



fs-TALIF in RAID ENR project report 2025

marcelo.baquero@epfl.ch

**Marcelo Baquero-Ruiz, Y. Andrebe, F. Carbone, A. Clement, I. Furno,
M. Goddijn, P. Lavanchy, M. Puppini and S.P.H. Vincent**

24 November 2025

1. Introduction
2. Implementation
3. Experiments
4. Other studies
5. Scientific output
6. Conclusions and outlook

1. Introduction
2. Implementation
3. Experiments
4. Other studies
5. Scientific output
6. Conclusions and outlook

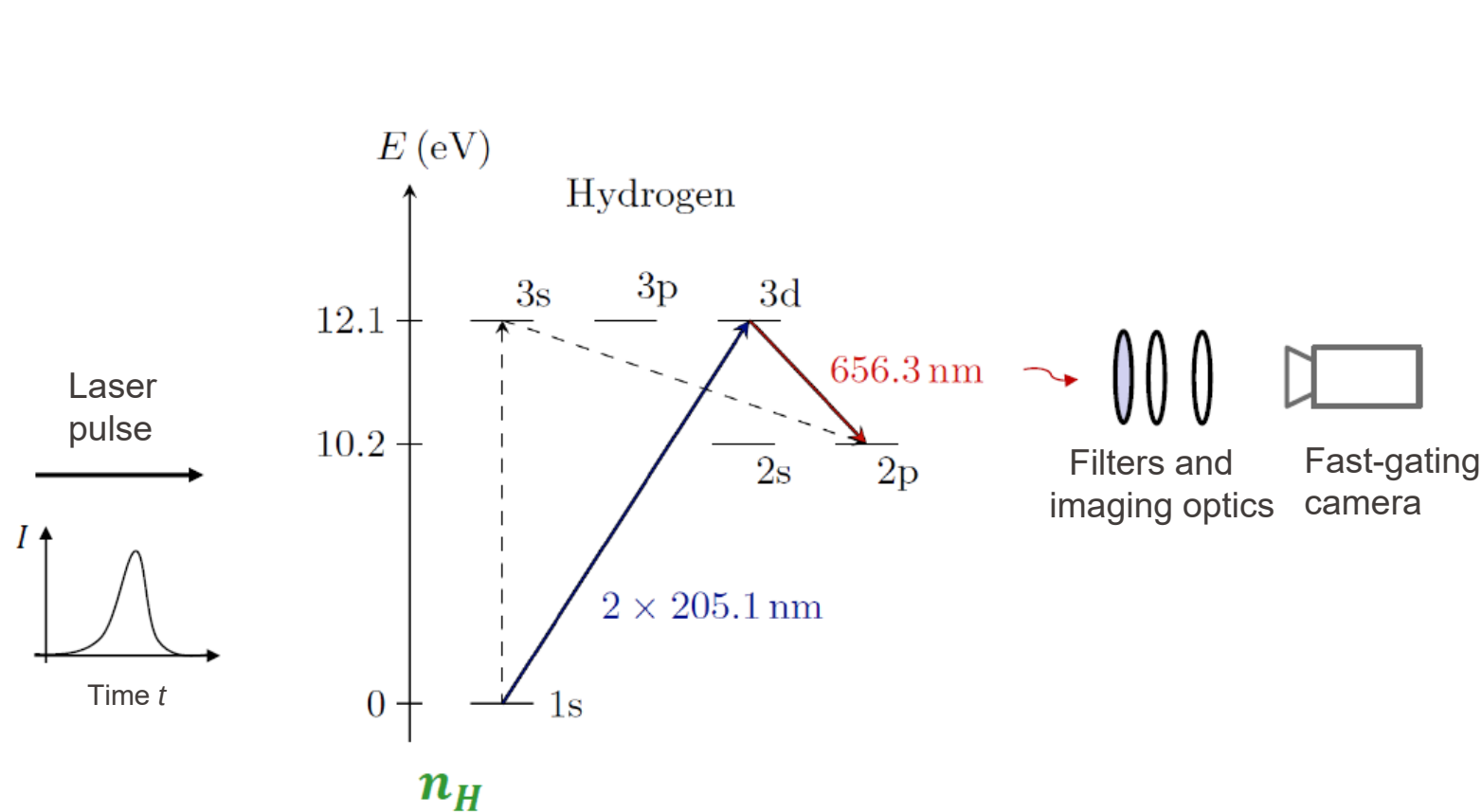
Introduction

Our objective is

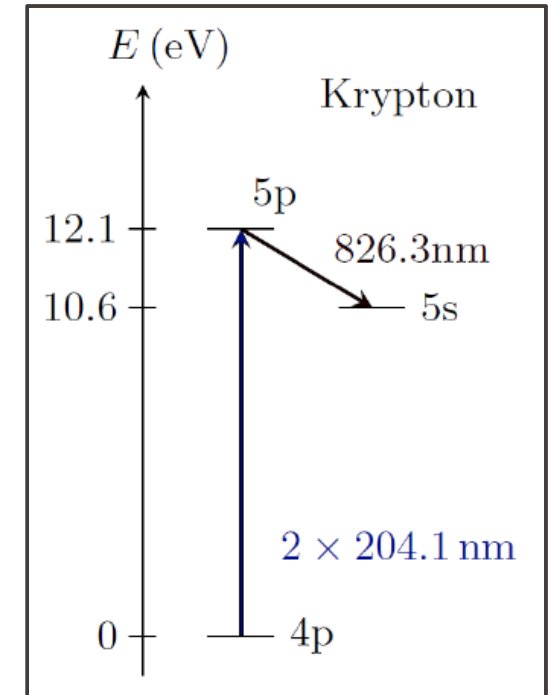
*Obtain **single laser-pulse** measurements of H densities in a dense hydrogen plasma using fs-TALIF.*

Provide an assessment of the feasibility of this method for application in magnetic confinement fusion research.

We use the Two-photon absorption laser-induced fluorescence (TALIF) technique



TALIF in Kr
(for calibration)



We use the Two-photon absorption laser-induced fluorescence (TALIF) technique

Total fluorescence photons
from observation volume
in one laser pulse

Fluorescence
branching ratio

Photon
stat. factor

Two-photon absorption
cross section

$$S_H(\nu_L) = n_H \beta_H a_H g_H(\nu_L) \frac{G \sigma_H^{(2)}}{(h\nu_L)^2} \int_{\text{obs.} V} \int_{\text{pulse}} I^2(\mathbf{r}', t') dt' d\mathbf{r}'$$

H density

Detection efficiency
(includes geometry,
transmission, etc.)

Normalized
absorption line
profile

Laser intensity
(squared)

We use the Two-photon absorption laser-induced fluorescence (TALIF) technique

Diagram illustrating the TALIF technique equation and its components:

$$S_F(\nu_F) = \eta_D \beta_{if} \sigma_H^{(2)} \frac{G}{(h\nu_L)^2} \int_{\text{obs. } V} \int_{\text{pulse}} I^2(\mathbf{r}', t') dt' d\mathbf{r}'$$

Labels and arrows pointing to the equation components:

- Total fluorescence photons from observation volume in one laser pulse (points to $S_F(\nu_F)$)
- Fluorescence branching ratio (points to β_{if})
- Photon stat. factor (points to G)
- Two-photon absorption cross section (points to $\sigma_H^{(2)}$)
- H density (points to $\sigma_H^{(2)}$)
- Detection efficiency (includes geometry, transmission, etc.) (points to η_D)
- Normalized absorption line profile (points to G)
- Laser intensity (squared) (points to $I^2(\mathbf{r}', t')$)

Valid only for "sufficiently" low laser intensities

TALIF in Kr gas is used to calibrate the H TALIF signal

Diagram illustrating the calibration of the H TALIF signal using Kr gas. The equation relates the H density (n_H) to the Kr density (n_{Kr}) through several ratios:

$$\frac{n_H}{n_{Kr}} = \frac{\beta_{Kr}}{\beta_H} \left(\frac{\sigma_{Kr}^{(2)}}{\sigma_H^{(2)}} \right) \left(\frac{a_{Kr}}{a_H} \right) \left(\frac{g_{Kr}(\nu_L)}{g_H(\nu_L)} \right) \left(\frac{E_{\text{pulse},Kr}^2}{E_{\text{pulse},H}^2} \right) \left(\frac{S_H(\nu_L)}{S_{Kr}(\nu_L)} \right)$$

Annotations and arrows indicate the physical quantities represented by the terms in the equation:

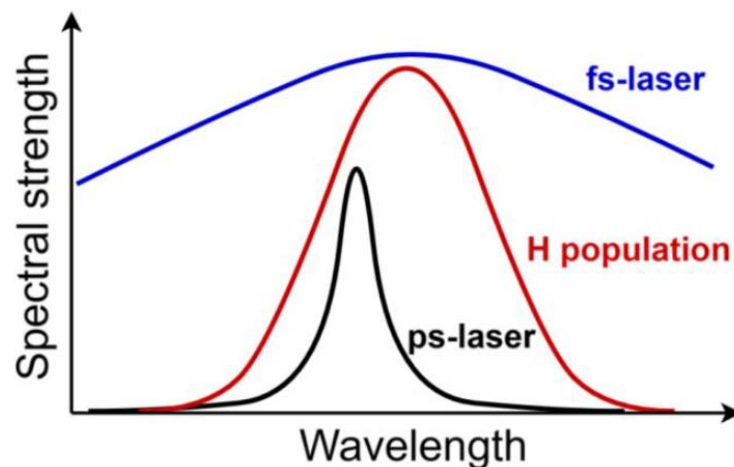
- H density:** Points to n_H (green arrow).
- Ratio of detection efficiencies:** Points to $\frac{\beta_{Kr}}{\beta_H}$ (black arrow).
- “Known” ratio of absorption cross sections:** Points to $\left(\frac{\sigma_{Kr}^{(2)}}{\sigma_H^{(2)}} \right)$ (black arrow).
- Branching ratio of H fluorescence:** Points to $\frac{a_{Kr}}{a_H}$ (red arrow).
- Ratio of relative absorptions at ν_L :** Points to $\left(\frac{g_{Kr}(\nu_L)}{g_H(\nu_L)} \right)$ (blue arrow).
- Ratio of pulse energies (squared) used in H and Kr measurements:** Points to $\left(\frac{E_{\text{pulse},Kr}^2}{E_{\text{pulse},H}^2} \right)$ (black arrow).
- Ratio of TALIF signals in H and in Kr:** Points to $\left(\frac{S_H(\nu_L)}{S_{Kr}(\nu_L)} \right)$ (black arrow).

We investigate TALIF with fs pulses

- fs pulses are spectrally very broad and significantly larger than the absorption linewidth of H in a **tokamak divertor or SOL** ($T_e \leq 10$ eV, $B \approx 1$ T).



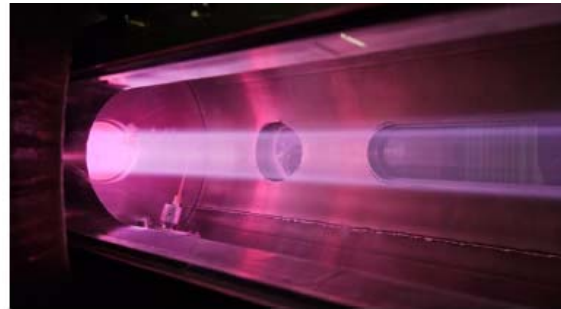
They can excite the entire H population with a single laser pulse.



- Furthermore, fs pulses can be extremely intense at similar pulse energies as ps lasers.
- Since fs pulses are spectrally very broad, a Kr TALIF signal can in principle be obtained with ν_L tuned to the H TALIF resonance.

Experiments are carried out in the Resonant Antenna Ion Device (RAID) at SPC

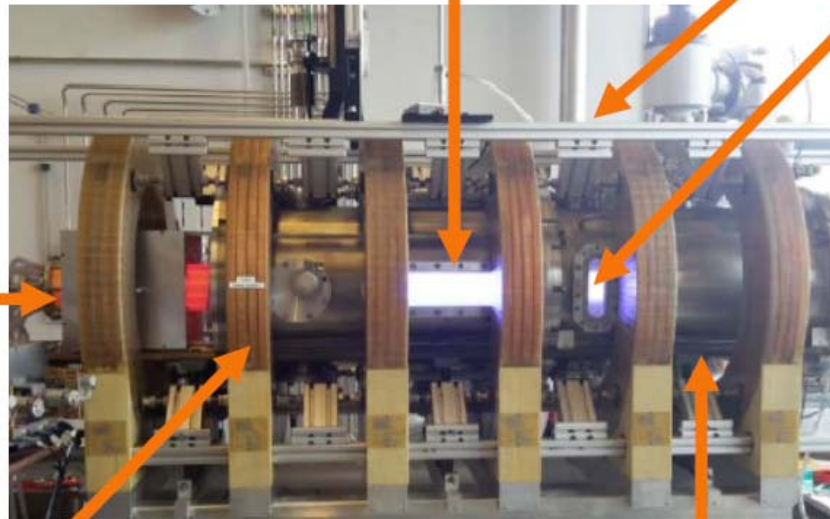
$$\begin{aligned} B_0 &= 100 - 800 \text{ G} \\ p_{\text{H}_2} &= 0.1 - 2.0 \text{ Pa} \\ P_{\text{RF}} &\leq 2 \times 10 \text{ kW} \\ n_e &\leq 5 \cdot 10^{18} \text{ m}^{-3} \text{ (for H)} \\ T_e &= 1 - 5 \text{ eV (for H)} \end{aligned}$$



Helicon plasma column

Laser beam
injection

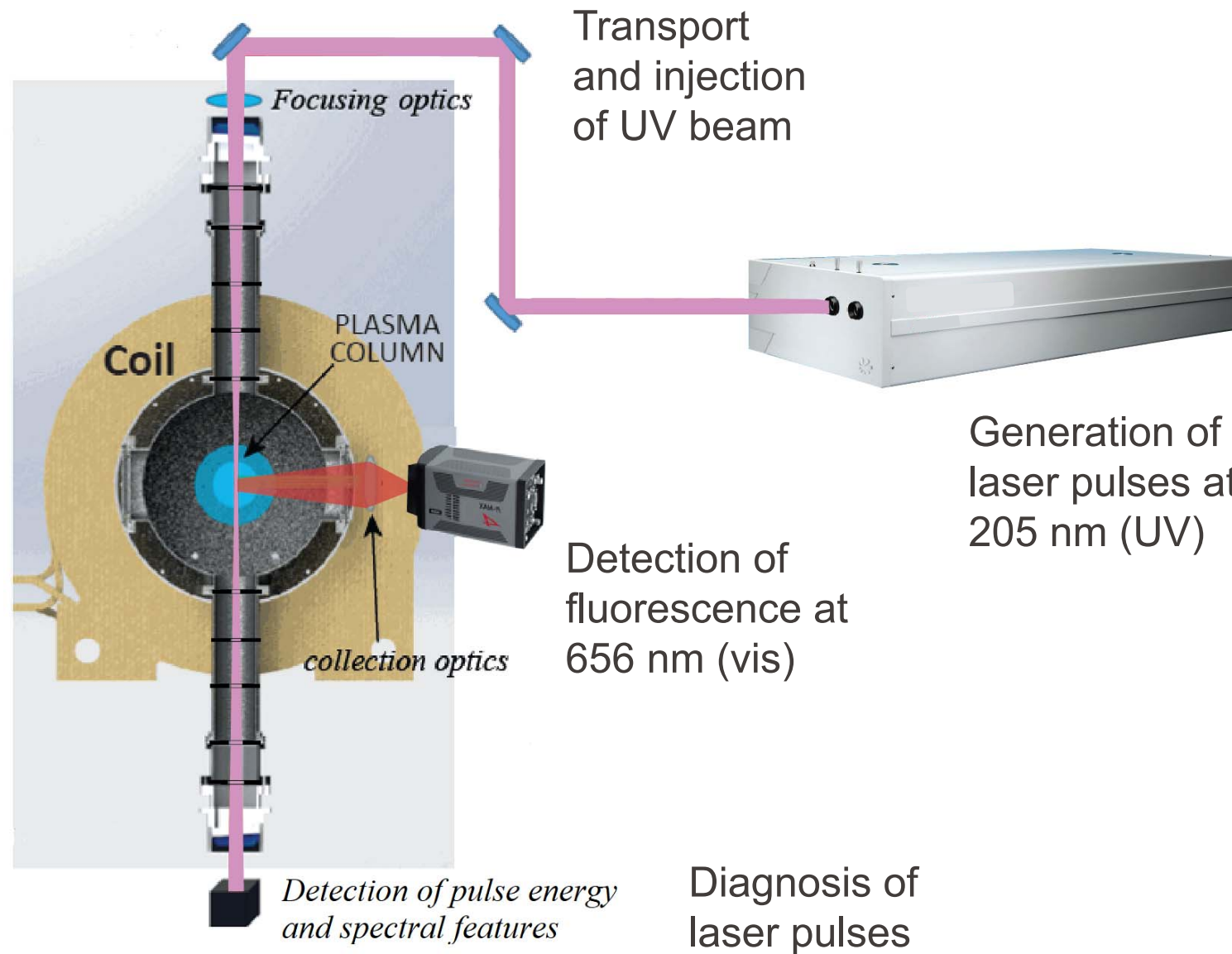
Window used to
observe fluorescence



Magnetic field coils

Vacuum vessel ($R = 20 \text{ cm}$, $L = 1.5 \text{ m}$)

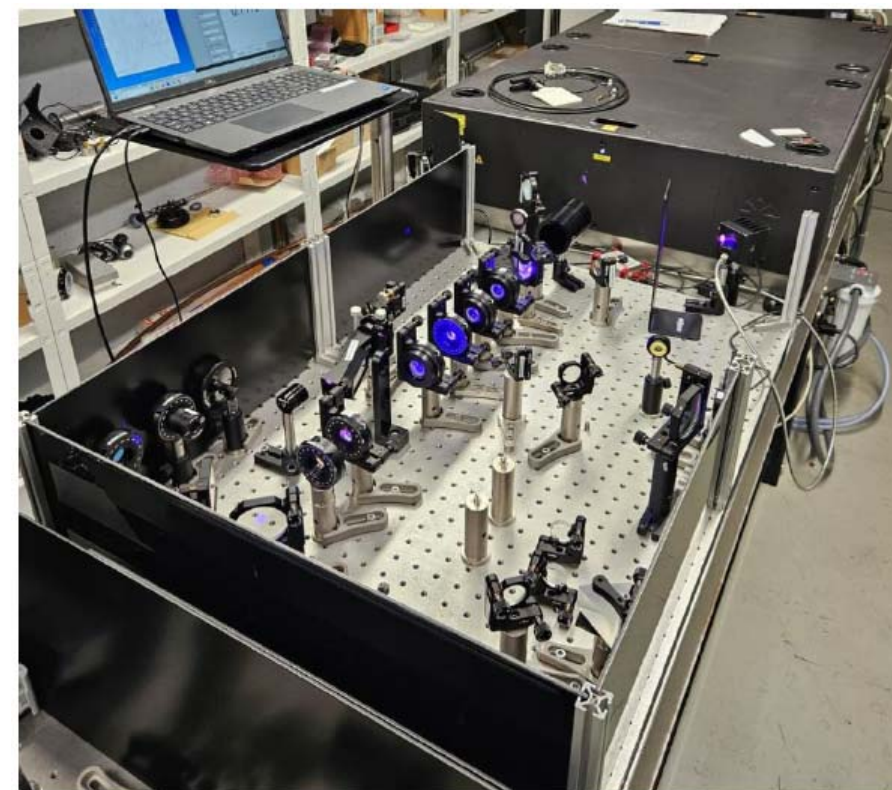
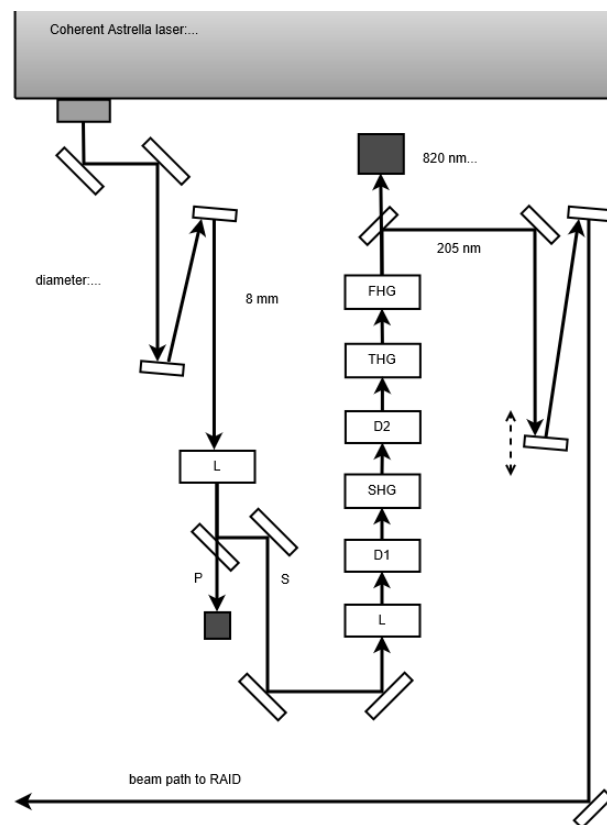
Introduction Proposed system



1. Introduction
2. Implementation
3. Experiments
4. Other studies
5. Scientific output
6. Conclusions and outlook

Implementation of laser (T1) FHG module and Astrella upgrade

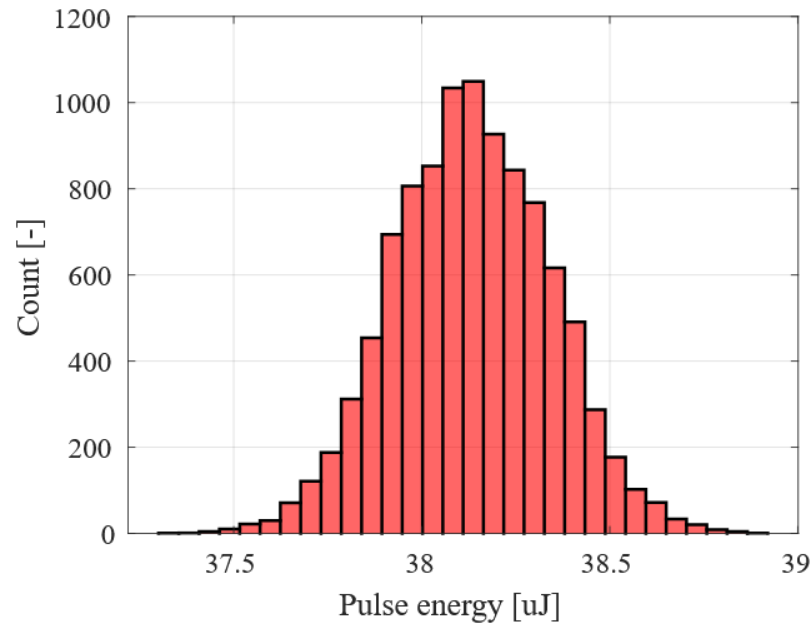
- We built a fourth harmonic generator (FHG) to obtain 205 nm from an upgraded Astrella system (1kHz rep-rate, 7 mJ per pulse at 820 nm, pulse duration = 70 fs).
- The FHG is based on the design proposed by Susnjar et al. *Opt. Comm.* (2023)



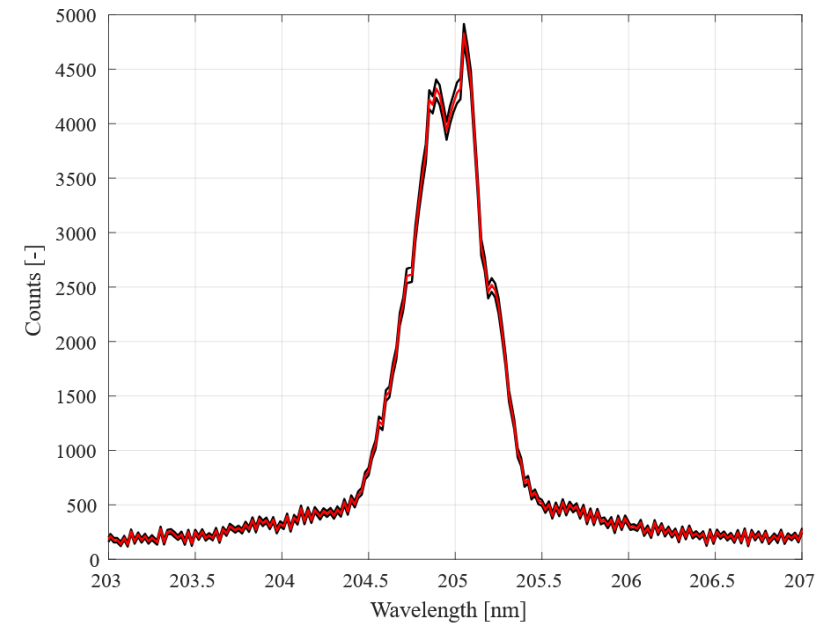
Implementation of laser (T1)

Characterization of UV fs pulses

- FHG and Astrella laser were configured for maximum energy while obtaining a wide and symmetric spectrum centered at 205.1 nm.



- Histogram of energy of 10000 pulses shows stable values.

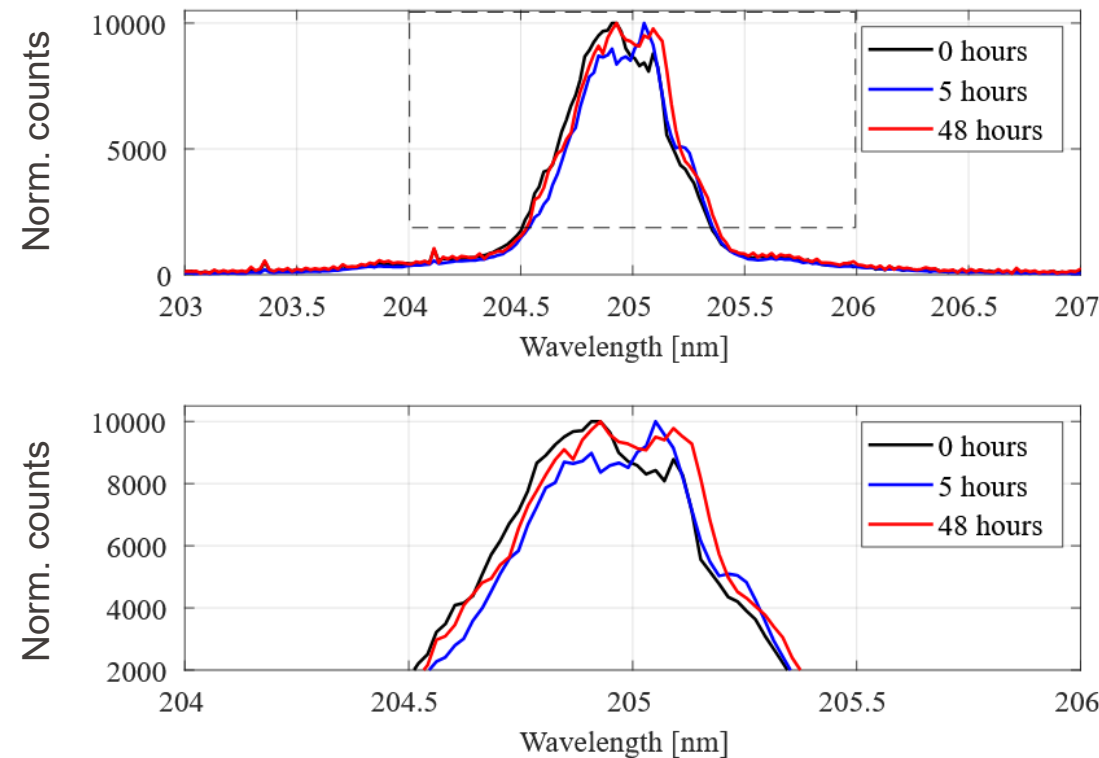


- Stable short-term spectral features. Red curve is the average of 1000 consecutive spectra, each one obtained from 100 pulses.

Implementation of laser (T1)

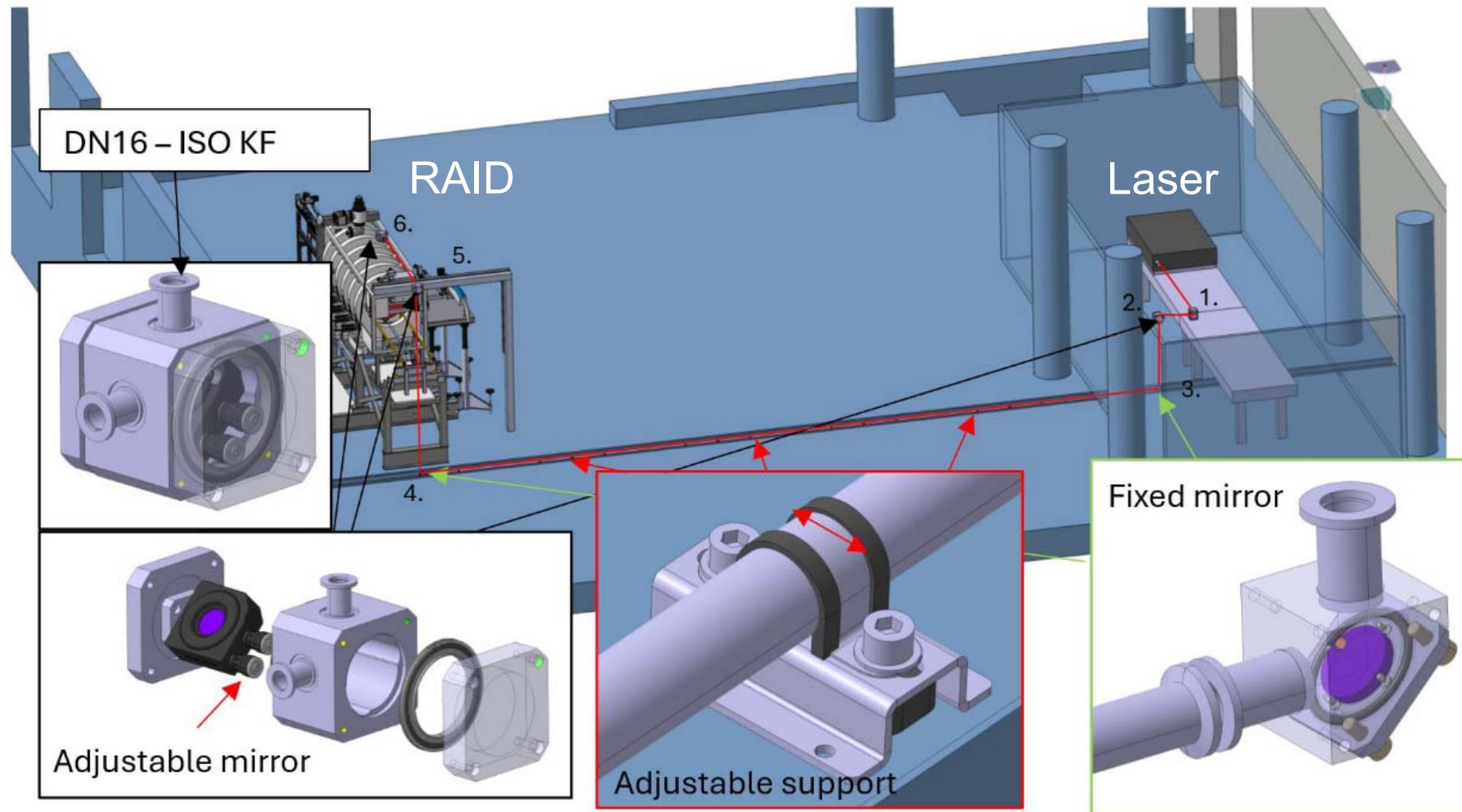
Characterization of UV fs pulses

- FHG and Astrella laser were configured for maximum energy while obtaining a wide and symmetric spectrum centered at 205.1 nm.



- Long-term spectral features are stable but require monitoring.

Implementation of beam path (T3)



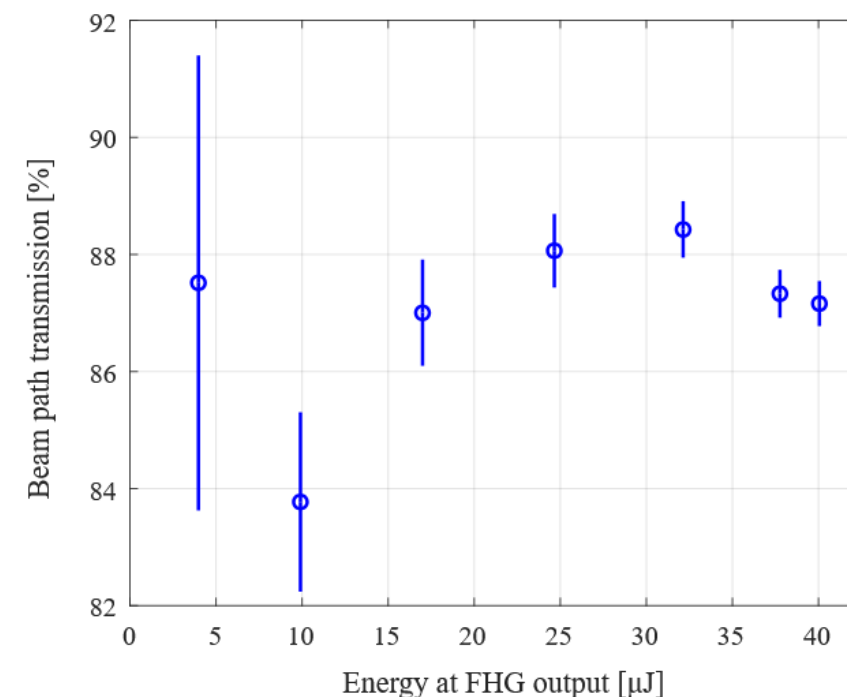
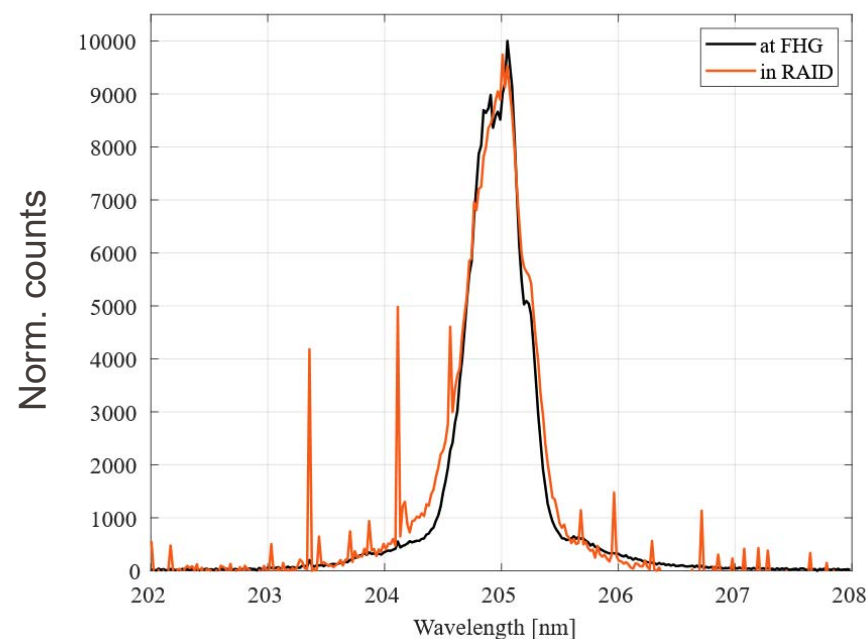


- Sealed beam path (designed for vacuum).
- Total length ≈ 19 m
- Refractive elements minimized to reduce dispersion and avoid damage of optics.
- Beam is loosely focused. Waist size can be changed by means of a reflective telescope at entrance of beam path.



Laser at exit of beam path

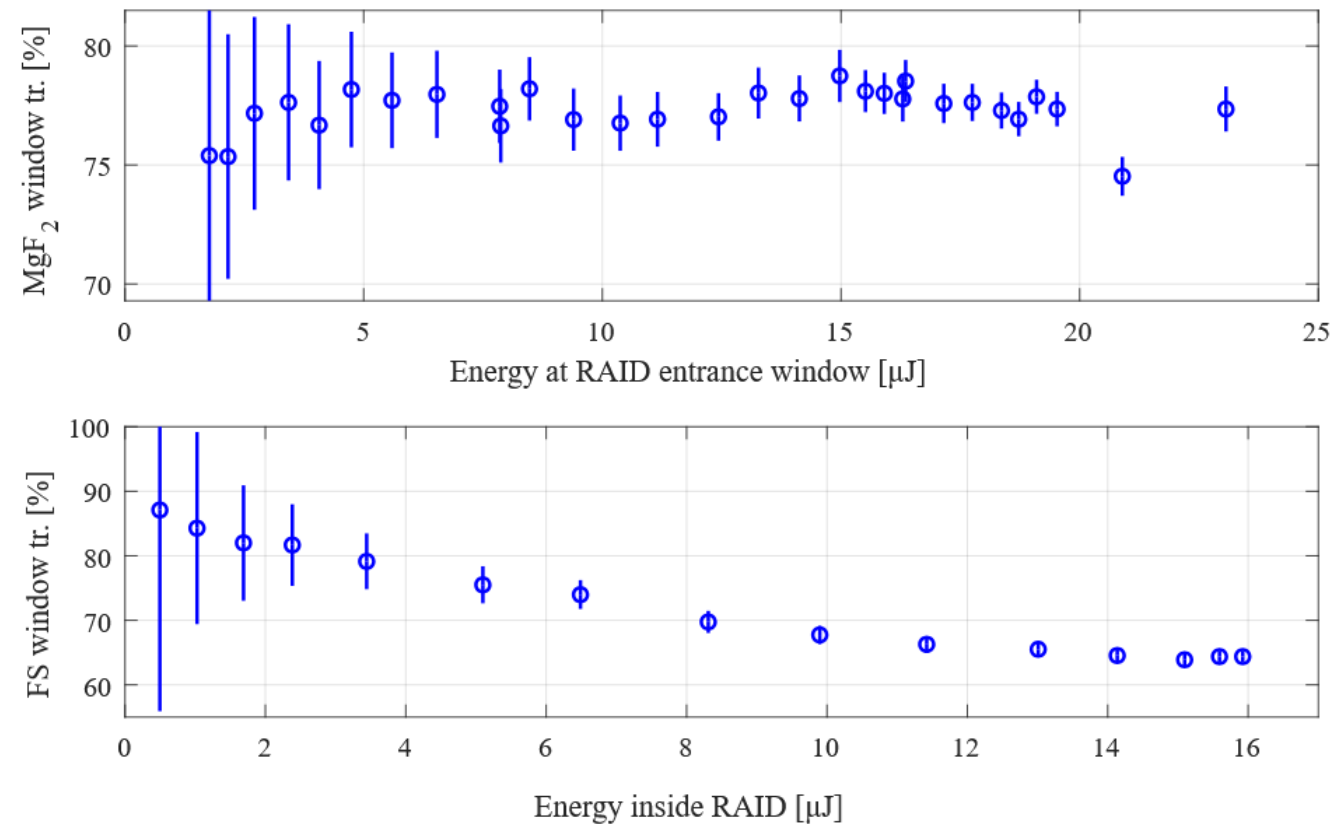
- Beam path is very efficient at transmitting the pulses and preserves their spectral features (temporal features will be studied in more detail).
- Refractive optics (windows, lenses, etc.) must be avoided as much as possible.



- Spectra measured at entrance and exit of beam path.

- Transmitted/input energy. Only $\approx 13\%$ of the pulse energy is lost in beam path.

- Windows lead to losses.



- Energy measurements are made after FS window (60% efficiency at highest energies).
- Pointing of laser may vary by a few mm over a few hours. We think this is due mainly to heating of the RAID vacuum chamber.

Detection of fluorescence (T2)

ICCD camera

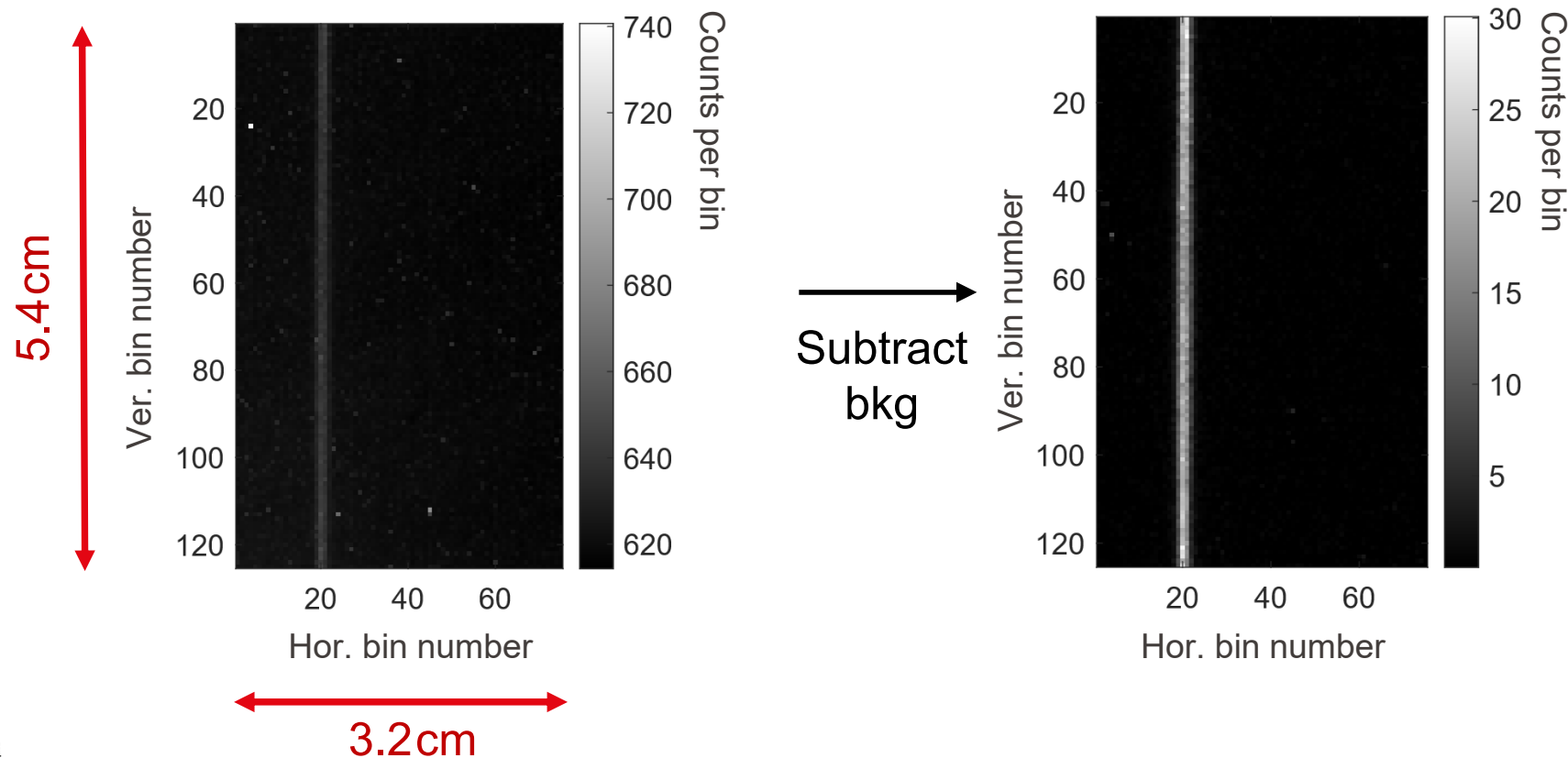


- Detection system based on imaging optics, optical filtering and a very short gating time (down to 3ns) ICCD camera.

1. Introduction
2. Implementation
3. Experiments
4. Other studies
5. Scientific output
6. Conclusions and outlook

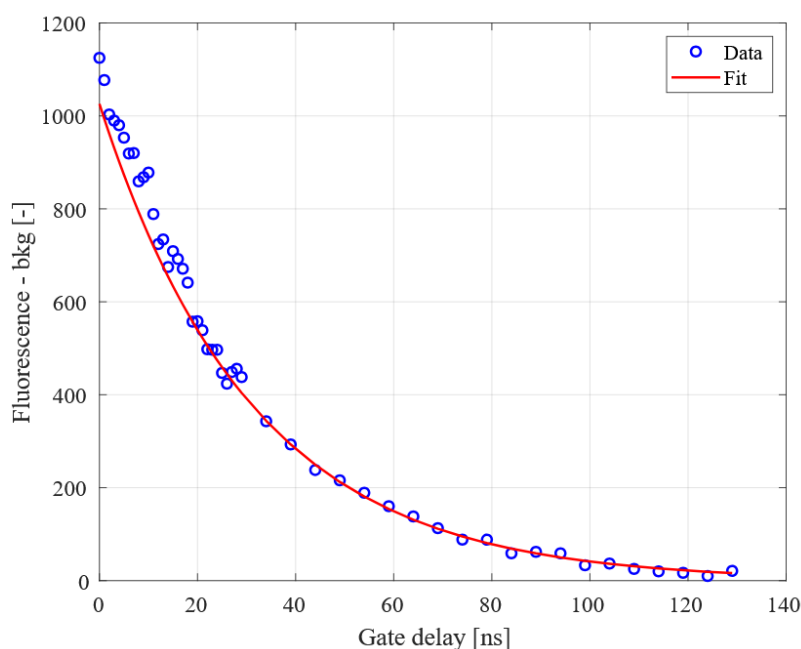
We obtain a signal with the laser tuned to 205.1 nm

- No plasma
- Gas fill pressure: 5.37 Pa
- Laser pulse energy: 10.3 μJ (measured; 17.1 μJ in RAID)
- 200 on-chip accumulations, 1000 frames averaged

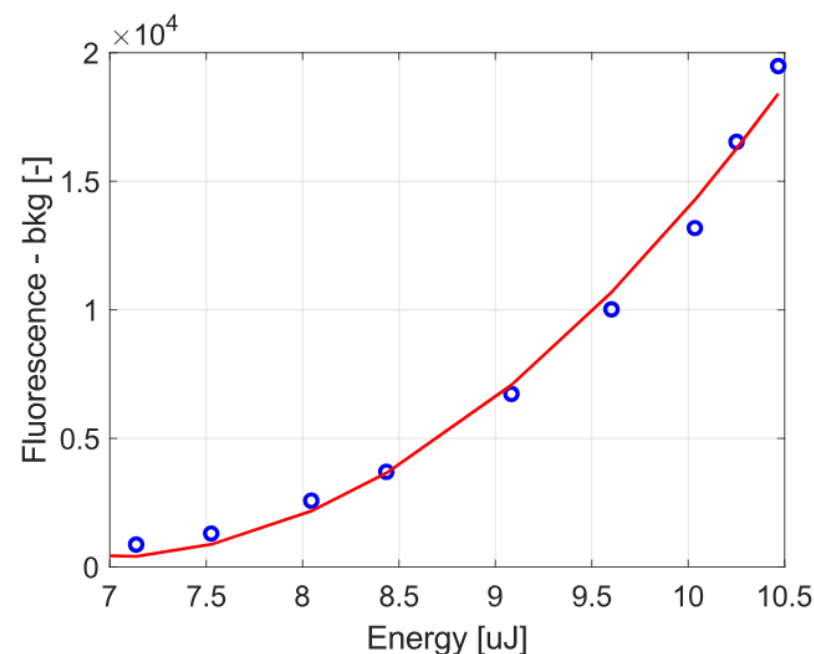


- Notice that beam is collimated within field of view.
- Diameter of beam ($1/e^2$) is $\approx 3\text{mm}$
- Signal is very well resolved in 1D

Signal has expected fluorescence decay rate
and seems to follow quadratic scaling with Intensity



- Decay rate of fluorescence ≈ 31 ns, close to expected value from literature (Kadi et al. *PPCF* (2024), Gazeli et al. *PoP* (2021)).

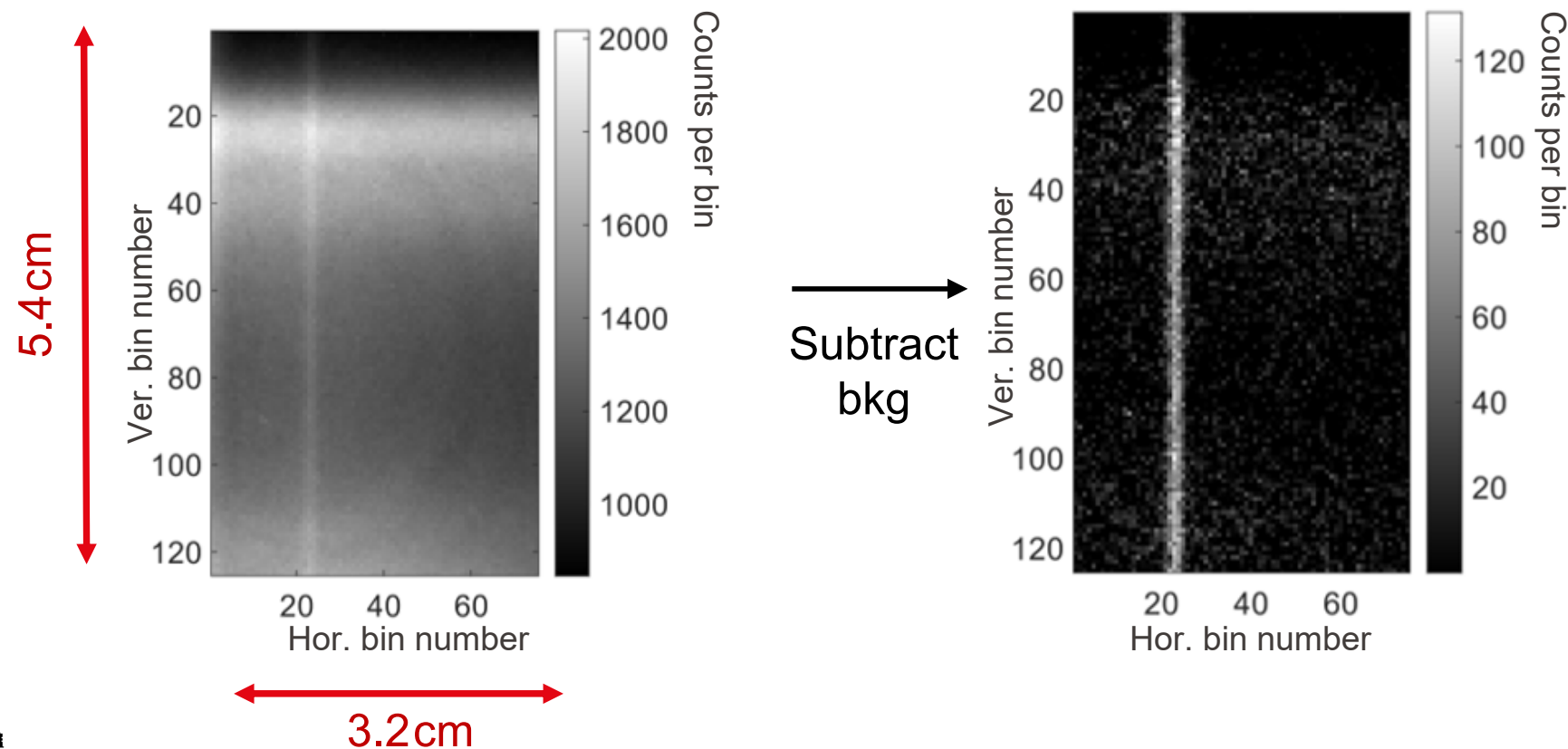


- Fluorescence scales quadratically with laser pulse energy starting from $\approx 6 \mu\text{J}$ (measured). That value includes the entire spectrum of the pulse.

fs-TALIF in H plasma

Clear signal with on-chip accumulation and frame averaging

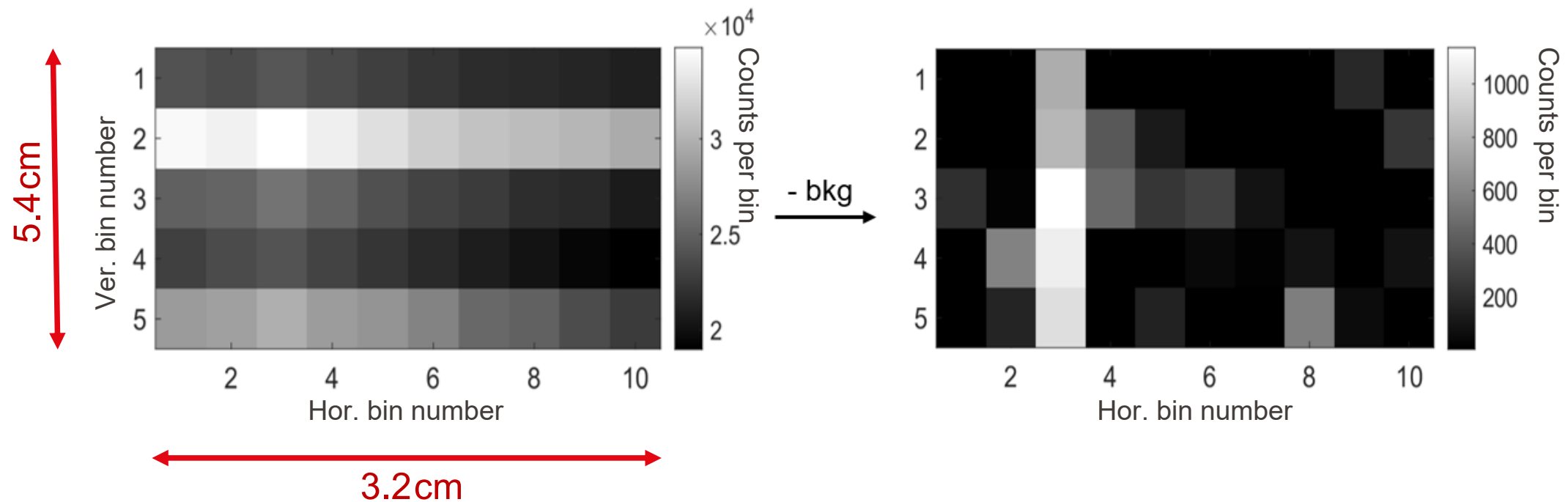
- H plasma, $P_{\text{RF}} = 1 \times 1 \text{ kW}$, $B_{\text{field}} = 260 \text{ G}$
- Gas fill pressure: 1.02-1.18 Pa
- Laser pulse energy: 10.3 μJ (measured; 17.1 μJ in RAID)
- 30 on-chip accumulations, 1000 frames averaged



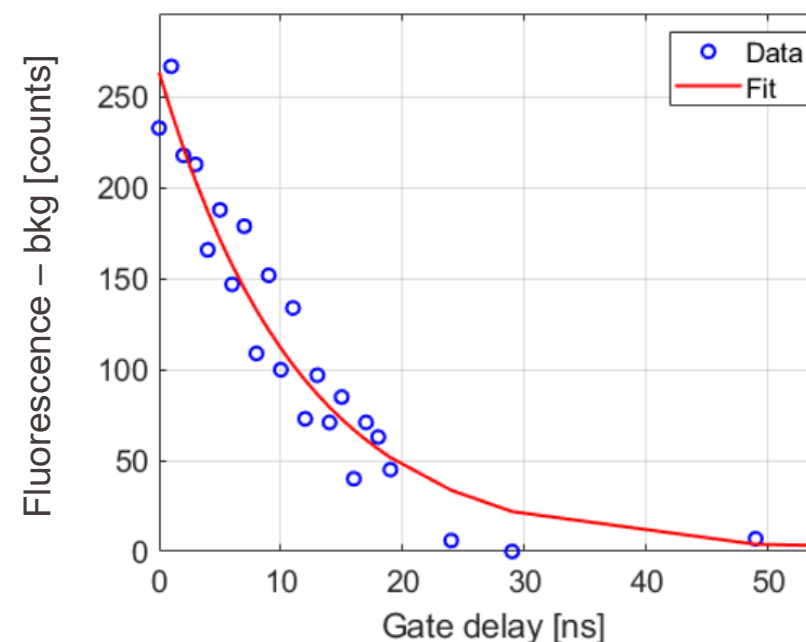
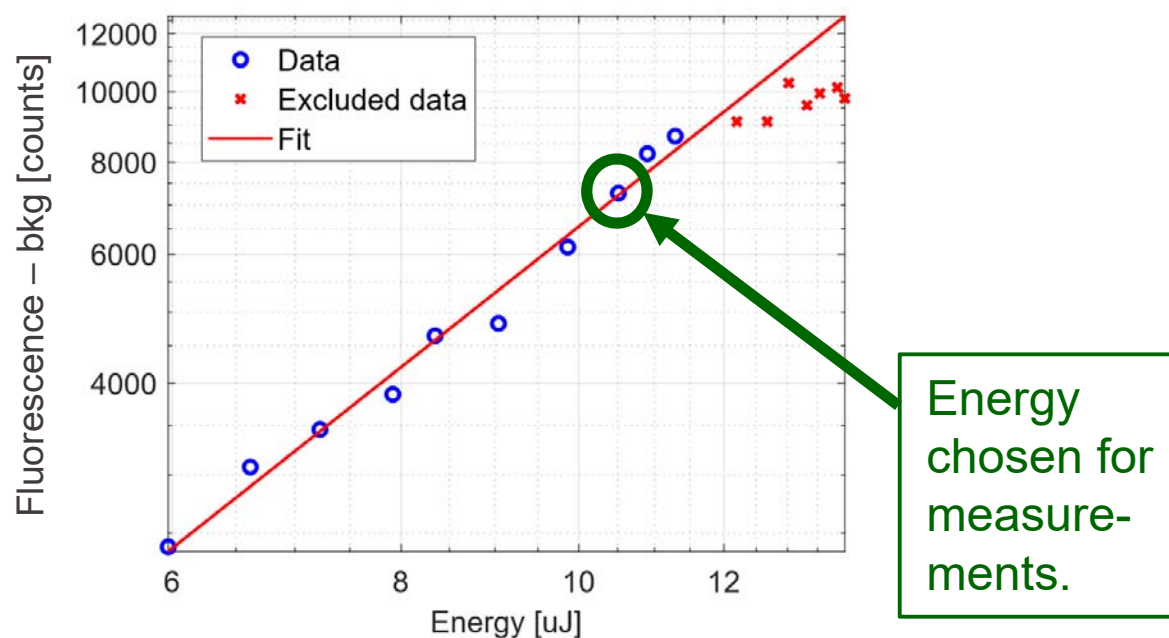
- Beam is collimated within field of view.
- Diameter of beam ($1/e^2$) is $\approx 3 \text{ mm}$
- Signal is very well resolved in 1D

Faster acquisition with larger bins

- H plasma, $P_{\text{RF}} = 1 \times 1 \text{ kW}$, $B_{\text{field}} = 260 \text{ G}$
- Gas fill pressure: 1.02-1.18 Pa
- Laser pulse energy: 10.3 μJ (measured; 17.1 μJ in RAID)
- 10 on-chip accumulations, 100 frames averaged



Beam waist chosen for maximum signal while staying in quadratic regime



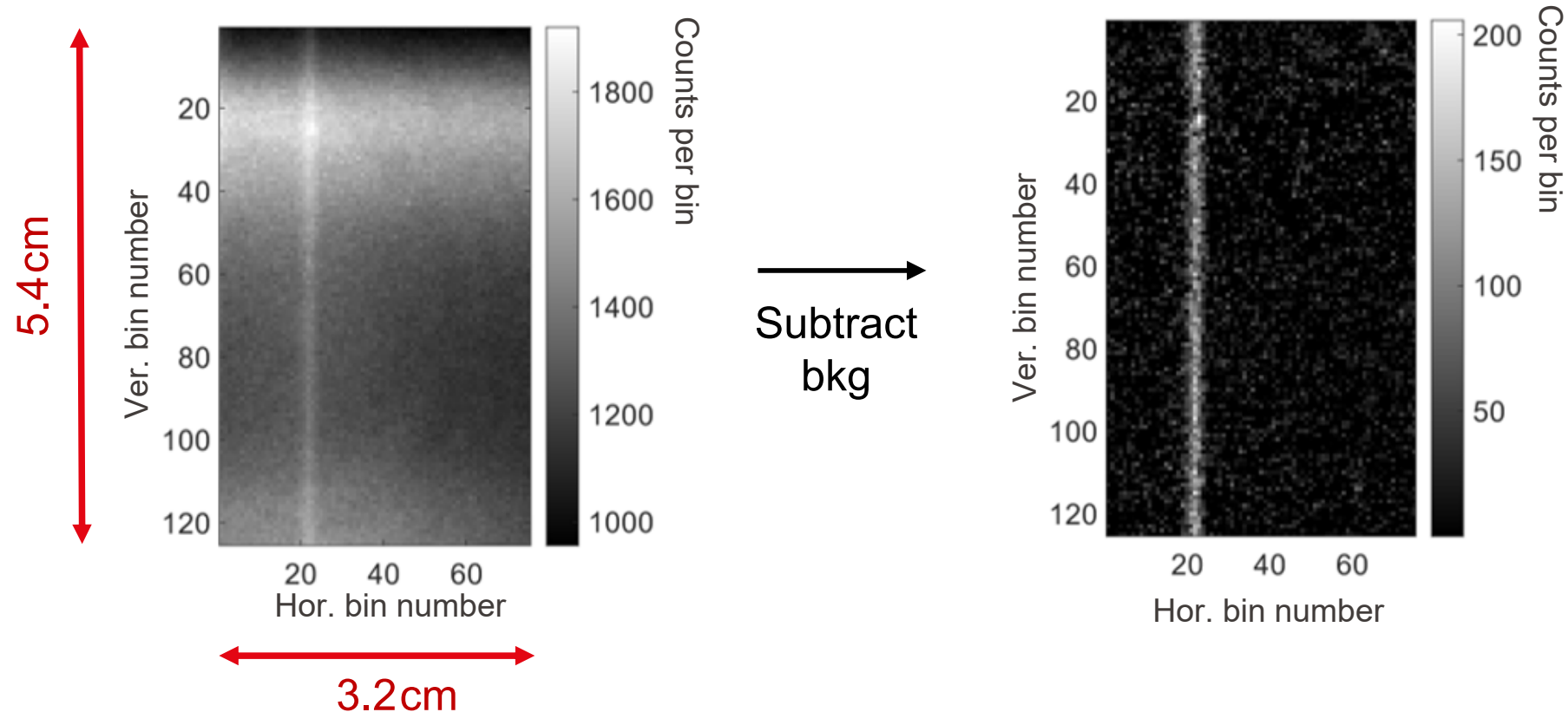
- Fluorescence signal scales quadratically with pulse energy up to 12 μJ (measured; 20 μJ in RAID).

- Decay rate of fluorescence $\approx 10\text{ns}$, close to expected value from literature (Kadi et al. *PPCF* (2024)).

fs-TALIF in D plasma

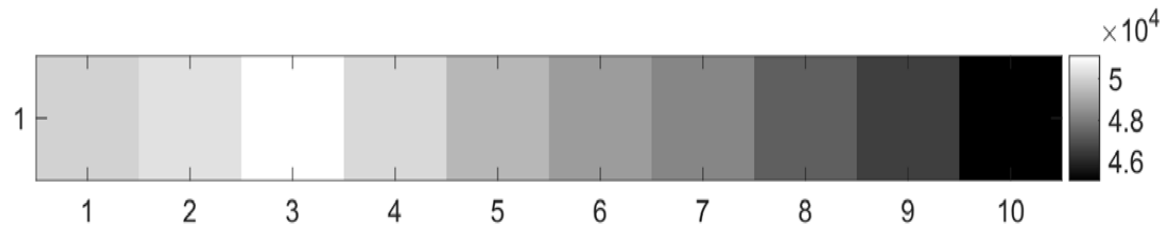
Results similar to H

- D plasma, $P_{\text{RF}} = 1 \times 1 \text{ kW}$, $B_{\text{field}} = 260 \text{ G}$
- Gas fill pressure: 1.1 Pa
- Laser pulse energy: 10.3 μJ (measured; 17.1 μJ in RAID)
- 30 on-chip accumulations, 1000 frames averaged

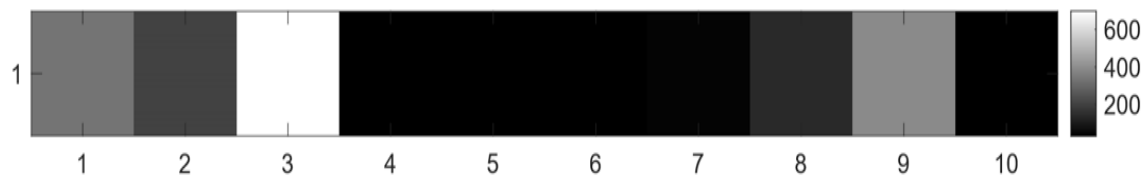


With coarser binning we can do

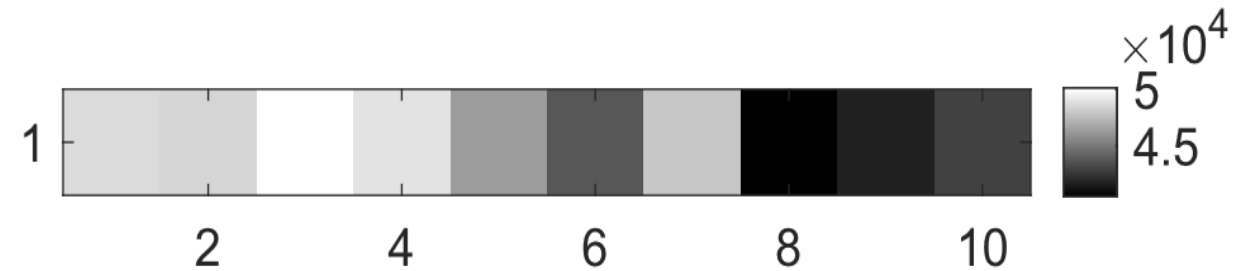
- D plasma, fill pressure: 0.58 Pa
- Full vertical bins
- 1 on-CCD accumulation, 100 frm. avg.



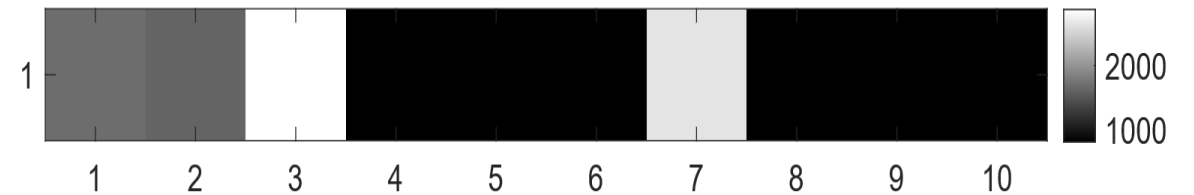
↓ - bkg



- D plasma, fill pressure: 0.58 Pa
- Full vertical bins
- 1 on-CCD accumulation, **no** frm. avg.



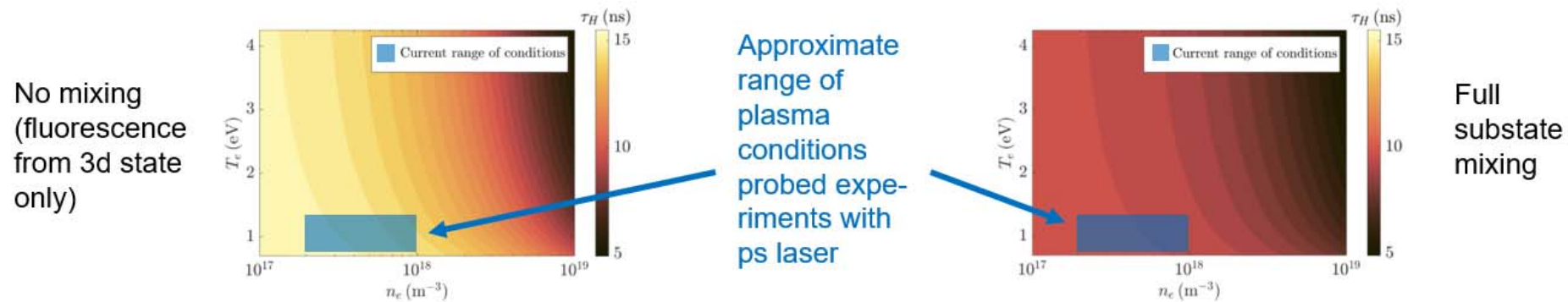
↓ - bkg



1. Introduction
2. Implementation
3. Experiments
4. Other studies
5. Scientific output
6. Conclusions and outlook

Studies of n=3 substate mixing (T5)

- Simulations of H fluorescence lifetimes suggest that electron collisional processes are not sufficient to explain observations made in (separate) experiments with our ps laser (Kadi et al. *PPCF* (2024)).



- If one includes full substate mixing of n=3 (i.e. all m levels are quickly populated statistically upon laser excitation), lifetimes become compatible with observations.
- If indeed present, full substate mixing makes value of branching ratio very robust up until $n_e \approx 3 \times 10^{18} \text{ m}^{-3}$.
 - Furthermore, the value of the branching ratio changes. Its impact in the calculation of n_H leads to an increase of $\approx 26\%$ in the density.
- Further theoretical and experimental studies are required.

- Even though widely used, the model for two-photon absorption given in introduction may have limitations in fs regime:
 - Pulse intensities are very large
 - Rate equations may neglect important quantum effects
 - Laser photon statistics may differ compared to ps and ns cases
- Using a quantum model (based on Bloch equations) we have established that model does apply for laser intensities up to (at least) a few GW/cm^2 .
- From literature, laser photon statistics factor may have a different value.
- Further theoretical studies are being pursued.
- Experimental studies with pulse compressor will allow us to study effects of pulse duration.

1. Introduction
2. Implementation
3. Experiments
4. Other studies
5. Scientific output
6. Conclusions and outlook

Conference contributions

Authors	Conference name	Title
S.P.H. Vincent, et al.	67th Annual Meeting of the APS Division of Plasma Physics (2025)	Poster: Femtosecond TALIF for H density measurements
I. Furno, et al.	30th IAEA Fusion Energy Conference (2025)	Poster: Active spectroscopy for atomic H and D measurements in fusion
M. Baquero-Ruiz, et al.	21st International Symposium on Laser-Aided Plasma Diagnostics (2025)	Poster: Femtosecond TALIF of atomic H in RAID
M. Goddijn, et al.	6th European Conference on Plasma Diagnostics (2025)	Poster: Femtosecond TALIF diagnostic on the RAID linear device
M. Goddijn, et al.	Joint Annual Meeting of the Austrian and Swiss Physical Societies (2025)	Oral: Femtosecond TALIF in fusion-relevant hydrogen plasmas

1. Introduction
2. Implementation
3. Experiments
4. Other studies
5. Scientific output
6. Conclusions and outlook

- We have developed a fs TALIF system to study H (and/or D) densities in RAID
 - Laser pulses are stable in energy and spectral features.
 - Frequent monitoring may nevertheless be required in applications.
- Using averaging in detection, we obtain clear fs-TALIF signals resolved in 1D
 - Laser beam is collimated over field of view.
 - Successful observations in H plasma
 - Successful observations in D plasma
 - Successful observations in Kr gas. Calibrations seem to be possible without changing wavelength.
- Very coarse binning allows us to run fast
 - Nonetheless, single laser pulse measurements are **not yet possible**.
 - We believe that we are at the limit of what we can achieve with the ICCD.

- We will pursue the measurements of n_H
 - We believe that the results presented here show very good potential.
- We will install a streak camera
 - Resolving the background immediately prior to the laser pulse may allow us to better distinguish signal over noise.
 - Camera and optics are ready. Experiments are planned for early 2026.
- It may be possible to increase pulse energies while having the spectrum centered at 205.1 nm
 - With system tuned closer to 206 nm we obtain at least twice the energy.
 - We think that thicker delay crystals in FHG may allow us to obtain this at 205.1 nm.
- Studies of full substate mixing and pulse absorption will continue
 - Installation of pulse compressor and experiments will be interesting to study effect of pulse duration.

Thank you for your attention