



EUROfusion

Flux-tube simulations of TAE modes (work in progress)

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Acknowledgements: D. Brioschi, T. Görler, T. Hayward-Schneider, R. Kleiber, A. Könies, A. Di Siena



MAX-PLANCK-INSTITUT
FÜR PLASMAPHYSIK



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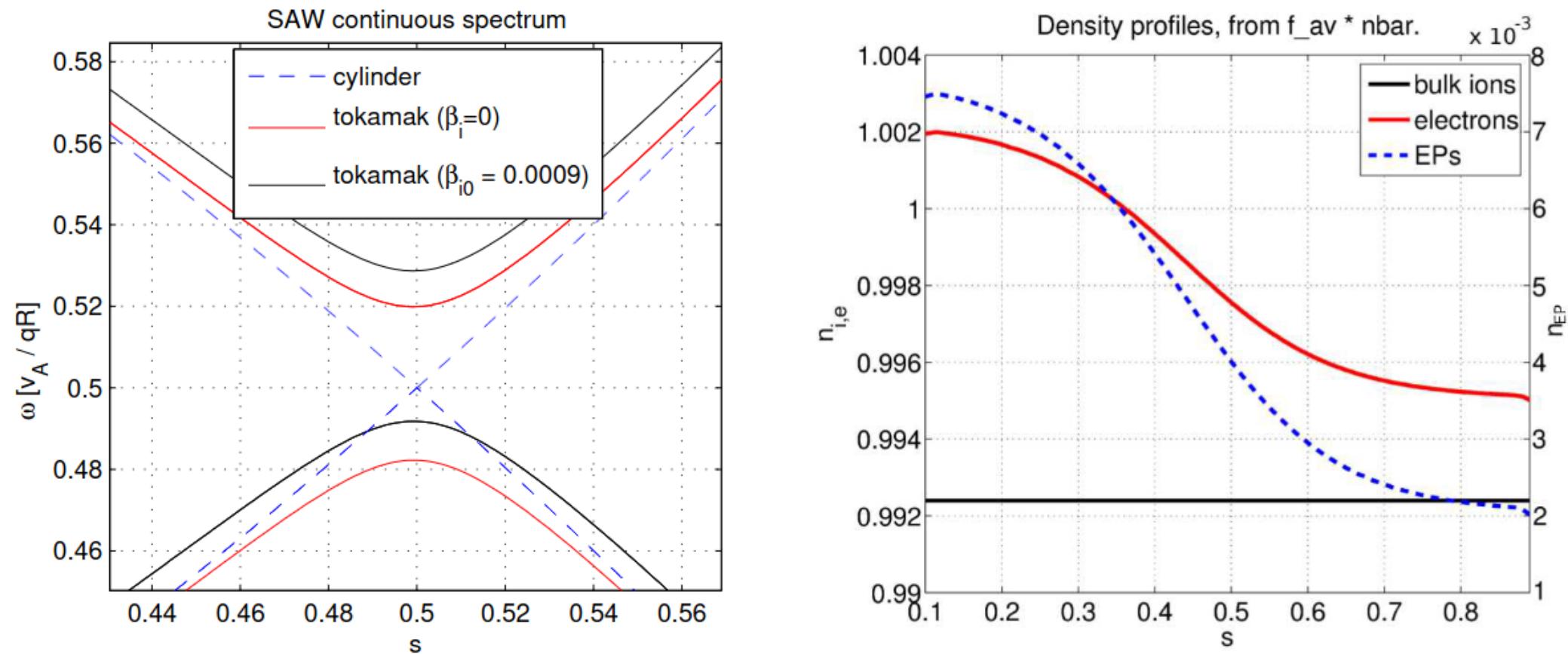
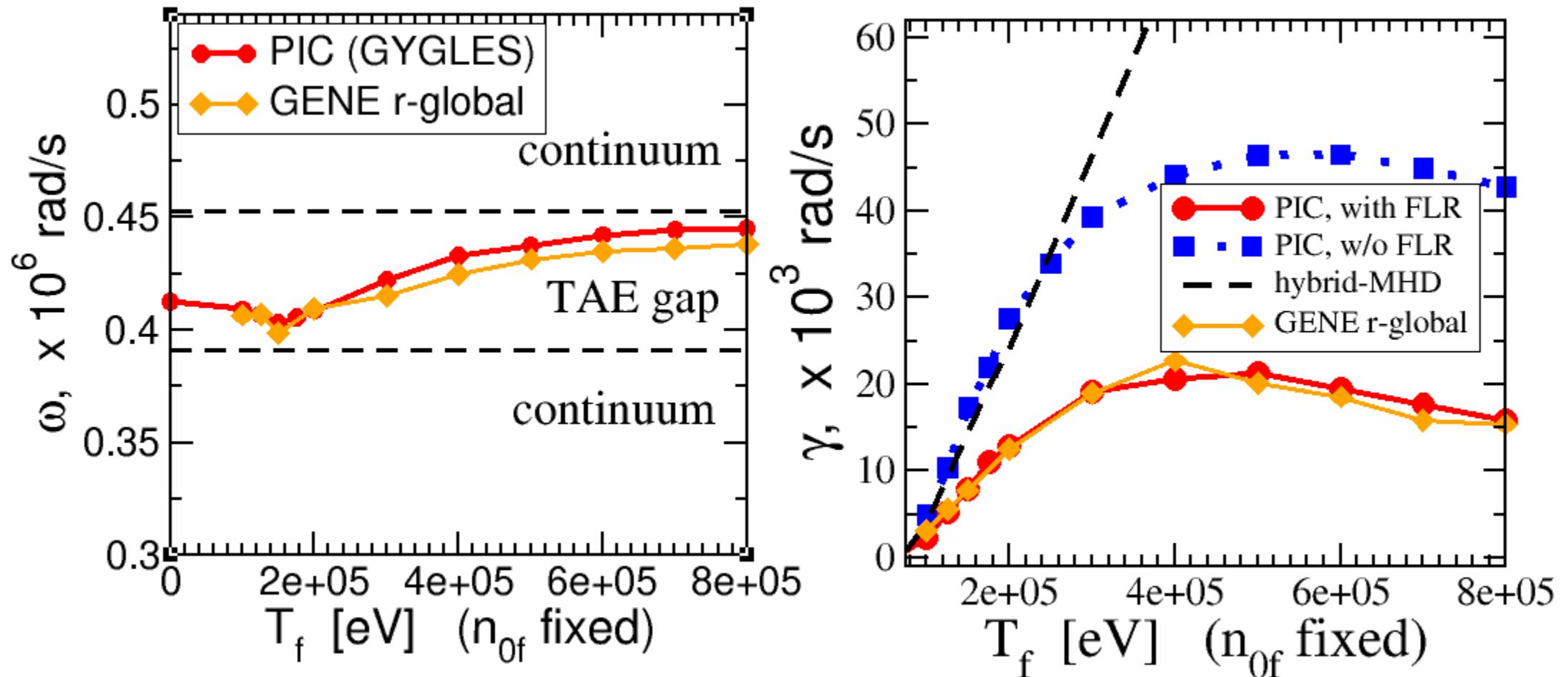
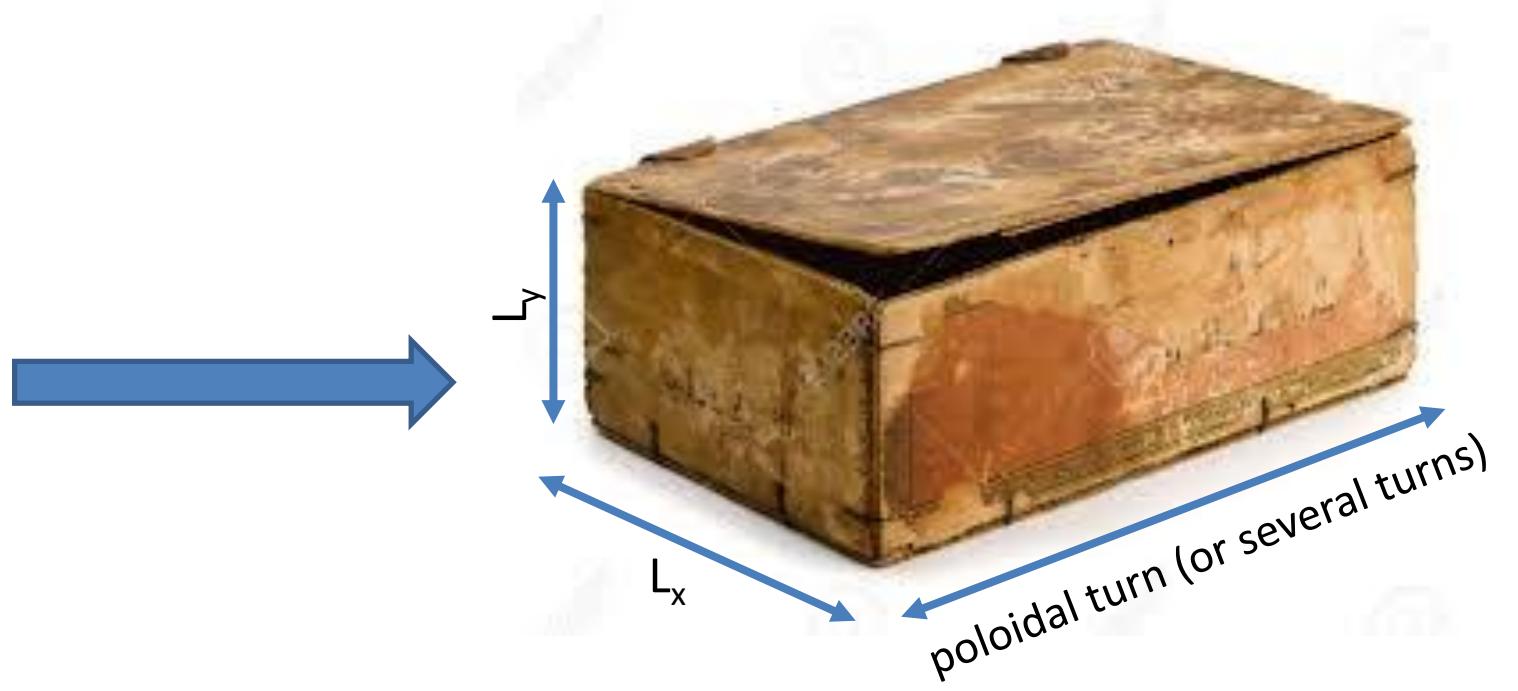
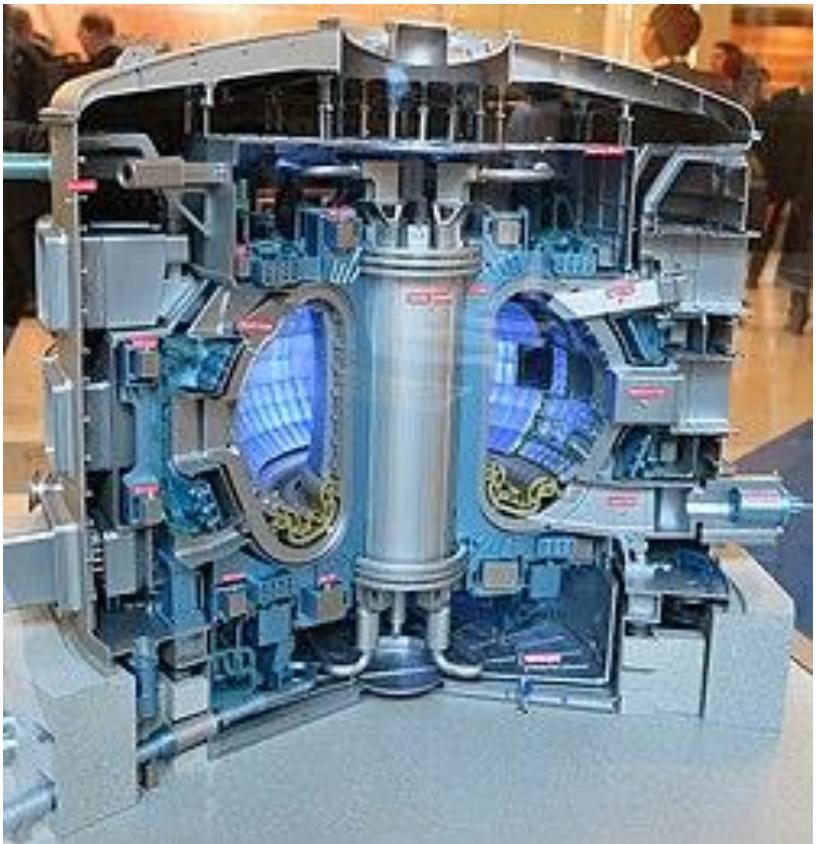


Figure 1: On the left, zoom of the continuous spectrum obtained by neglecting toroidicity and compressibility (blue dashed lines), or by keeping toroidicity only (red line) or by keeping both toroidicity and compressibility (black line), with approximated formula of Ref. [21]. On the right, density profiles in units of $n_{e0} = 2 \cdot 10^{19} m^{-3}$, for a case with $\langle n_{EP} \rangle / n_{e0} = 0.004$.



Radially-global GENE linear frequency and growth rate compare very well to GY GLES (even maybe slightly better than Davide's comparison with ORB5 performed last year).

GY GLES simulations: [Phys. Plasmas 16, 082105 \(2009\)](https://doi.org/10.1063/1.3160520) ; ITPA benchmark: [A. Könies et al 2018 Nucl. Fusion 58 126027](https://doi.org/10.1088/1741-4326/aa9a20)



In the flux-tube approach, one replaces a torus with a box

Poloidal box size is determined by spectral resolution: $L_y = 2 \pi / k_{y,\min}$

Radial box size is determined by „twist-and-shift“ condition (parallel bc): $L_x = 1/(k_{y,\min} \text{ shear})$

ITPA geometry:

$$k_{y,\min} \rho_s = 0.022$$

$$\text{shear} = 0.0457$$

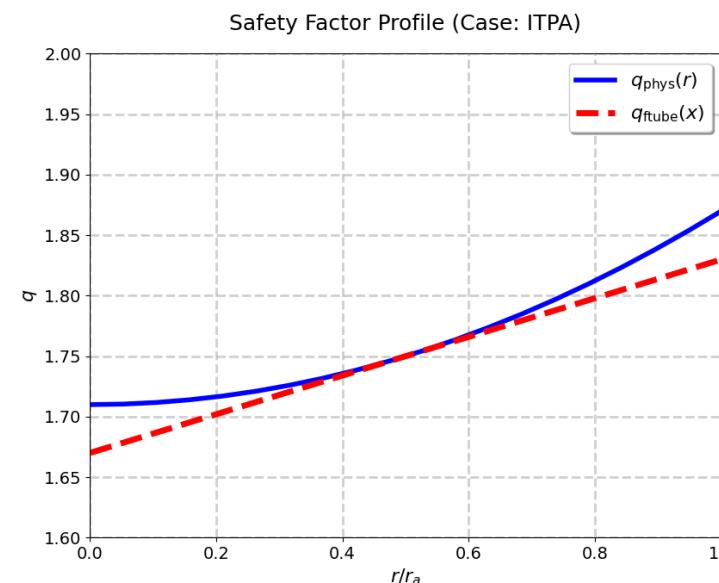
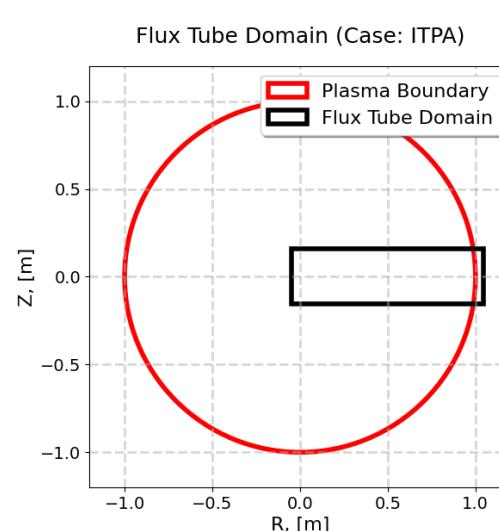
$$L_x = 507.9 \rho_s, N_{\text{pol}} = 2$$

$$L_y = 291.7 \rho_s$$

$$r_a / \rho_s = 927.0$$



ITPA TAE benchmark: flux-tube GENE



ITPA geometry:

$$k_{y,\min} \rho_s = 0.022$$

$$\text{shear} = 0.0457$$

$$L_x = 1015.8 \rho_s \quad (N_{\text{pol}} = 1)$$

$$L_y = 291.7 \rho_s$$

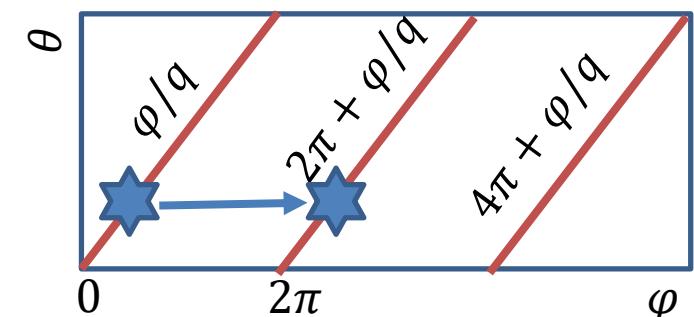
$$r_a / \rho_s = 927.0$$

In the flux-tube approach, one replaces a torus with a box

Poloidal box size is determined by spectral resolution:

$$L_y = 2\pi / k_{y,\min} \quad , \quad k_{y,\min} \rho_s = m_0 \rho_*$$

Radial box size is determined by „twist-and-shift“ condition (parallel bc): $L_x = 1/(k_{y,\min} \text{ shear} N_{\text{pol}})$



Twist-and-shift parallel boundary condition:

Shift: $F(x, \alpha, \theta + 2\pi N_{\text{pol}}) = F(x, \alpha - 2\pi q N_{\text{pol}}, \theta) ; \alpha = q(x)\theta - \phi ; k_y r_0 = m_0$

$$\exp(i k'_x x) \exp \left[i k_y \left(y - \frac{2\pi q(x) r_0}{q_0} N_{\text{pol}} \right) \right] \sim \exp(-2\pi i N_{\text{pol}} k_y r_0) \exp(i k'_x - 2\pi i N_{\text{pol}} k_y \hat{s} x)$$

Twist: $k'_x = k_x + 2\pi k_y \hat{s} N_{\text{pol}} \Rightarrow k_{x,\min} = 2\pi N_{\text{pol}} k_{y,\min} \Rightarrow L_x = 1/(k_{y,\min} \hat{s} N_{\text{pol}})$



In the flux-tube approach, one replaces a torus with a box

Poloidal box size is determined by spectral resolution: $L_y = 2 \pi / k_{y,\min}$, $k_{y,\min} \rho_s = m_0 \rho_*$

Radial box size is determined by „twist-and-shift“ condition (parallel bc): $L_x = 1/(k_{y,\min} \text{ shear } N_{\text{pol}})$

ITPA geometry:

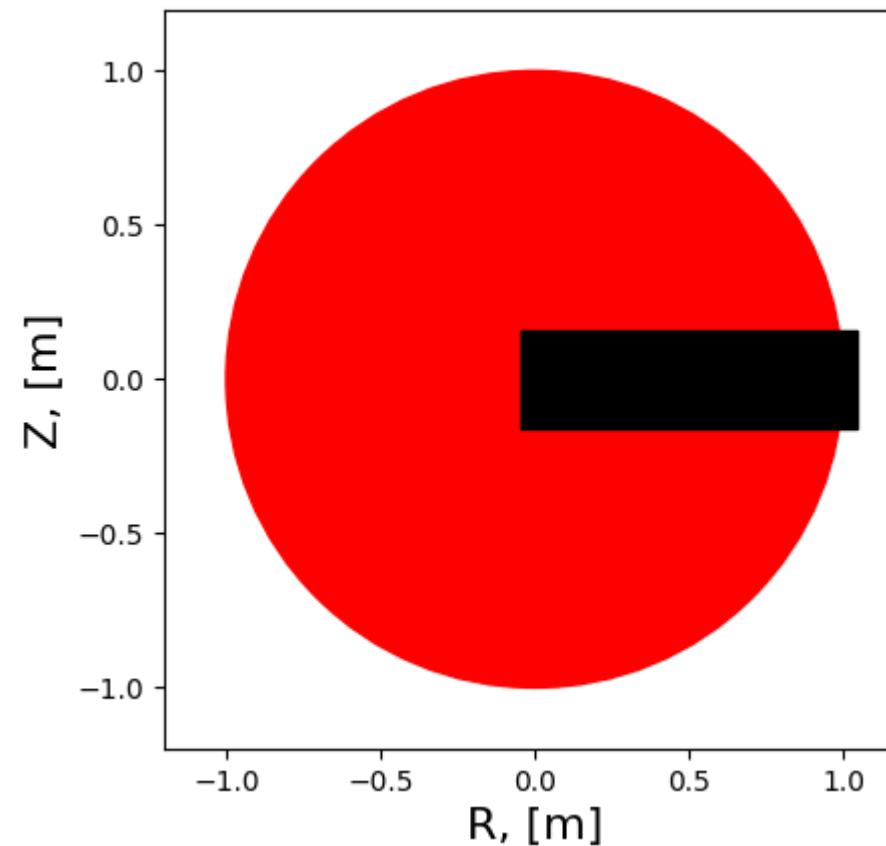
$$k_{y,\min} \rho_s = 0.022$$

$$\text{shear} = 0.0457$$

$$L_x = 1015.8 \rho_s (N_{\text{pol}} = 1)$$

$$L_y = 291.7 \rho_s$$

$$r_a / \rho_s = 927.0$$



Only shear (twist-and-shift bc) and density gradient computed at r_0 for radial (perpendicular) dependencies

Parallel variation same for all „x“

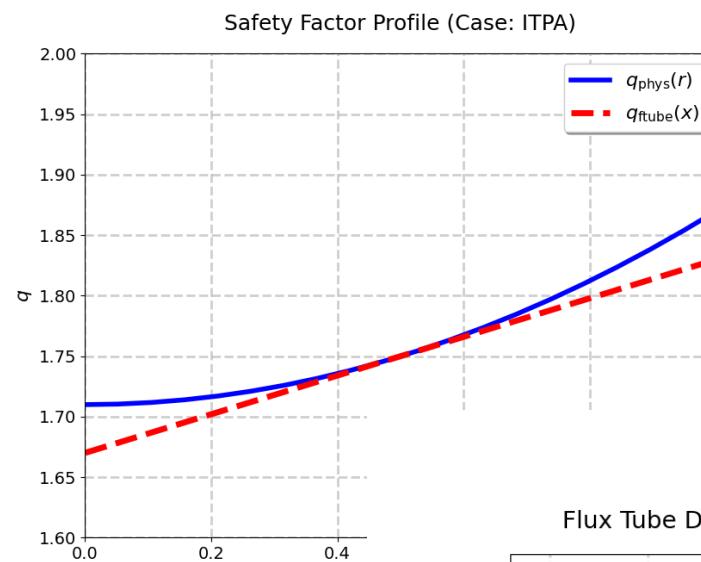
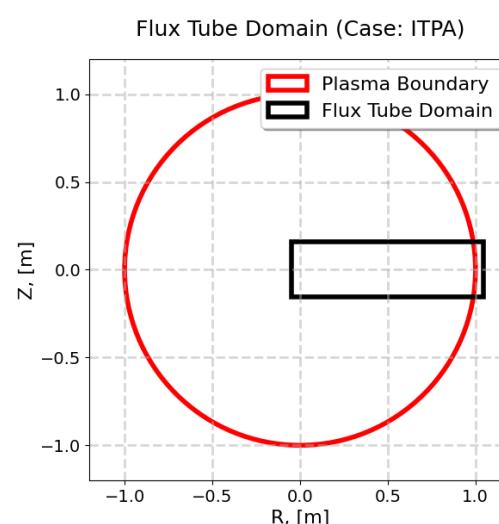
Minimalistic „adapted“ L_x for linear run (nonlinear L_x can be much larger)

EP orbits are „smaller“ than the flux tube

The flux tube is not „small“

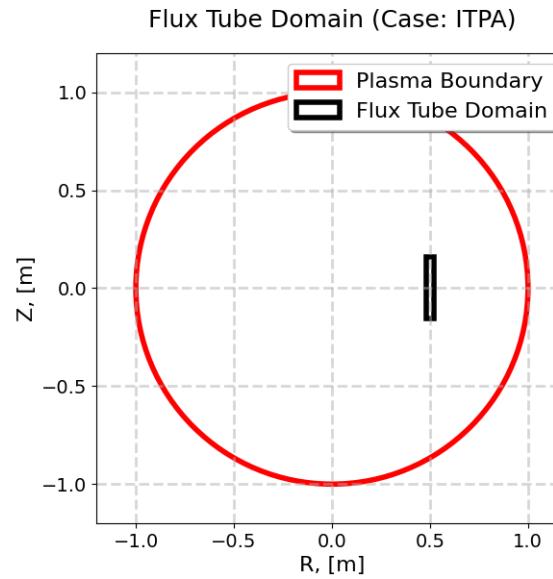


ITPA TAE benchmark: flux-tube GENE

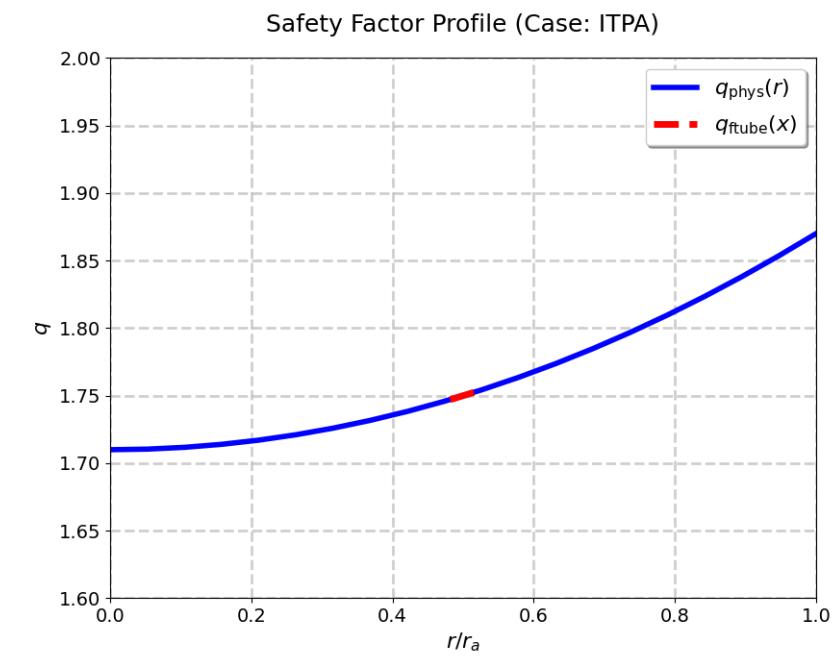


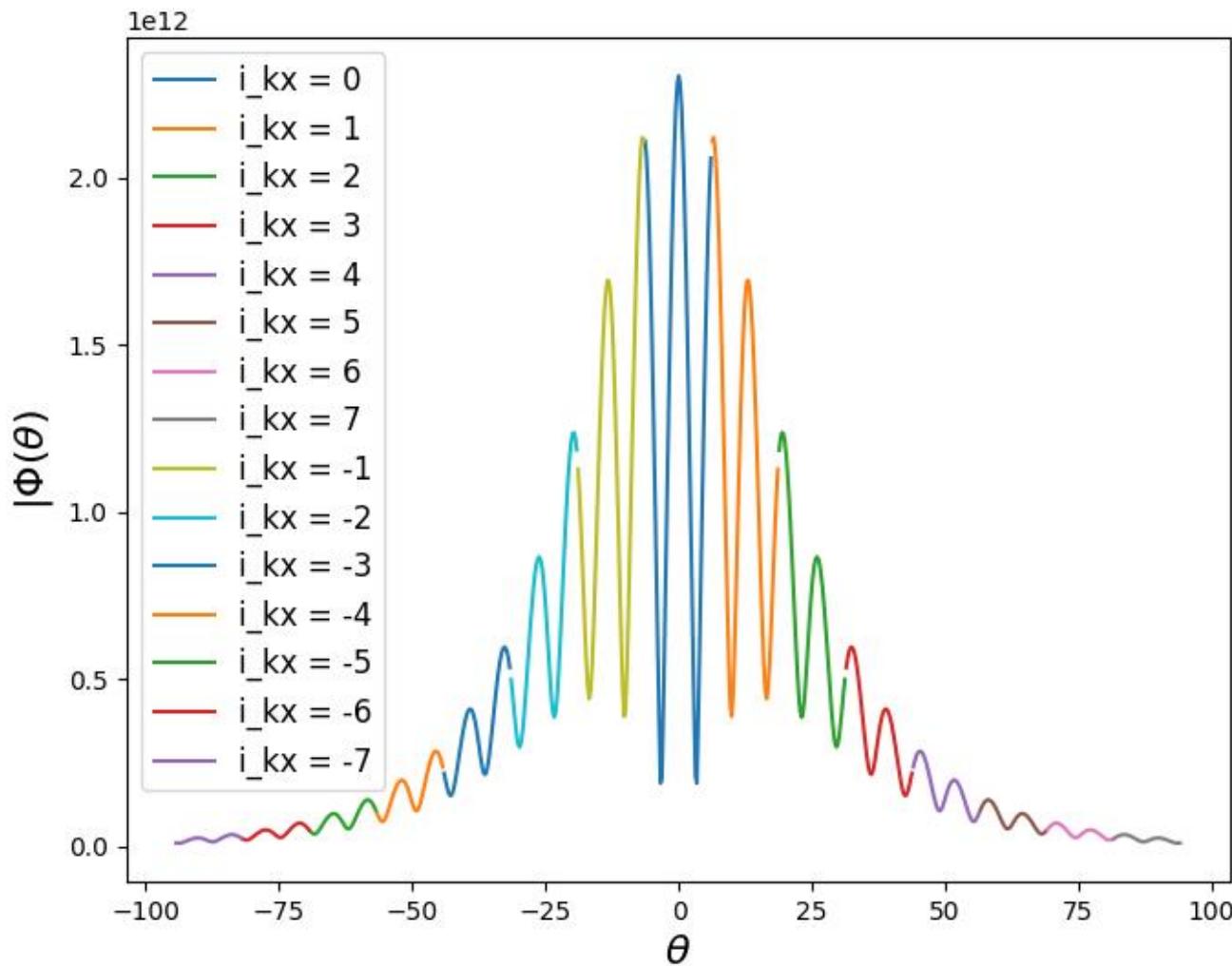
ITPA geometry:
 $k_{y,\min} \rho_s = 0.022$
shear = 0.0457
 $L_x = 1015.8 \rho_s$ ($N_{\text{pol}} = 1$)
 $L_y = 291.7 \rho_s$
 $r_a / \rho_s = 927.0$

ITPA geometry:
 $k_{y,\min} \rho_s = 0.022$
shear = 0.0457
 $L_x = 33.8 \rho_s$ ($N_{\text{pol}} = 30$)
 $L_y = 291.7 \rho_s$
 $r_a / \rho_s = 927.0$



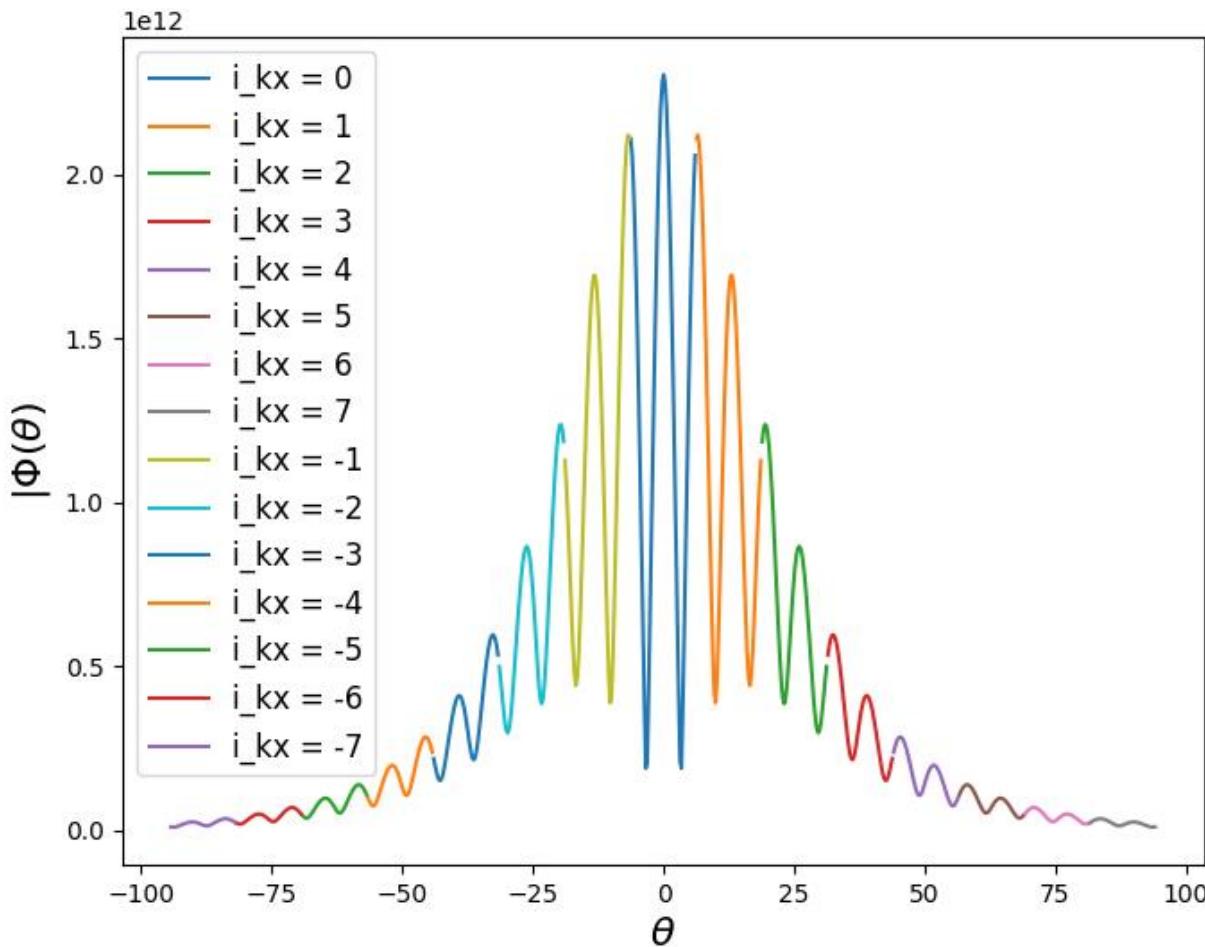
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 $L_y = 2 \pi / k_{y,\min}$, $k_{y,\min} \rho_s = m_0 \rho_*$
Radial box size is determined by „twist-and-shift“ condition (parallel bc): $L_x = 1 / (k_{y,\min} \text{ shear } N_{\text{pol}})$





TAE instability (physical gap mode)

Around **31 poloidal turns** are needed to resolve the mode: $-31\pi < \vartheta < -31\pi$



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TAE mode results from a specific interaction of two shear Alfvén waves

Multiple-scale parallel mode structure

The width of the continuum mode has to be resolved (implying high k_x or, alternatively, large L_x)

For ideal MHD, SAW width is a delta function resulting in infinite domain in the ballooning space

Kinetically, gyro-radius introduces a finite scale $\sim \rho_s$



Extraction of the poloidal mode structure:

$$x = q_0/(B_0 r_0)(\psi - \psi_0) = r - r_0 ; y = (r_0/q_0)[q(x)\theta - \zeta] ; z =$$

$$q = q_0(1 + \hat{s}x/r_0) ; \hat{s} = (r_0/q_0)dq/dr$$

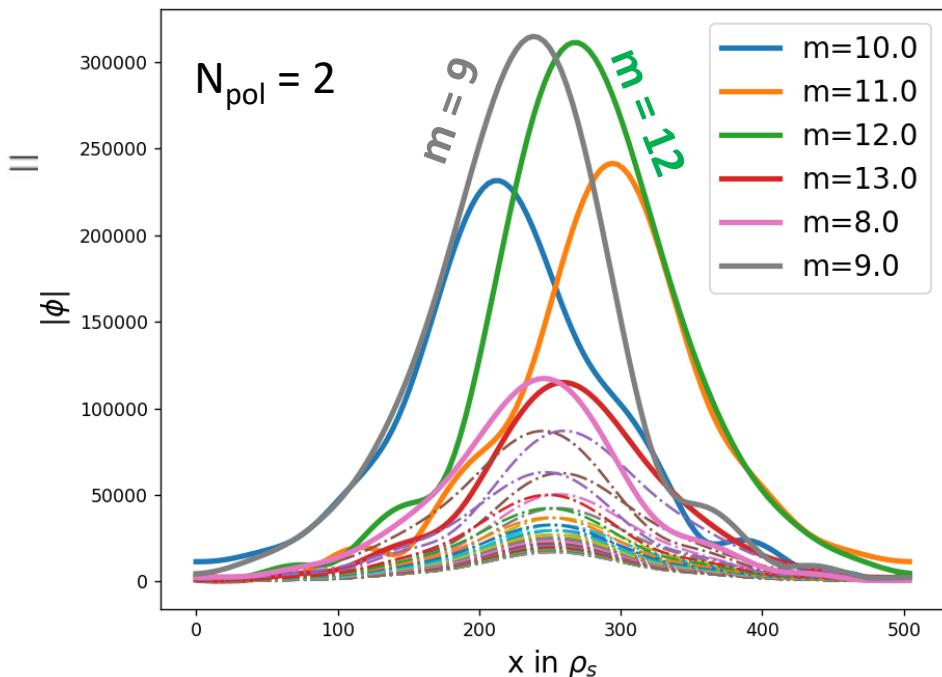
$$\phi_n = \Phi_n(x, z) \exp(ik_y r_0 \theta) \exp(ik_y x \hat{s} \theta) \exp[-ik_y (r_0/q_0) \zeta]$$

$$n_0 = k_y r_0 / q_0 ; m_0 = k_y r_0 = k_y \rho_s / \rho_* ; \rho_* = \rho_s / r_0$$

$$\phi_n = \sum_{m'} \left[\sum_{k_x} \hat{\Phi}_{k_x m' n} \exp(ik_x x) \right] e^{im' \theta} \exp(ik_y x \hat{s} \theta) \exp(im_0 \theta) \exp(-in_0 \zeta), \theta \in [-N_{pol} \pi, N_{pol} \pi)$$

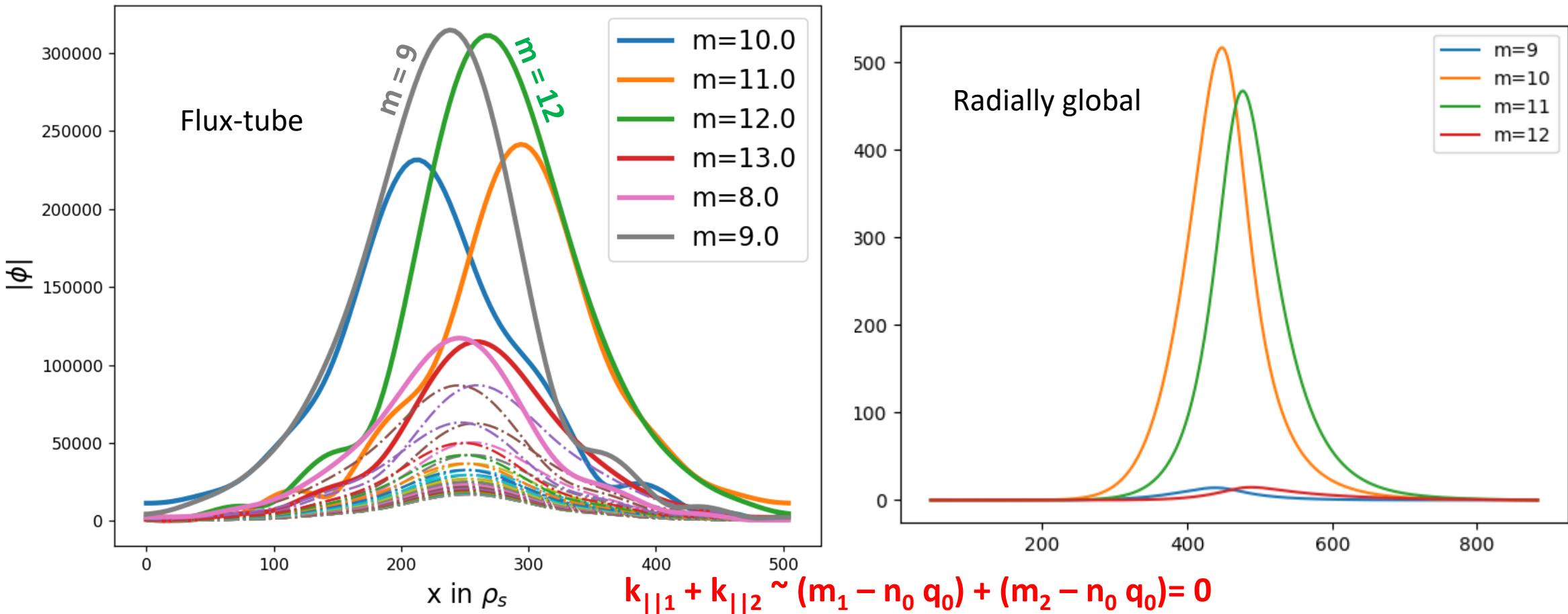
$$\mathbf{k}_{||1} + \mathbf{k}_{||2} \sim (m_1 - n_0 q_0) + (m_2 - n_0 q_0) = 0$$

- Mode structure of coupled poloidal harmonics (a “standing wave”) is reproduced:
 $m_1 + m_2 = 2 n_0 q(s_0) = 21$ for $n_0 = 6$ and $q(s_0) = 1.75$
- **In contrast to global result, harmonics with $m_2 - m_1 > 1$ are present and strong!**
- The radial structure of the mode is different : periodic BC in radial direction (box „multiplied“)



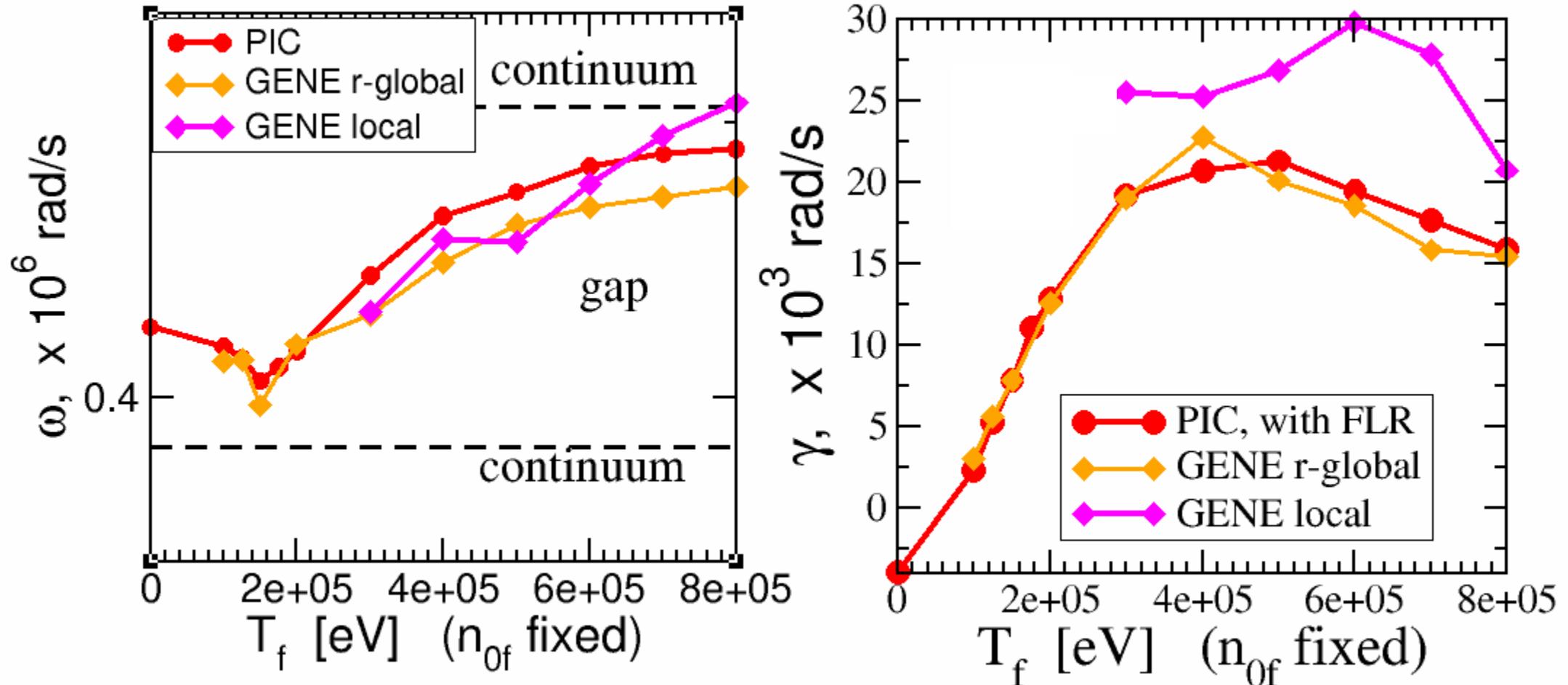
Toroidal mode number:
 $n_0 = k_y r_0 / q_0$

Poloidal mode number:
 $m_0 = k_y r_0$



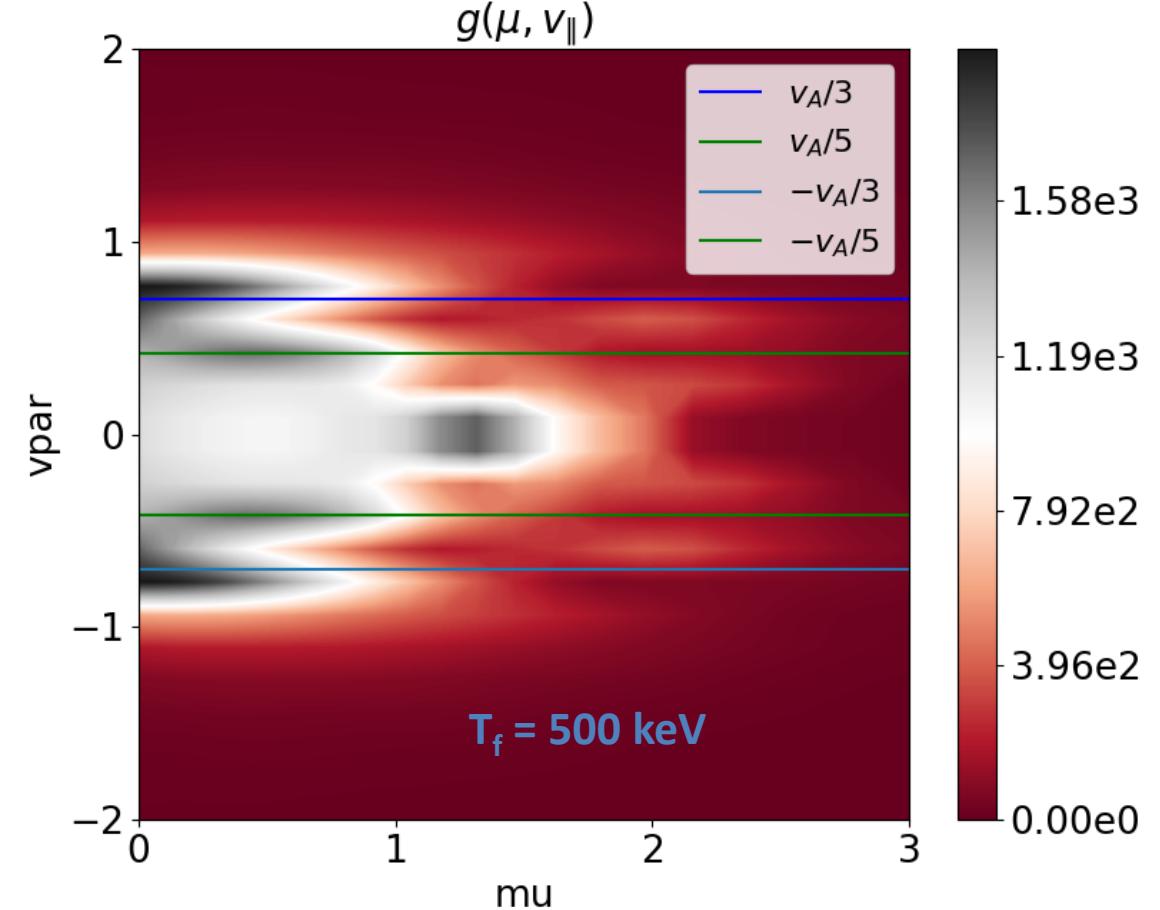
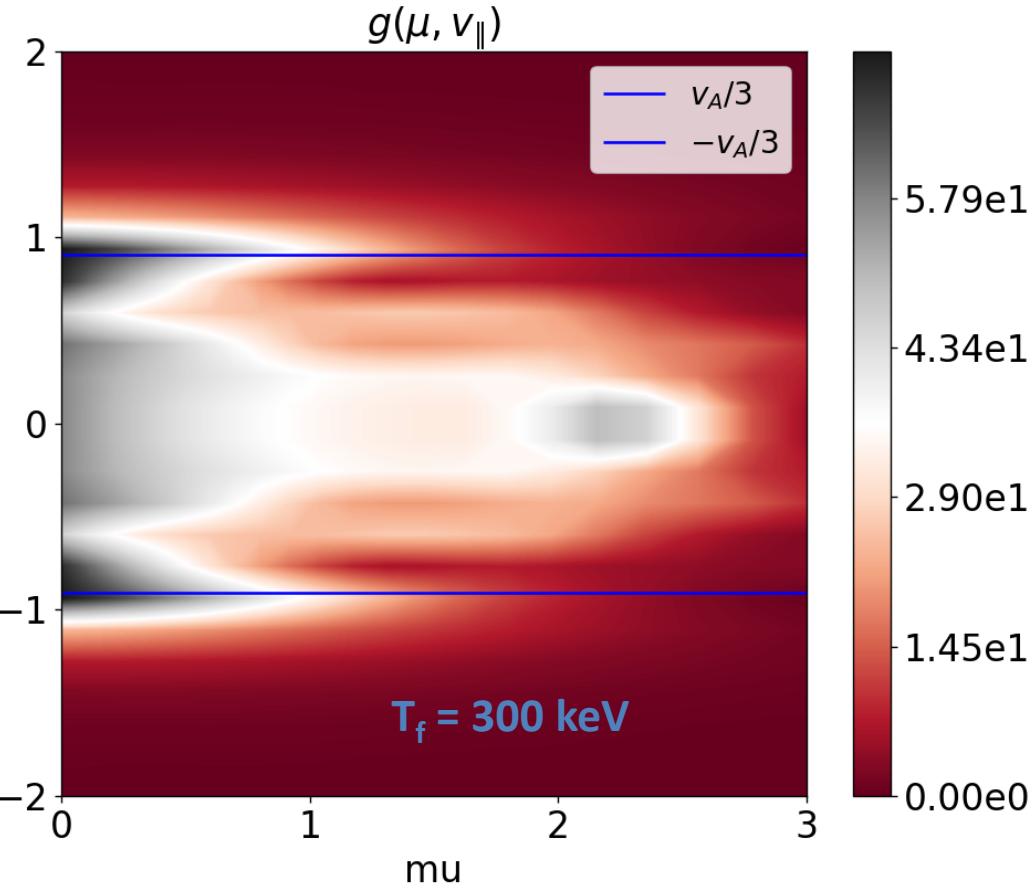
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Toroidal mode number:
 $n_0 = k_y r_0 / q_0$
 Poloidal mode number:
 $m_0 = k_y r_0$



Good agreement in the TAE frequency; comparable numbers in the growth rate

Parallel mode structure at the accumulation point is the key for the mode; included (to some extent) in flux tube



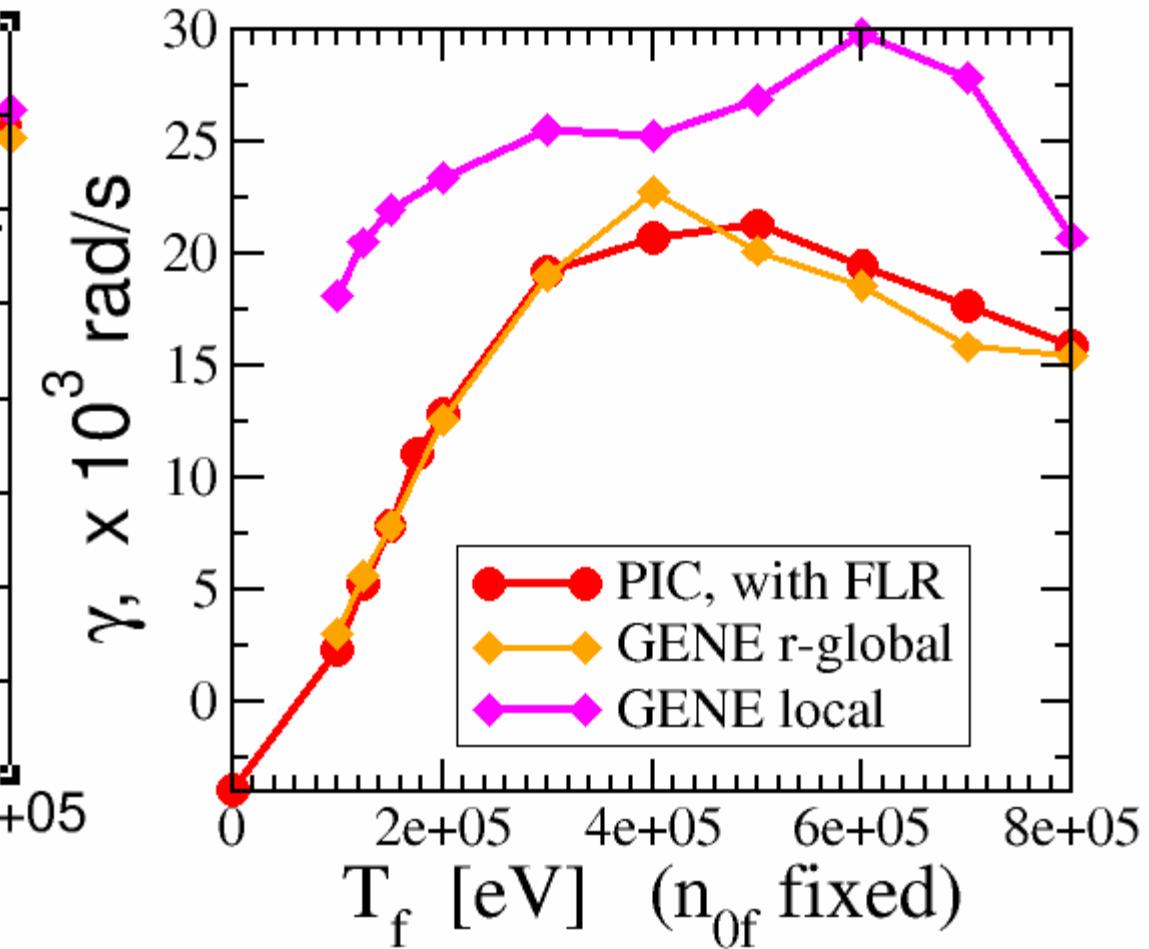
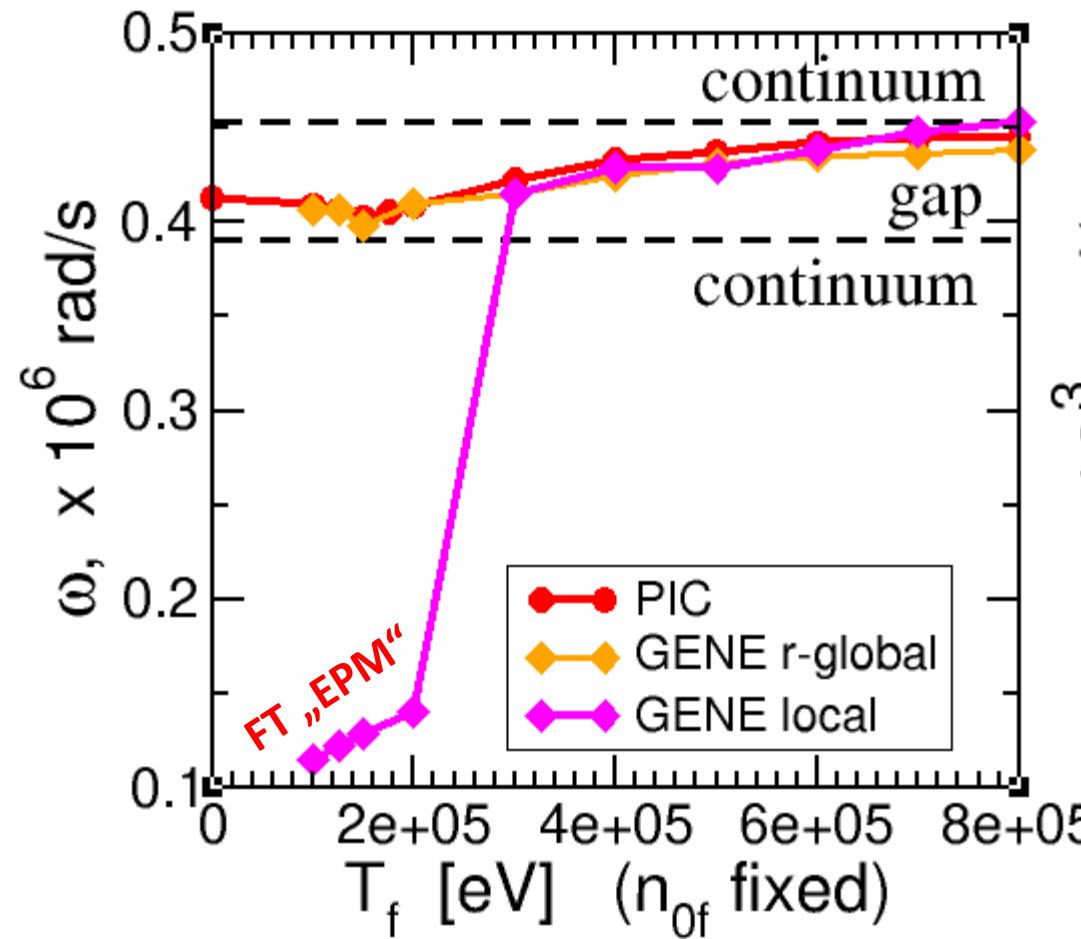
Energetic-particle velocity space: non-adiabatic distribution function

Resonances with energetic passing particles: the main resonance v_A and the „side-band“ resonances

$v_A/(2 s - 1)$: $v_A/3$, $v_A/5$ (?), $v_A/7$ (?) ...

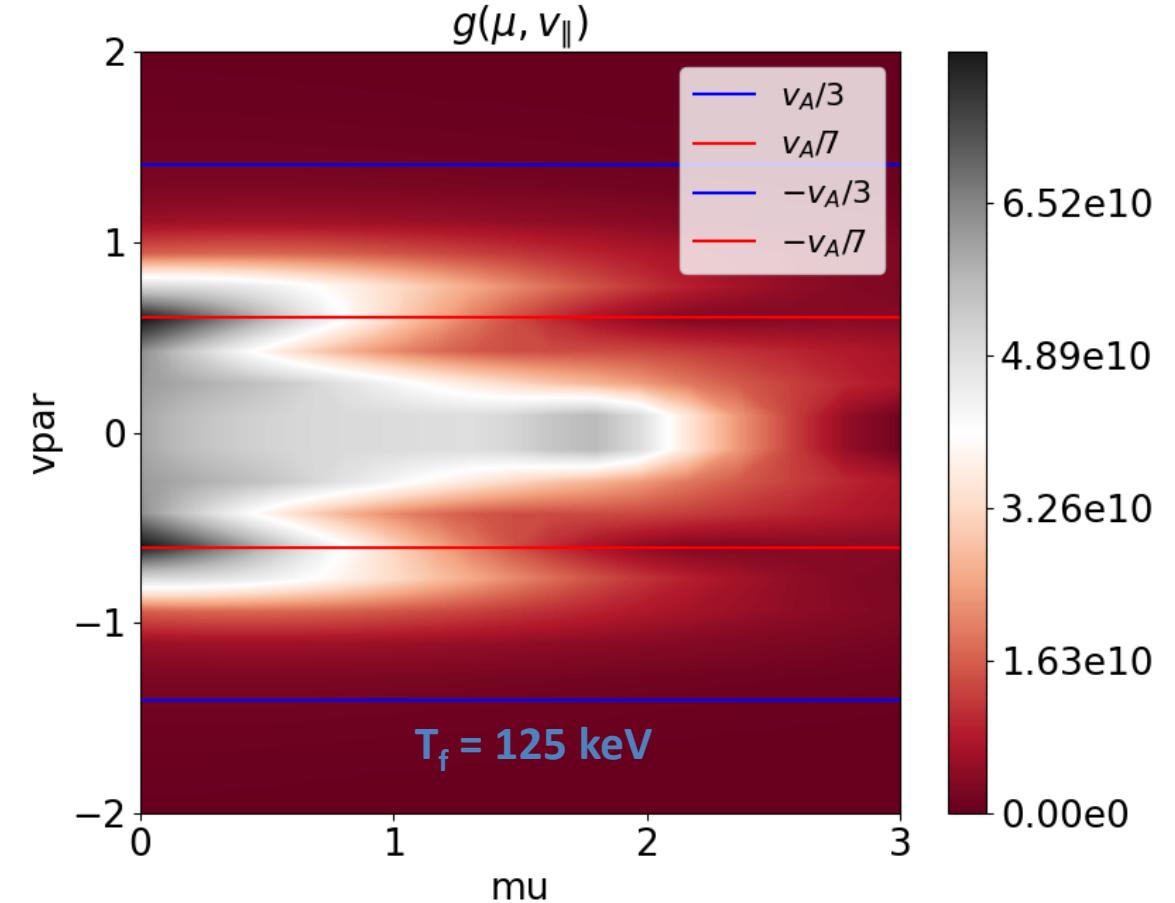
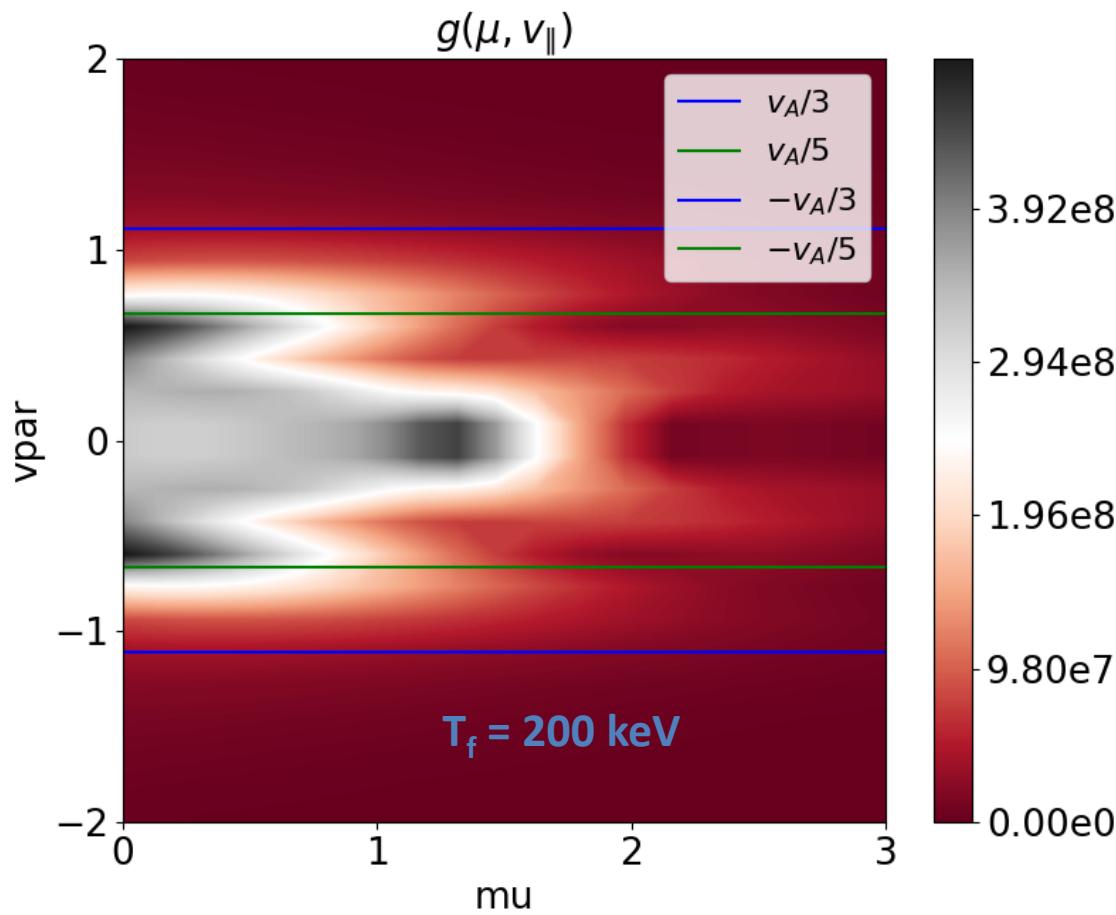
In global simulations: only $v_A/3$ or v_A are usually relevant

Resonance condition:
 $k_z v_A = k_{z1} v_{EP}$, $k_{z1} \neq k_z$
 (generally)



For smaller temperatures: jump to an „EPM“ in the flux-tube case (deep into the continuum)
 The „flux-tube EPM“ appears to be well resolved (codewise). Is it real?

ITPA TAE benchmark: flux-tube GENE – radially global GENE comparison

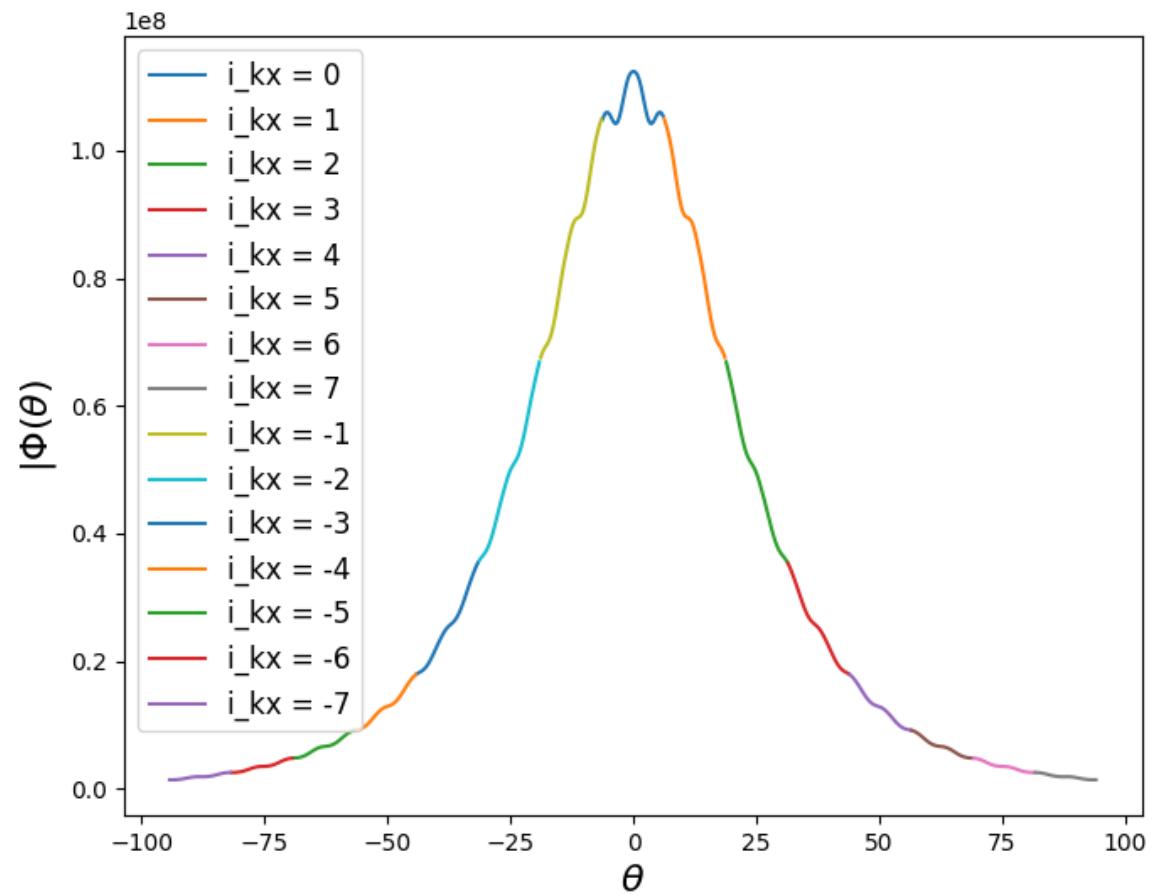
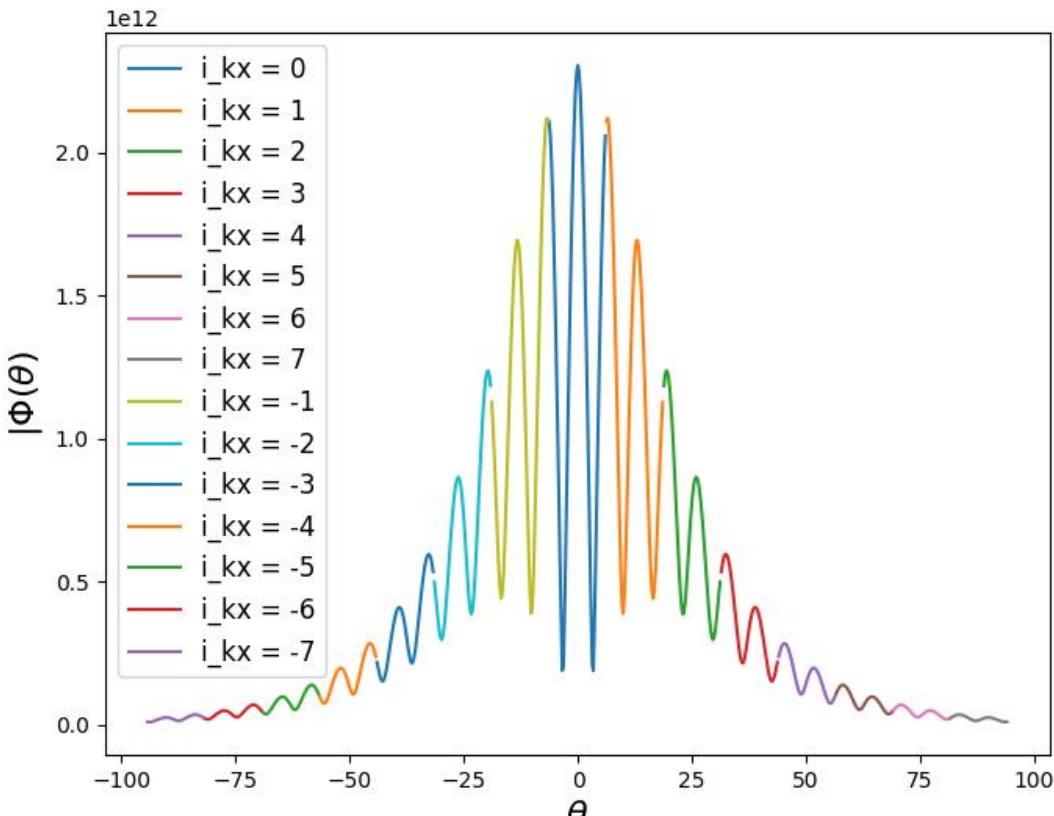


Energetic-particle velocity space: non-adiabatic distribution function

Higher-order „side-band“ resonances are dominant: $v_A/5, v_A/7$; lower frequencies better for drive
In global simulations: only $v_A/3$ is usually relevant ; frequency remains in the TAE gap

Resonance condition:
 $k_z v_A = k_{z1} v_{EP}$, $k_{z1} \neq k_z$
 (generally)

Additional „FT“ resonance because of the FT distortions in the mode structure?



TAE instability (physical gap mode)
 More than **31 poloidal turns** are needed to
 resolve the mode: $-31\pi < \vartheta < -31\pi$

Flux-tube „EPM“. Is it physical?
**Parallel and radial mode structures
 result from the same dataset**



Conclusions (not yet final)

- Reasonable agreement in frequency for higher Tf; similar order of growth rates
- Mode structure is very broad in parallel direction; multiple scales (fine-scale component + envelope)
- The radial mode structure is different (additional couplings): may be a consequence of the periodic boundary conditions (radially) effectively „multiplying“ the box in the radial direction
- Side-band resonances appear ($v_A/5, v_A/7$) not so present globally: may be a consequence of the issues with the parallel FT mode structure (failure of the „outgoing wave“ boundary condition; reflections)
- These additional resonances drive an „EPM“ at smaller $T_f < 300$ keV