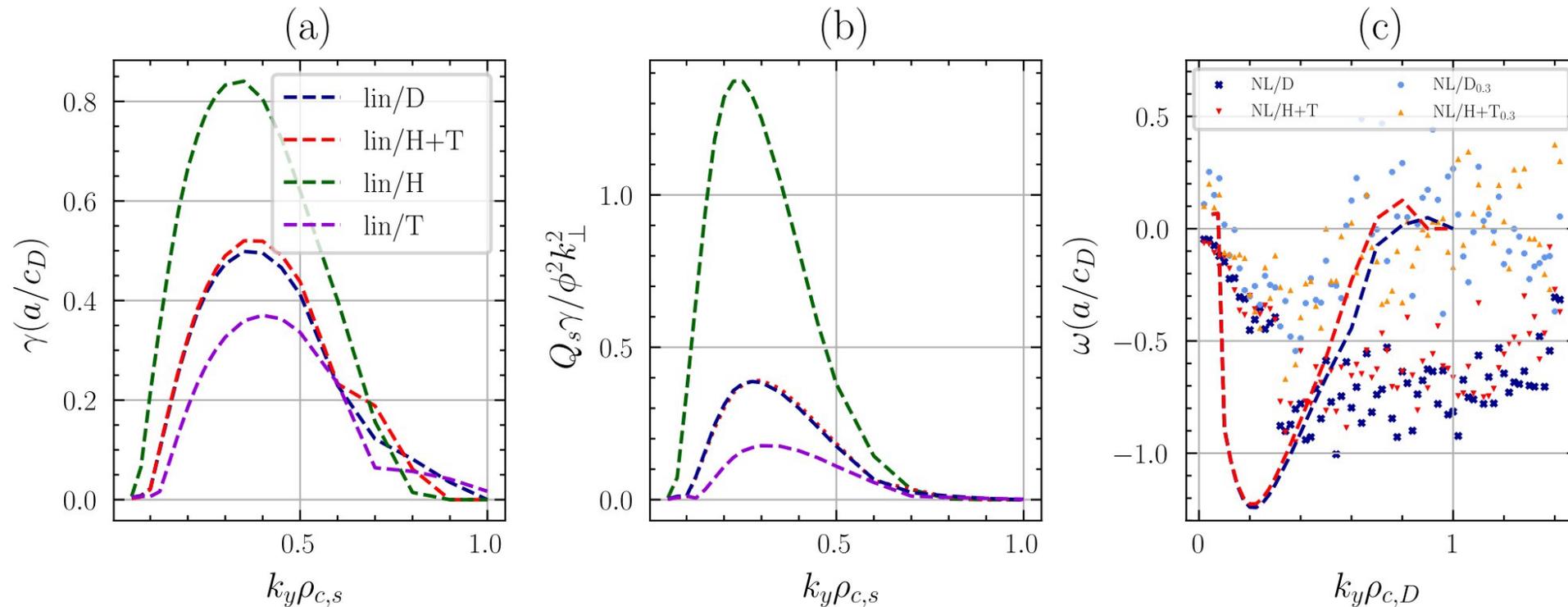


▲ : HRTS data  
 ● : LiBE data  
 — : Custom fit used as input for GENE/TGLF

# TURBULENCE NATURE – LINEAR SCANS

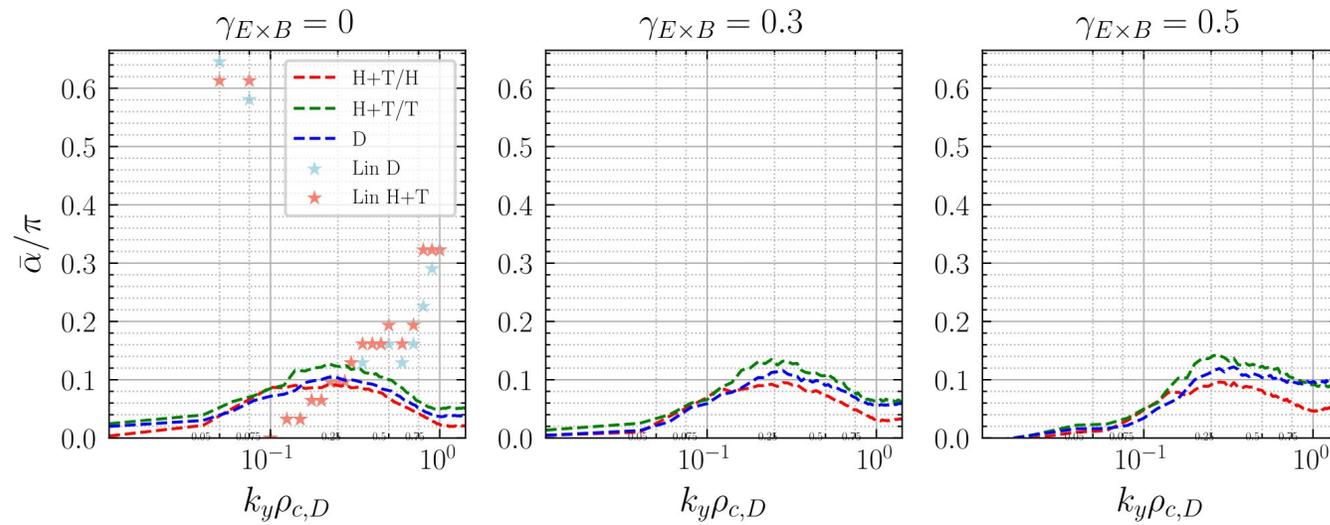
- Anti-gyroBohm scaling of  $\gamma$  and QL heat fluxes
- Negative frequencies, similar between linear and NL simulations

**No differences between H+T and pure deuterium !!**



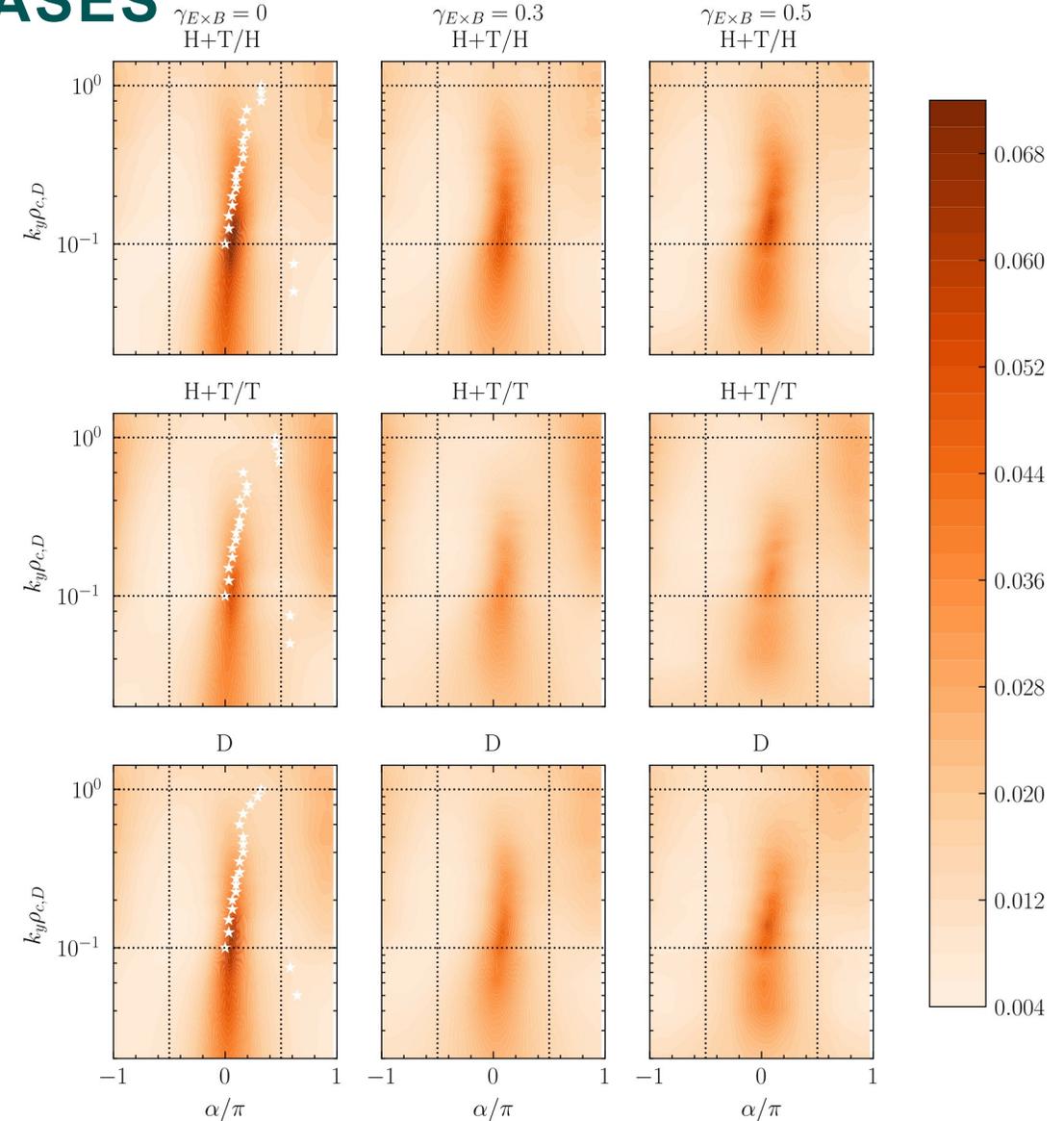
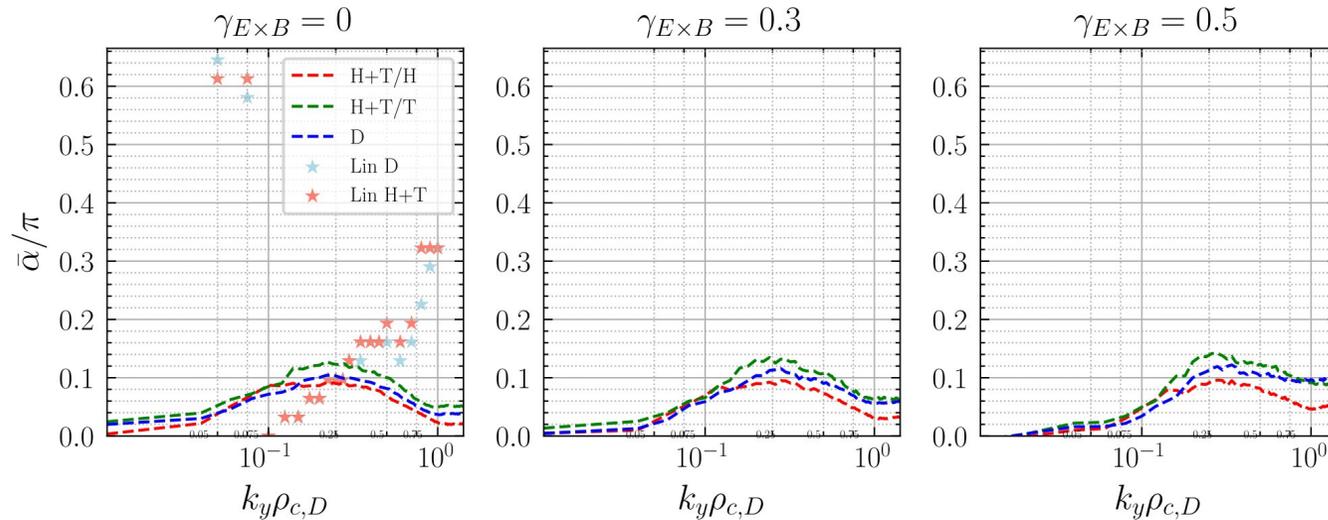
# TURBULENCE NATURE – CROSS-PHASES

- Linear and non-linear simulations cross-phases :  $\alpha(\phi \times T_{\perp}) \rightarrow 0$



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- Linear and non-linear simulations cross-phases :  $\alpha(\phi \times T_{\perp}) \rightarrow 0$
- $\alpha(\phi \times T_{\perp}) \rightarrow 0$  on all  $k_y \rho_{c,D}$  range, broadening with  $\gamma_{E \times B}$

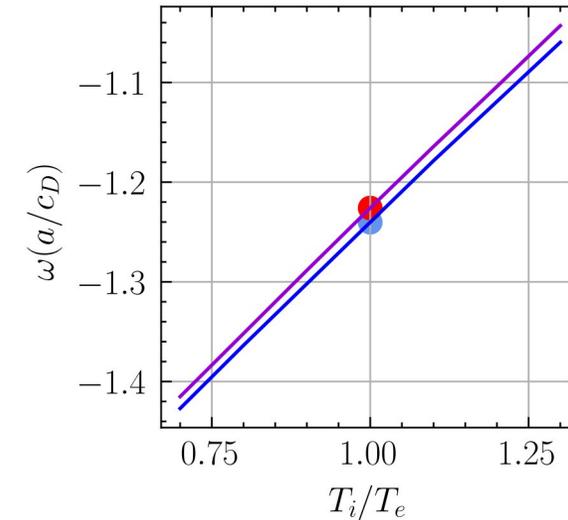
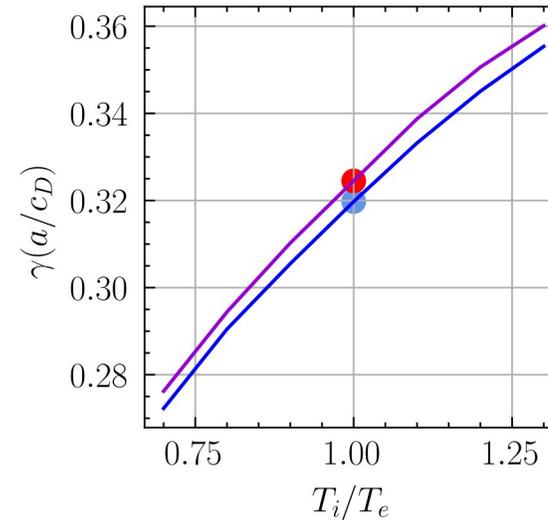
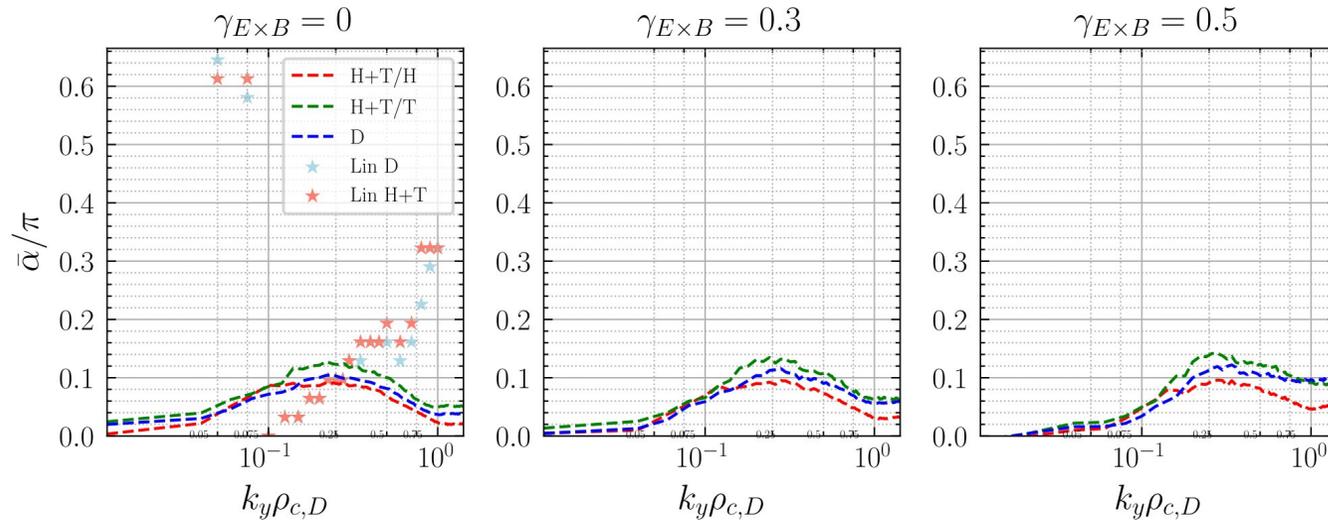
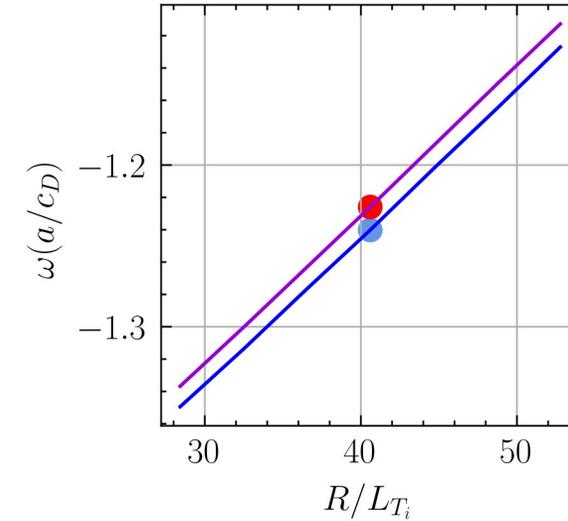
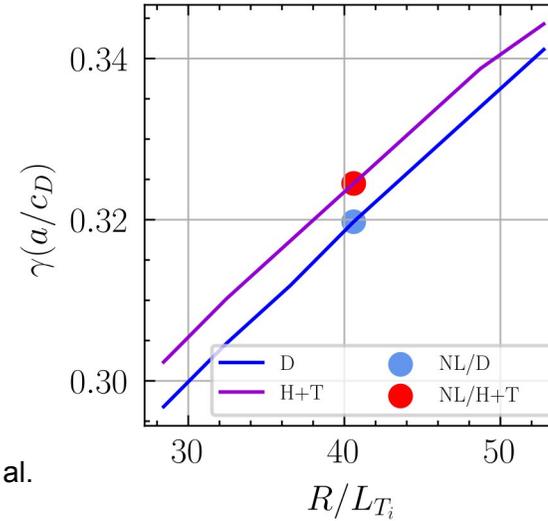


# TURBULENCE NATURE – $T_i$ SENSITIVITY

- Linear and non-linear simulations cross-phases :  $\alpha(\phi \times T_\perp) \rightarrow 0$
- $\alpha(\phi \times T_\perp) \rightarrow 0$  on all  $k_y \rho_{c,D}$  range, broadening with  $\gamma_{E \times B}$
- No major scaling with  $T_i$  profile

## Resistive electron drift-wave modes

[Bonanomi et al. NF, 2019], [Bonanomi et al. PoP, 2021], [Bonanomi et al. PoP, 2024], [Bonanomi et al. PPCF, 2025]

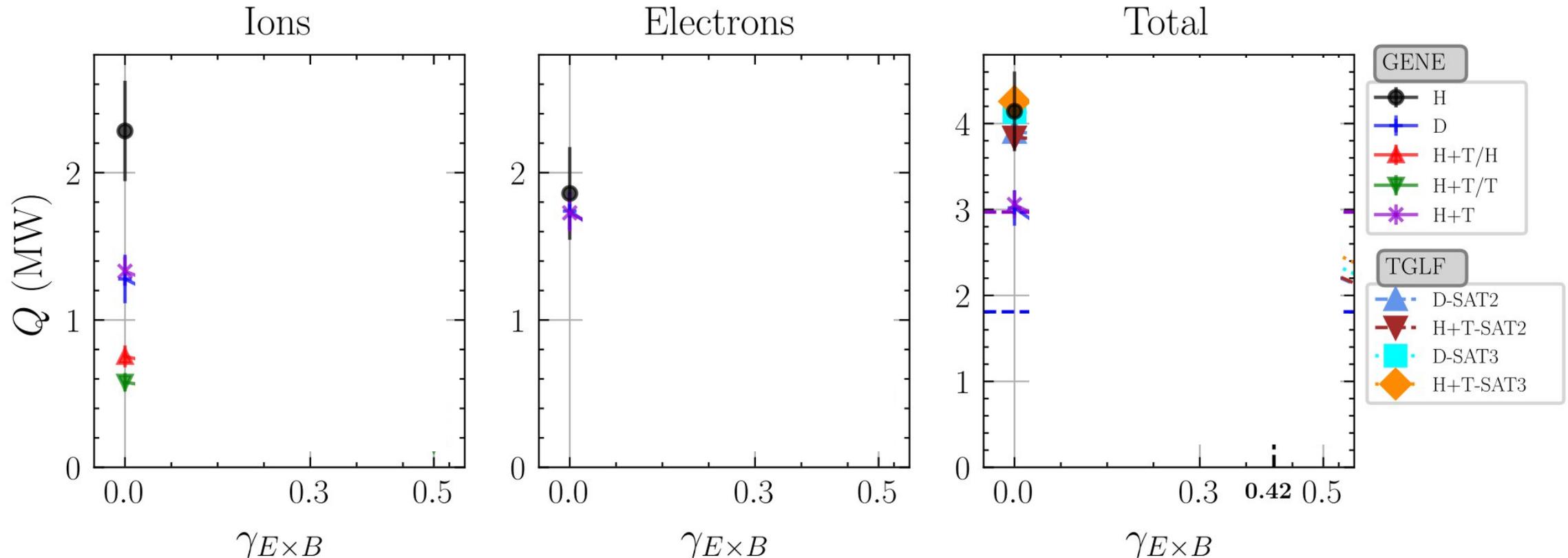


# NON-LINEAR RESULTS – HEAT FLUXES – PART 1

- Anti-gyroBohm scaling retrieved for non-linear heat fluxes

$$Q_H^{\gamma_{E \times B}=0} > Q_D^{\gamma_{E \times B}=0} = Q_{H+T}^{\gamma_{E \times B}=0}$$

Good agreement with experimental flux levels BUT no differences at same effective mass between H+T and pure D...



# NON-LINEAR RESULTS – $E \times B$ SHEAR

- Anti-gyroBohm scaling retrieved for non-linear heat fluxes

$$Q_H^{\gamma_{E \times B} = 0} > Q_D^{\gamma_{E \times B} = 0} = Q_{H+T}^{\gamma_{E \times B} = 0}$$

Good agreement with experimental flux levels BUT no differences at same effective mass between H+T and pure D...

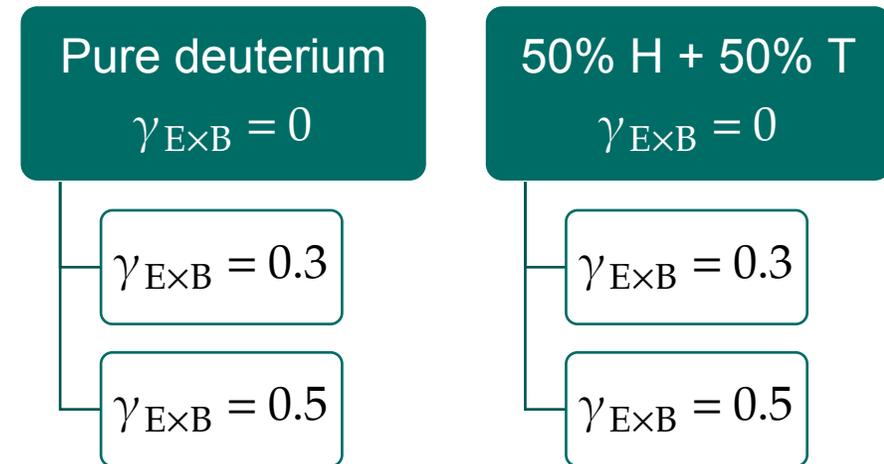
What about  $\gamma_{E \times B}$  effect ? -> not measured, but **expected to be non-zero !**

$$\gamma_{E \times B} = \frac{\rho_{\text{tor}}}{q} \frac{\partial}{\partial \rho_{\text{tor}}} \left( \frac{E_r}{B_p R} \right) \frac{\rho_a}{c_D}$$

- **Estimated** using  $T_i = T_e$  and the **diamagnetic component only**

$$\gamma_{E \times B}^{\text{est}} \approx 0.42$$

- Fork the simulations by adding  $\gamma_{E \times B}$  after saturation is reached

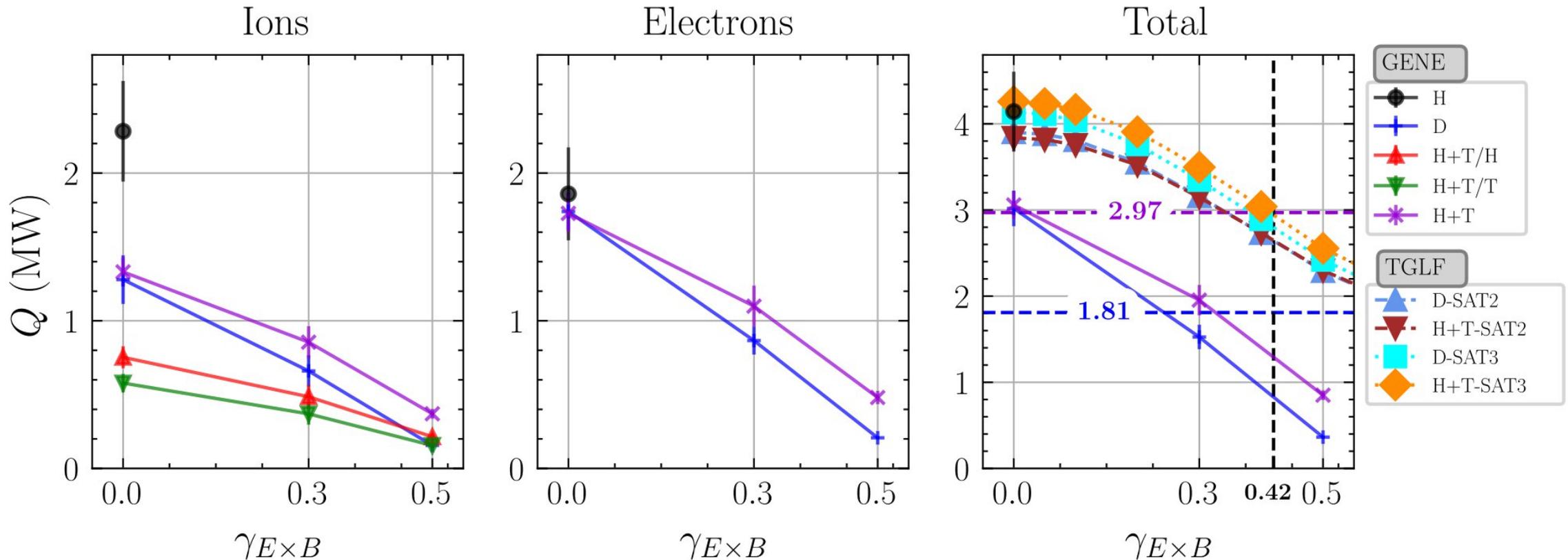


# NON-LINEAR RESULTS – HEAT FLUXES – PART 2

- Non-zero  $\gamma_{E \times B}$  allows to retrieve a behaviour similar to experiments

$$Q_D^{\gamma_{E \times B} \neq 0} < Q_{H+T}^{\gamma_{E \times B} \neq 0}$$

- This behaviour is **not** retrieved with TGLF SAT2 and 3



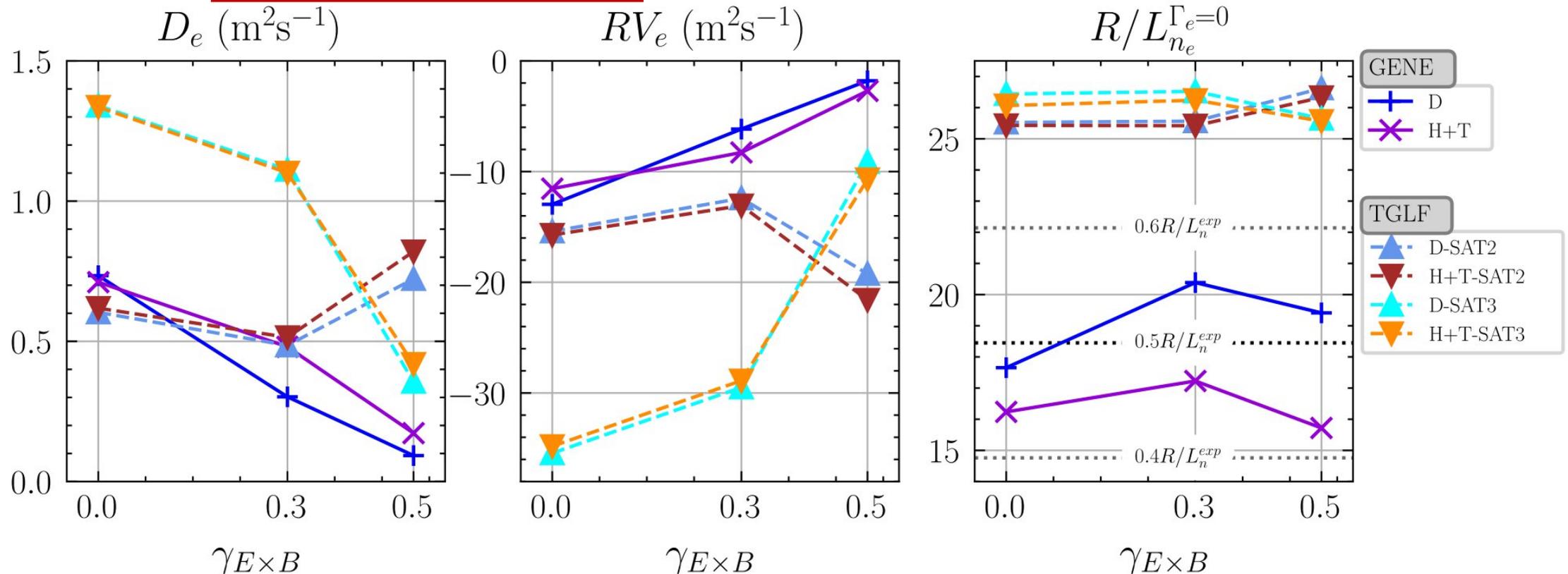
# NON-LINEAR RESULTS – PARTICLE TRANSPORT

- Two electron population; one with  $R/L_n = 0$ , one with  $R/L_n = 2 \times R/L_{ni}$
- Allows to extract diffusion and convection, as well as zero flux gradients

$$\Gamma_e = -D_e \nabla n_e + n_e V_e$$

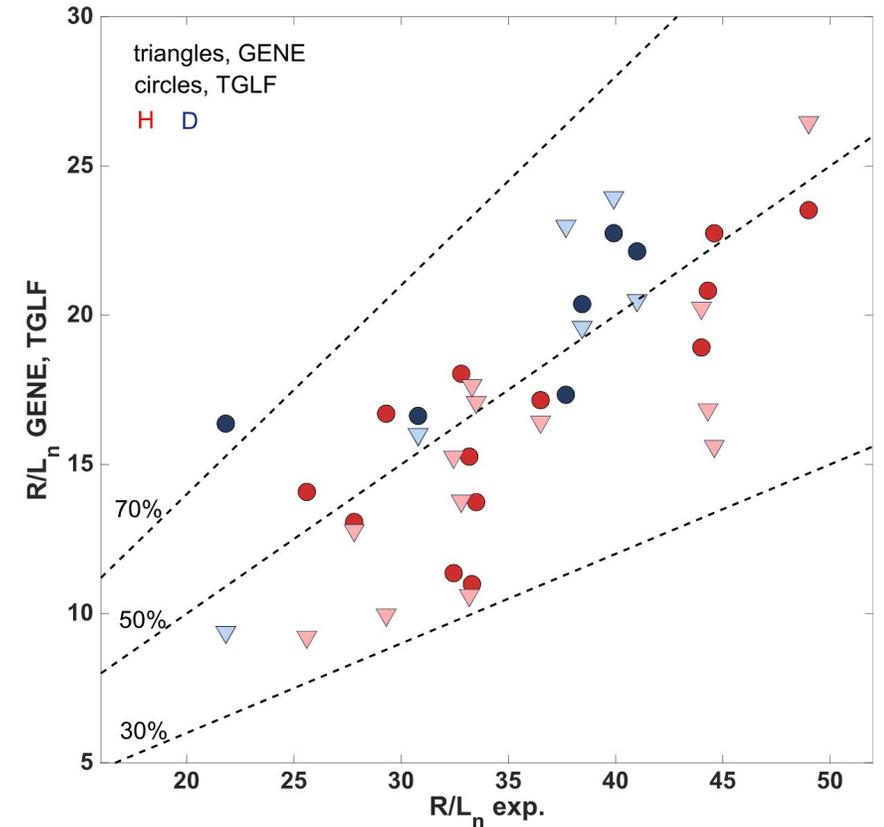
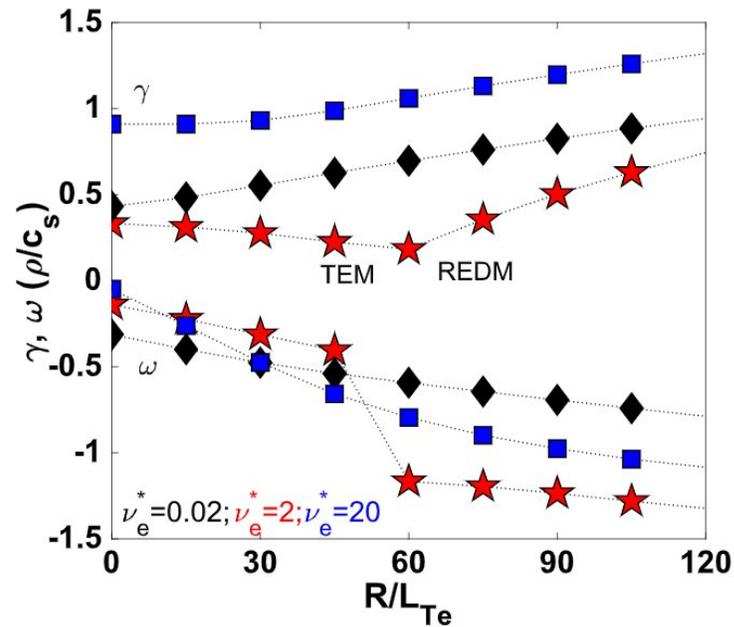
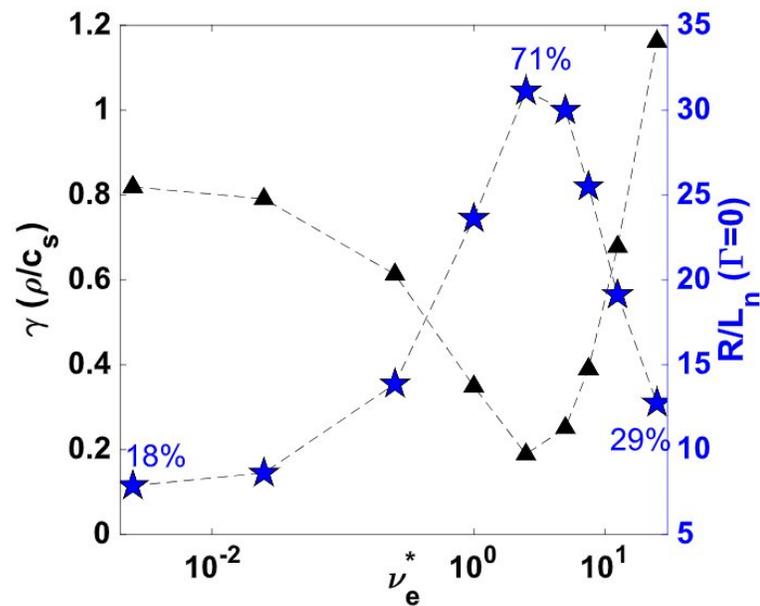
Overestimated convection for TGLF SAT2/3 compared to GENE

Turbulence represents here  $\sim$  half the density peaking ! (rest are sources)



# NON-LINEAR RESULTS – PARTICLE TRANSPORT

- GENE and TGLF with L-mode ASDEX Upgrade experimental data also show **half the density peaking is coming from turbulence**
- $R/L_{ni} = 36.29$  in our case, similar conditions as previous study from Nicola Bonanomi
- Comparable turbulent modes, scales with  $R/L_{Te}$



[Bonanomi et al. *PPCF*, 2025]

## CONCLUSION AND FUTURE WORK

❑ Linear simulations: **electron drift wave modes** with decreasing  $\gamma$  for increasing  $A_{\text{eff}}$ , consistently with  $P_{\text{LH}}$

$\gamma_{\text{E} \times \text{B}} = 0$  : linear, non-linear GENE and TGLF simulations yield no differences at  $A_{\text{eff}} = 2$

$\gamma_{\text{E} \times \text{B}} > 0$  : only non-linear GENE simulations provide  $Q_{\text{D}} < Q_{\text{H}+\text{T}}$ , similarly to  $P_{\text{LH}}$  experimental scaling ( $P_{\text{LH}}^{\text{D}} < P_{\text{LH}}^{\text{H}+\text{T}}$ )

❑ Separation of both cases due to **asymmetrical response of turbulence levels**, stronger for D than H+T

❑ **Differential effect on fluxes** (heat & particle) increases with  $\gamma_{\text{E} \times \text{B}}$  magnitude

❑ Heat fluxes are in the experimental range (of  $P_{\text{LH}}$ ) at  $\sim 3\text{MW}$

❑ **No such effect is observed in TGLF SAT2 and SAT3**

Next steps:

- Explore similar simulations at  $A_{\text{eff}} = 2.5$  for H(25%)+T(75%), D(50%)+T(50%) and « fake » 2.5 pure isotope
- Explore more in depth turbulence nature (look at ZF for instance)
- Check the latest version of TGLF and try integrated modelling of these pulses

# ISOTOPE EFFECT – PARALLEL ELECTRON DYNAMICS

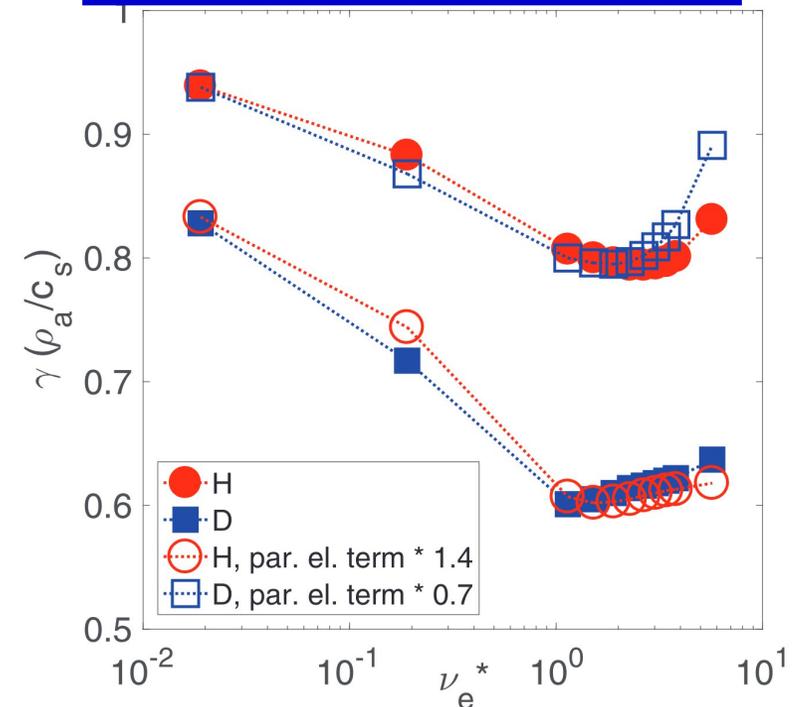
What is driving this isotope effect ?

$$C_m \left\{ -\hat{v}_{e,th} \frac{\hat{C}}{JB_0} \left[ v_{\parallel} \left( \frac{\partial f_{e,1}}{\partial z} - \frac{e}{T_e} F_{e,0} \frac{\partial \phi}{\partial z} + F_{e,0} \mu \frac{\partial B_{1,\parallel}}{\partial z} \right) - \mu \frac{\partial}{\partial z} \left( B_0 \frac{\partial f_{e,1}}{\partial v_{\parallel}} \right) \right] \right\}$$

$$C = \left( 0,51 v_{ei} \frac{L_{\perp}}{c_s} \frac{m_e}{m_i} \left( \frac{qR}{L_{\perp}} \right)^2 \right)$$

- Parallel electron dynamics term  $\propto 1/m_i$  (i.e. can be recast as  $\propto 1/A_{eff}$ )
- Increased coupling between electrons and ions for lower ion masses
- Artificially changing  $C_m$  allowed match turbulent states between H and D

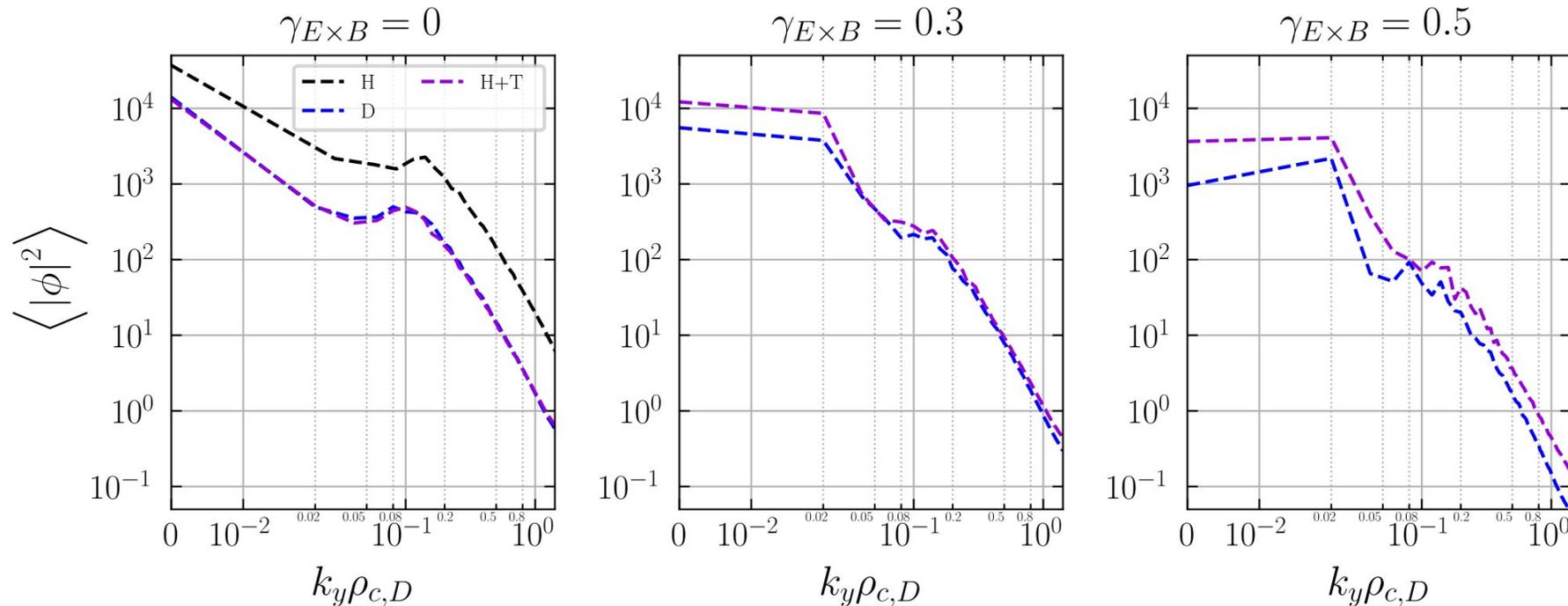
What about this term when we have two species ?



[Bonanomi et al. *PoP*, 2021]

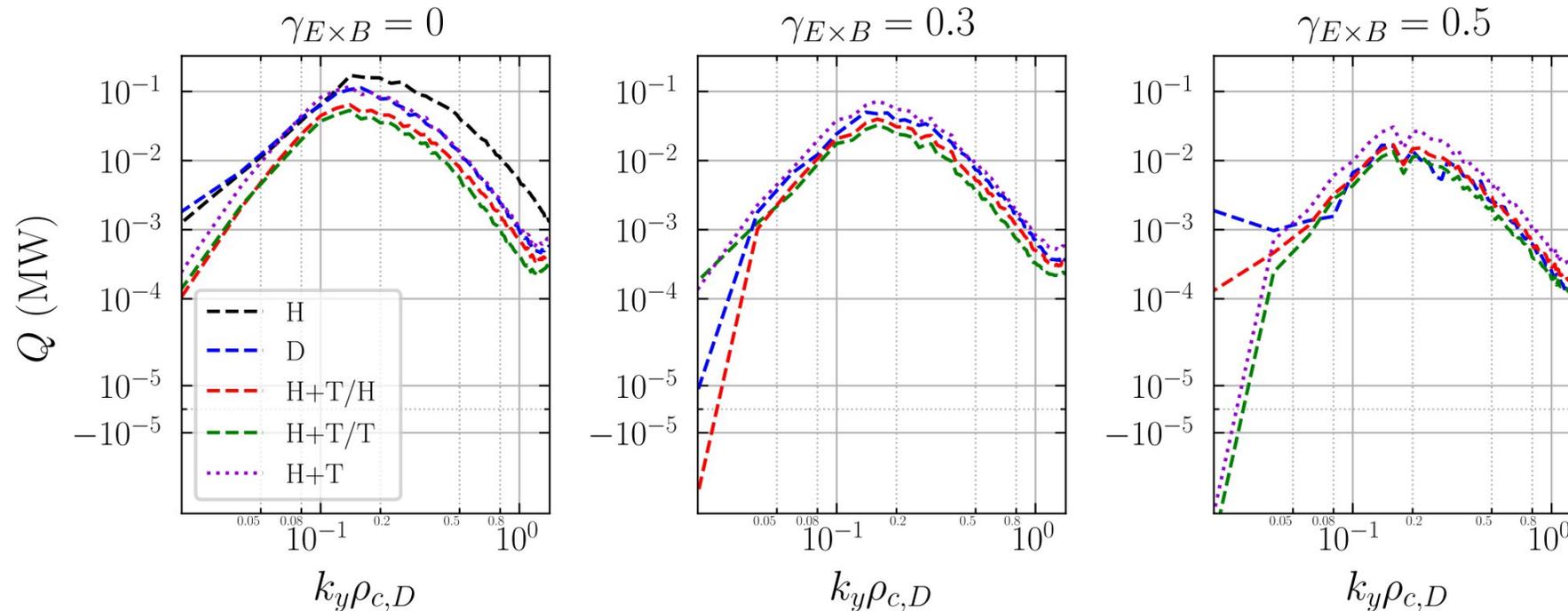
# NON-LINEAR RESULTS – POTENTIAL SPECTRA

- Identical spectra / turbulence between cases, H case has stronger modes
- Reduction of amplitude for bigger scales
- Generalized to smaller scales when increasing  $E \times B$  shear
- Suspicion of **different ZF activities**



# NON-LINEAR RESULTS – IONS HEAT FLUX SPECTRA

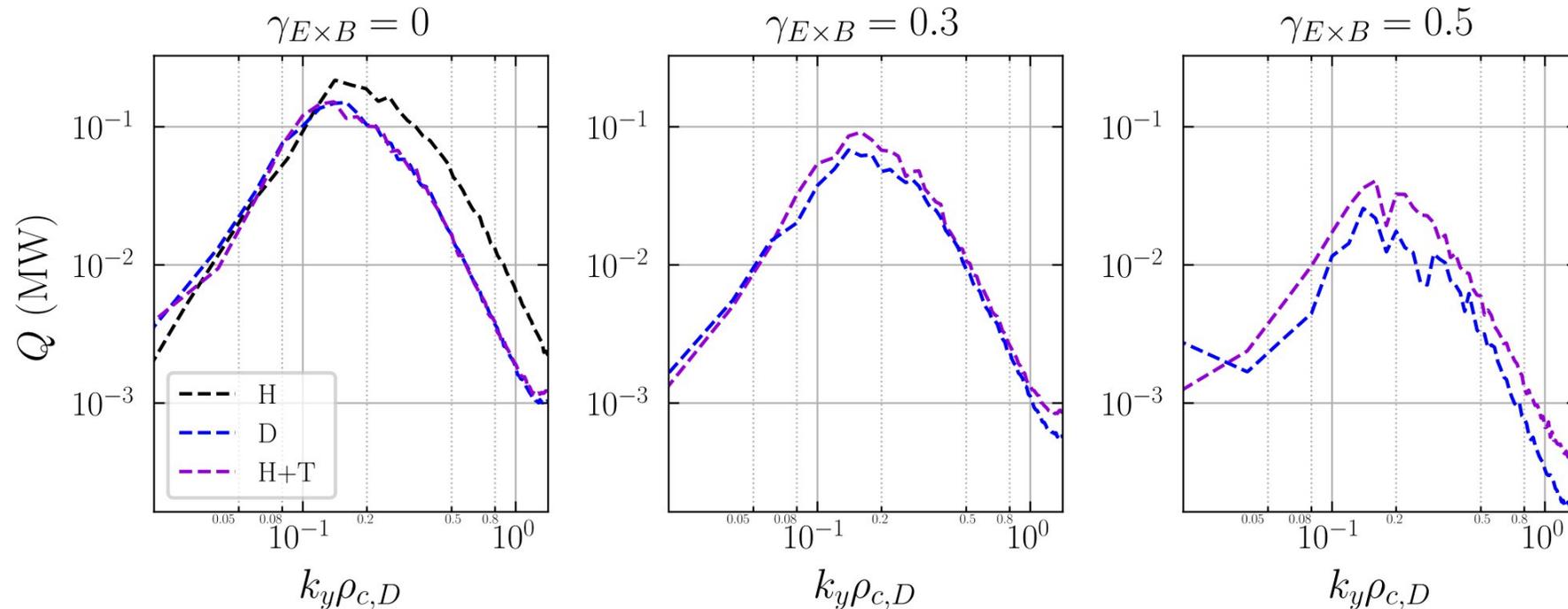
- Between pure H and H+T/D, smaller scales makes the difference for the ion heat flux
- For larger scales, difference is larger for H+T and D, but magnitudes are low
- The main difference is, as expected, in the range  $0.07 < k_y \rho_{c,D} < 0.4$  for  $\gamma_{E \times B} = 0.3$
- This is extended to  $0.06 < k_y \rho_{c,D} < 1.2$  for  $\gamma_{E \times B} = 0.5$



# NON-LINEAR RESULTS – ELECTRONS HEAT FLUX SPECTRA

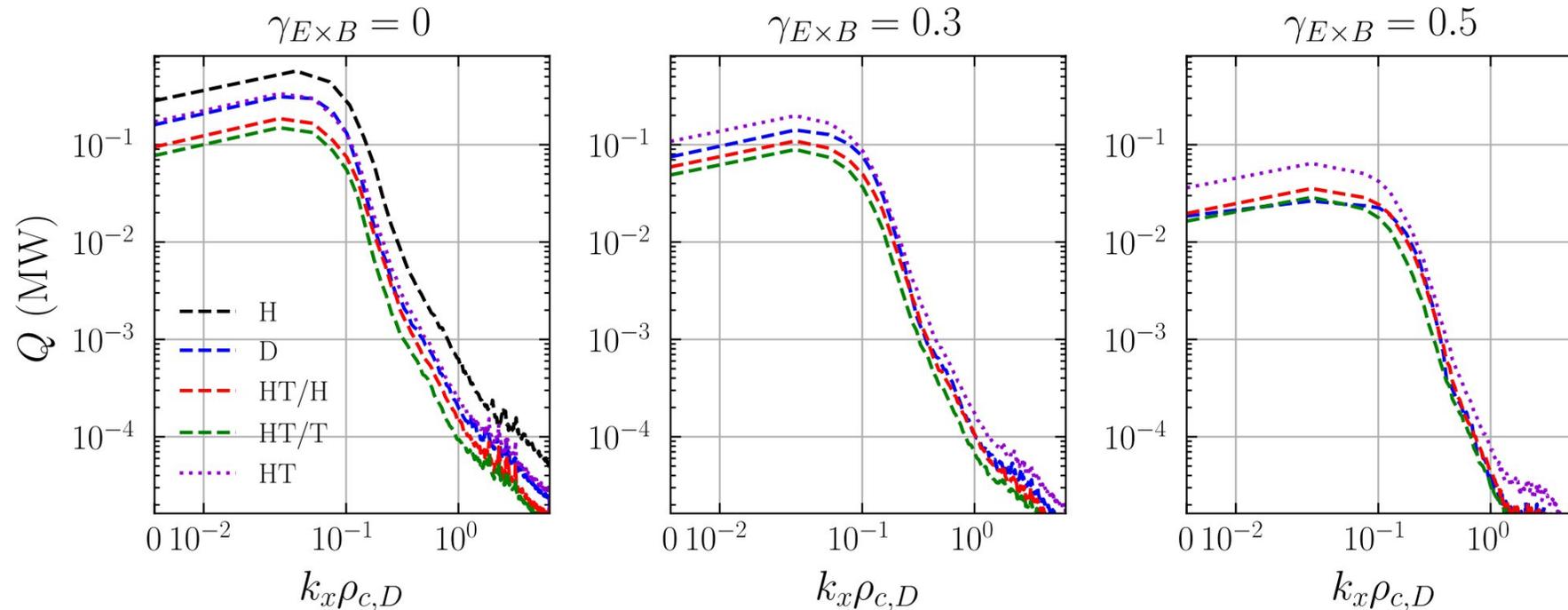
Similarly as before:

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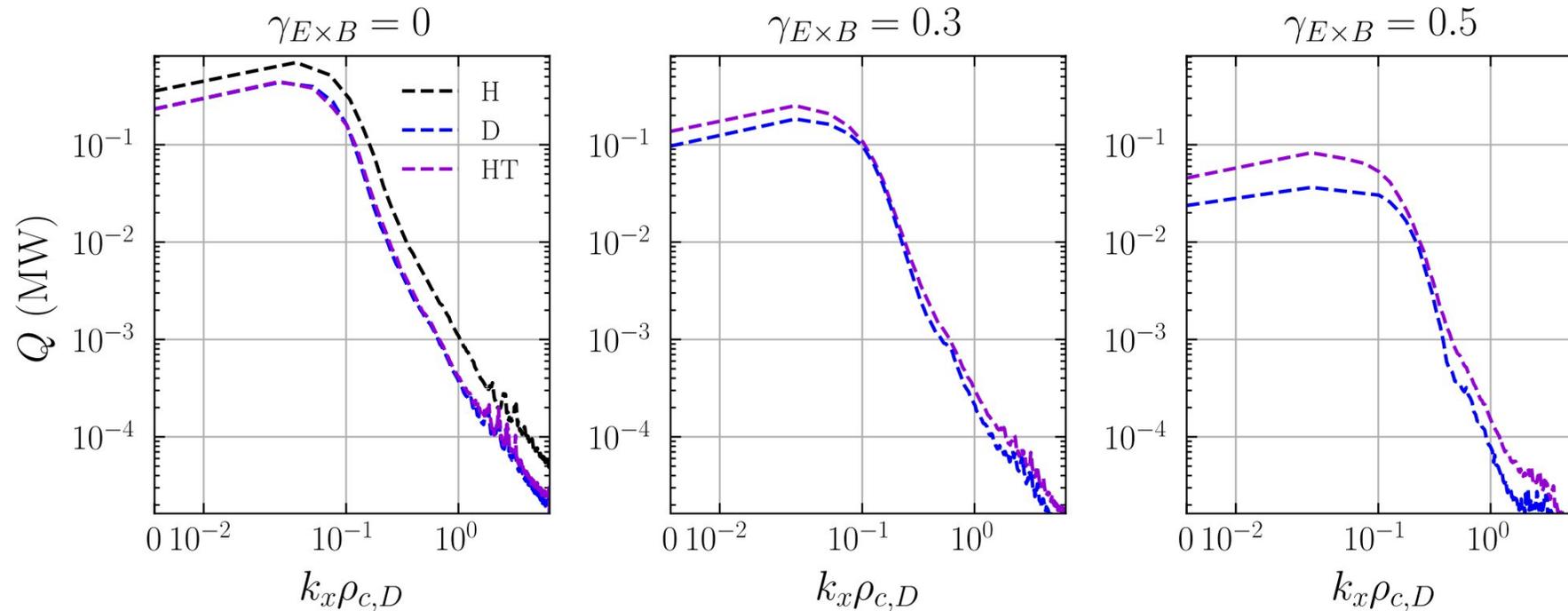
- Difference between H and H+T/D is present on all scales
- Again, larger scales is where the difference is larger for H+T and D
- The main difference is in the range  $0 < k_x \rho_{c,D} < 0.1$  for  $\gamma_{E \times B} = 0.3$
- This is extended to  $0 < k_x \rho_{c,D} < 0.2$  for  $\gamma_{E \times B} = 0.5$



# NON-LINEAR RESULTS – ELECTRONS HEAT FLUX SPECTRA

Similarly as before:

- The main difference is in the range  $0 < k_x \rho_{c,D} < 0.1$  for  $\gamma_{E \times B} = 0.3$
- This is extended to  $0 < k_x \rho_{c,D} < 0.2$  for  $\gamma_{E \times B} = 0.5$



# NON-LINEAR RESULTS – PARTICLE FLUX

- Expected fluxes are also in a “reasonable” range
- Also more H flux than T in the H+T mixtures
- Could the additional particle transport lead to higher convective heat flux ? (need to check)
- Similar trend as heat flux though (reason why not shown)
- Again, TGLF is not “that bad”

