# **KU LEUVEN**



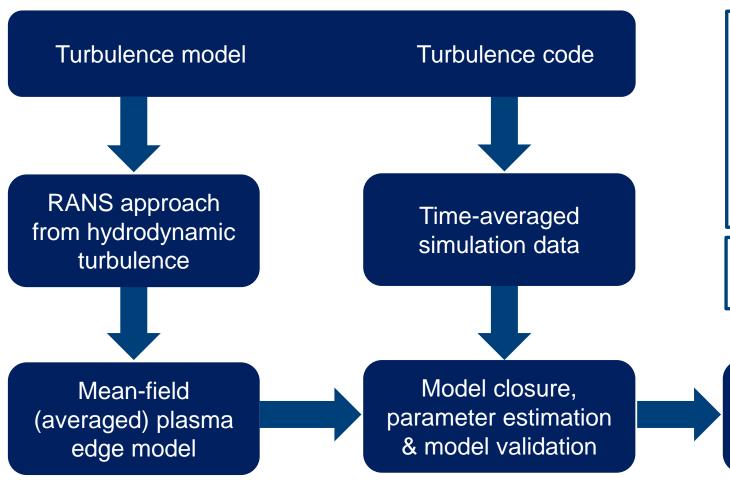
### Progress in modeling TCVX23 with k-model

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# Approach



Split in mean + fluctuating components

$$x = \bar{x} + x', \qquad \bar{x} = \lim_{T \to \infty} \frac{1}{T} \int_0^T x dt$$

$$x = \tilde{x} + x'', \qquad \tilde{x} = \frac{\overline{nx}}{\overline{n}}$$

Time-average governing equations (e.g.  $\partial_t \overline{n} + \nabla \cdot \overline{\Gamma} = \overline{S}$ )

Mean-field transport model consistent with turbulence model



#### Application to electrostatic turbulence (no B-fluctuations)

#### Average fluxes

$$\overline{\boldsymbol{Q}}_{i} = \overline{Q}_{i||} \mathbf{b} + \frac{5}{2} \left( \overline{n} \ \widetilde{\boldsymbol{u}}_{*} + \overline{n} \widetilde{\boldsymbol{u}}_{p} + \overline{n} \widetilde{\boldsymbol{u}}_{v} \right) \widetilde{T}_{i} + \frac{3}{2} \overline{n} \ \overline{\boldsymbol{u}}_{E \times B} \widetilde{T}_{i} + \frac{3}{2} \left( \overline{n' \boldsymbol{u}'_{E \times B}} \widetilde{T}_{i} + \overline{n \boldsymbol{u}''_{E \times B}} \widetilde{T}''_{i} \right) - \overline{\kappa_{i}} \nabla T_{i}$$
(for  $\overline{\boldsymbol{Q}}_{i}$ , neglecting some fluctuations of smaller drifts)

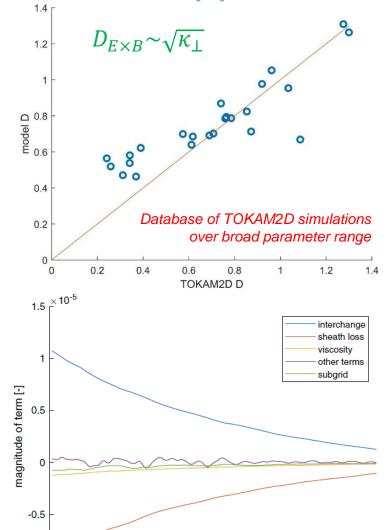
where 
$$\bar{n} \ \widetilde{\boldsymbol{u}}_* = \frac{1}{eB^2} \nabla \ \bar{p}_i \times \boldsymbol{B}$$
,  $\bar{n} \widetilde{\boldsymbol{u}}_p = \frac{\mathbf{b}}{eB} \times \left( \frac{\partial}{\partial t} \left( m \bar{n} \widetilde{\boldsymbol{u}} \right) + \nabla \cdot \left( m \bar{n} \boldsymbol{u} \overline{\boldsymbol{u}} \right) \right)$ ,  $\bar{n} \widetilde{\boldsymbol{u}}_{\nu} = \frac{\mathbf{b}}{eB} \times \nabla \cdot \overline{\Pi}_i$ , and  $\bar{p}_{i/e} = \bar{n} \widetilde{T}_{i/e}$ ,  $\bar{\boldsymbol{u}}_{E \times B} = \frac{1}{B^2} \boldsymbol{B} \times \nabla \ \bar{\phi}$ 

• Fluctuating  $E \times B$ -terms need closure; stick to diffusive closure (for now):

$$\circ \quad \overline{\mathbf{\Gamma}}_{i/e,E\times B} = \overline{n_i' \mathbf{u}_{E\times B}'} \qquad \sim -D_{E\times B} \nabla \overline{n}_i$$



#### RANS approach for electrostatic interchange turbulence



• Proposal: relate to turbulent kinetic energy  $\kappa_{\perp}$  as measure of local intensity of the turbulence/transport

$$\overline{n}_i \kappa_{\perp} = \frac{\overline{n_i m_i \mathbf{u}_{E \times B}^{\prime\prime 2}}}{2}$$

• Diffusive transport model based on 2D interchange simulations [Coosemans et al., CPP 2022, e202100193.]

$$O \qquad D_{E \times B} = C_D \rho_L \sqrt{\frac{\kappa_\perp}{m_i}}$$

$$\circ \quad \chi_{i/e,E\times B} = C_{i/e}D_{E\times B} \sim \frac{3}{2}D_{E\times B}$$

• Need transport equation for  $\kappa_{\perp}!$ 

50

radial position [gyro-radii]

100

150

### Sources and transport of $\kappa_{\perp}$ for interchange turbulence

[R. Coosemans et al., J. Plasma Phys. 90 (2024) 90590020.]

[1e5 Wm<sup>-3</sup>]

• Transport equation for  $\kappa_{\perp}$  for 2D electrostatic interchange turbulence

$$\frac{\partial}{\partial t} \bar{n} \kappa_{\perp} + \nabla \cdot \left( \overline{\Gamma}_{i} \kappa_{\perp} + \overline{\phi' J'_{||}} \right) = \overline{S}_{IC} + \overline{S}_{||}$$

• Total heat flux due to  $E \times B$  fluctuations drives production of  $\kappa_{\perp}$ 

$$\overline{S}_{IC} = -\frac{2}{3} \overline{\boldsymbol{Q}}_{i+e,E\times B} \cdot \nabla \ln B^2$$

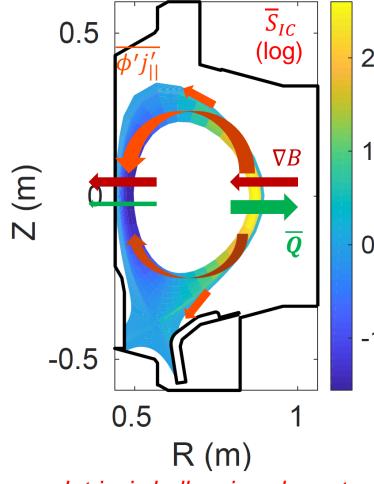
- Source in 'bad-curvature' regions, sink in 'good-curvature' regions
- Parallel transport of  $\kappa_{\perp}$  governed by plasma conductivity

$$\overline{\phi'j'_{||}} \sim -\sigma_{||}\nabla_{||} \frac{\overline{\phi'^2}}{2} \sim -C_{\sigma 1}\sigma_{||}\rho_L^2 B^2 \nabla_{||}\kappa_{\perp}$$

- $_{\circ}$  Strongly exceeds parallel convection with  $ilde{u}_{||}$
- Proxy to account for mean ExB flow shear  $S_m$  suppressing anomalous transport

$$D_{E \times B} = C_D \frac{\kappa_{\perp}/m_i}{\sqrt{\kappa_{\perp}/m_i}/\rho_L + C_S|S_m|}$$

...coupled to 'regular' mean field equations [Dekeyser et al., CPP 2022, e202100190.]

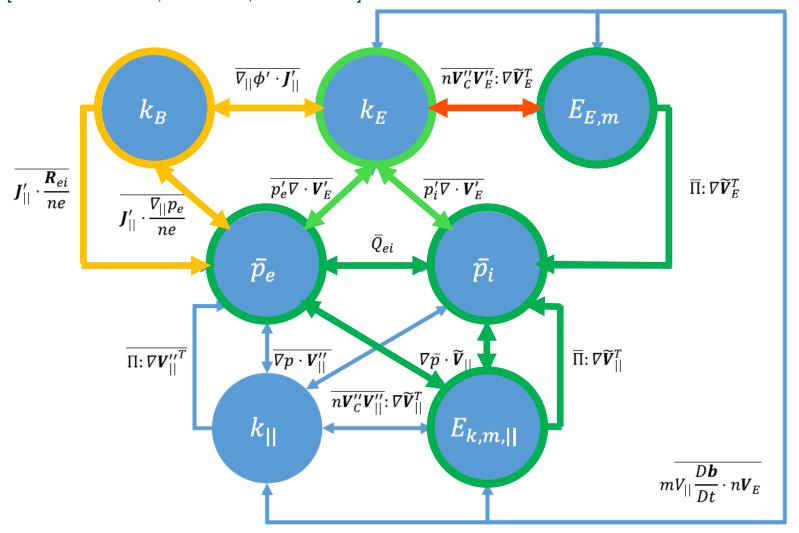


Intrinsic ballooning character and 'self-saturation'



## Energy transfer channels

[Coosemans et al., JPP 2024, 905900202.]



Default in mean-field codes

*k*-model

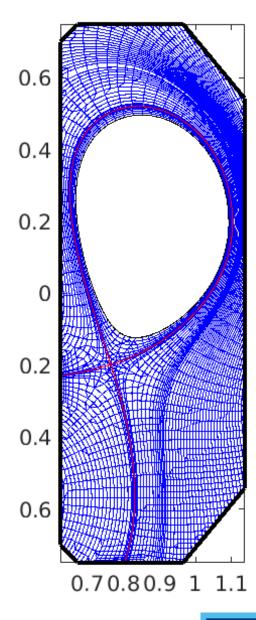
Simplified parallel dynamics in *k*-model

Ad-hoc RS model



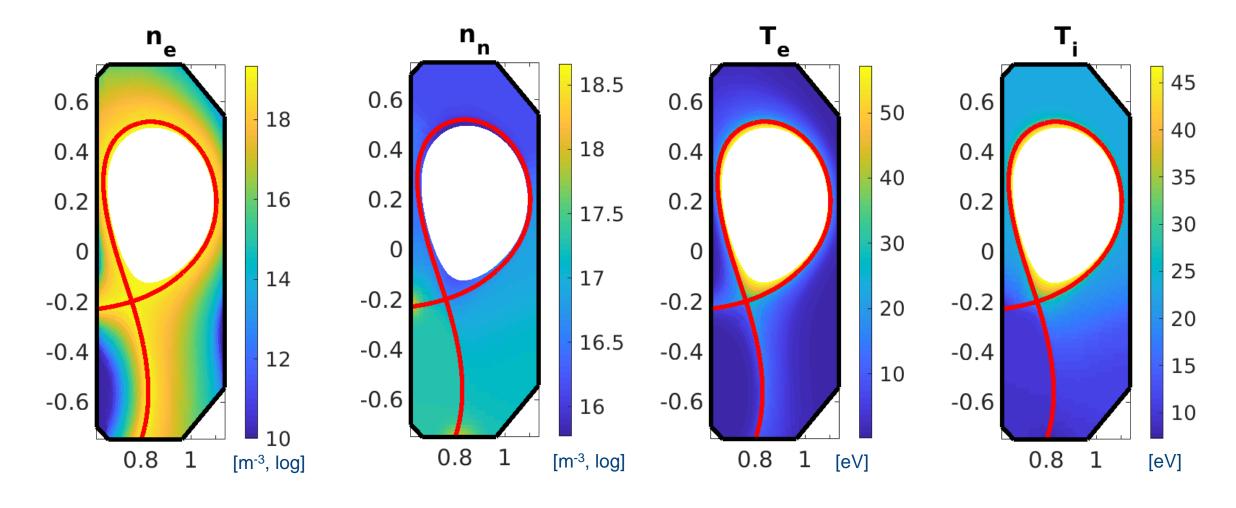
#### TCV-X23 showcase

- First-of-a-kind simulations including
  - Extended grids
  - AFN model
  - *k*-model for anomalous transport
  - Drifts
- Grid built with GOAT, exploiting new boundary layer grid option [S. Van den Kerkhof et al., PET 2025]
- Simulation setup based on TCV-X23 experiment [D. Mancini et al. Nucl. Fusion 64 (2024) 016012.]
- $\kappa_{\perp}$  -model parameters optimized via LSQ [S. Carli et al., PET 2025] based on TCV-X21 data [D.S. Oliveira et al 2022 Nucl. Fusion 62 096001]; compare with constant  $D_{\perp} \sim 0.4~m^2/s$





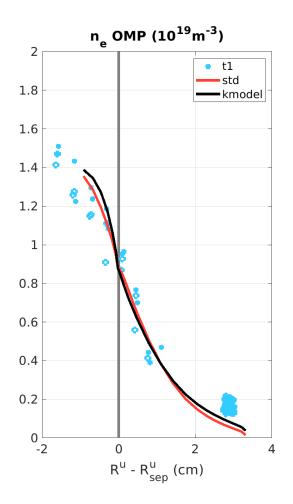
# TCV X23 edge plasma parameters

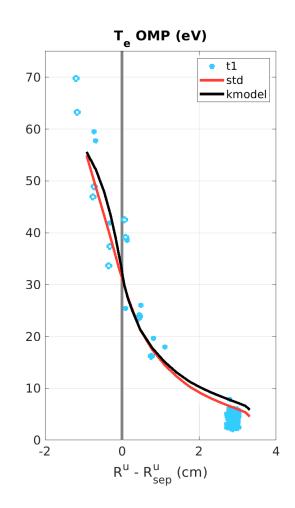


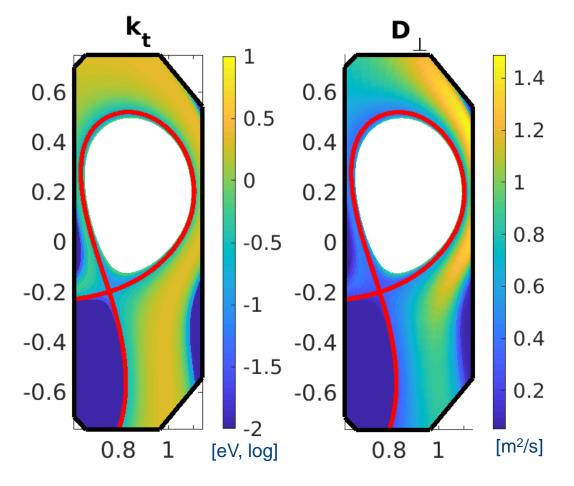


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# TCV-X23 showcase: turbulent kinetic energy and $D_{\perp}$









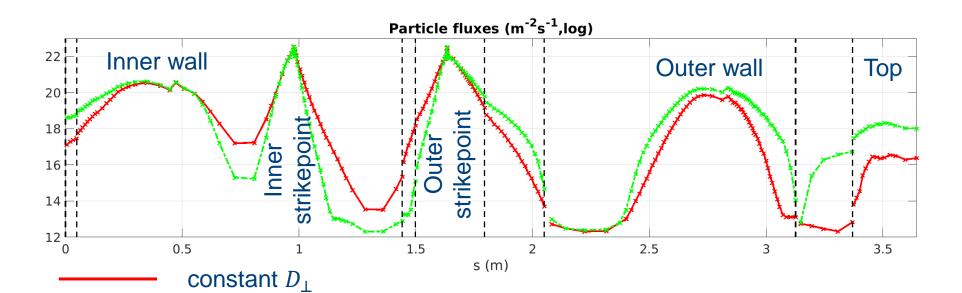
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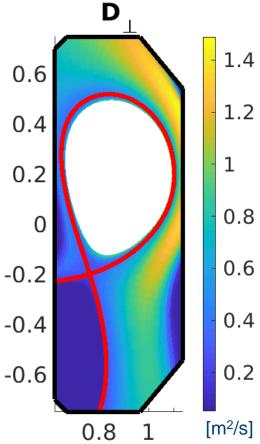
#### TCV-X23 showcase: wall fluxes

- Clear signature of magnetic field structure on wall loads
- Peak loads near strike points similar

 $\kappa_{\perp}$ -model

• Up to factor two higher wall fluxes near outboard side for  $\kappa_{\perp}$  -model, despite nearly identical density/temperature profiles







20/10/2025

### Summary and challenges

- Anomalous transport model based on RANS approach integrated in SOLPS-ITER
  - Includes features such as ballooned transport and self-saturation, but also remaining uncertain closure constants
  - Need for extensive validation with turbulence simulation data and experiment, incl. predicted turbulence characteristics
  - Need for continued development: additional energy exchange mechanisms, updated closures, BCs,...
- First application to TCV-X23 case
  - Including drifts and grids-to-the-wall
  - As for TCV-X21: limited ballooning in near SOL for the selected model parameters
  - Pronounced increase in k towards outer wall, and 'notable' changes in wall fluxes. Potential to discriminate between models?



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Thank you! Questions?

