

# Flux-tube simulations of TAE modes (work in progress)

A. Mishchenko

Acknowledgements: D. Brioschi, T. Görler, T. Hayward-Schneider, R. Kleiber, A. Könies, A. Di Siena

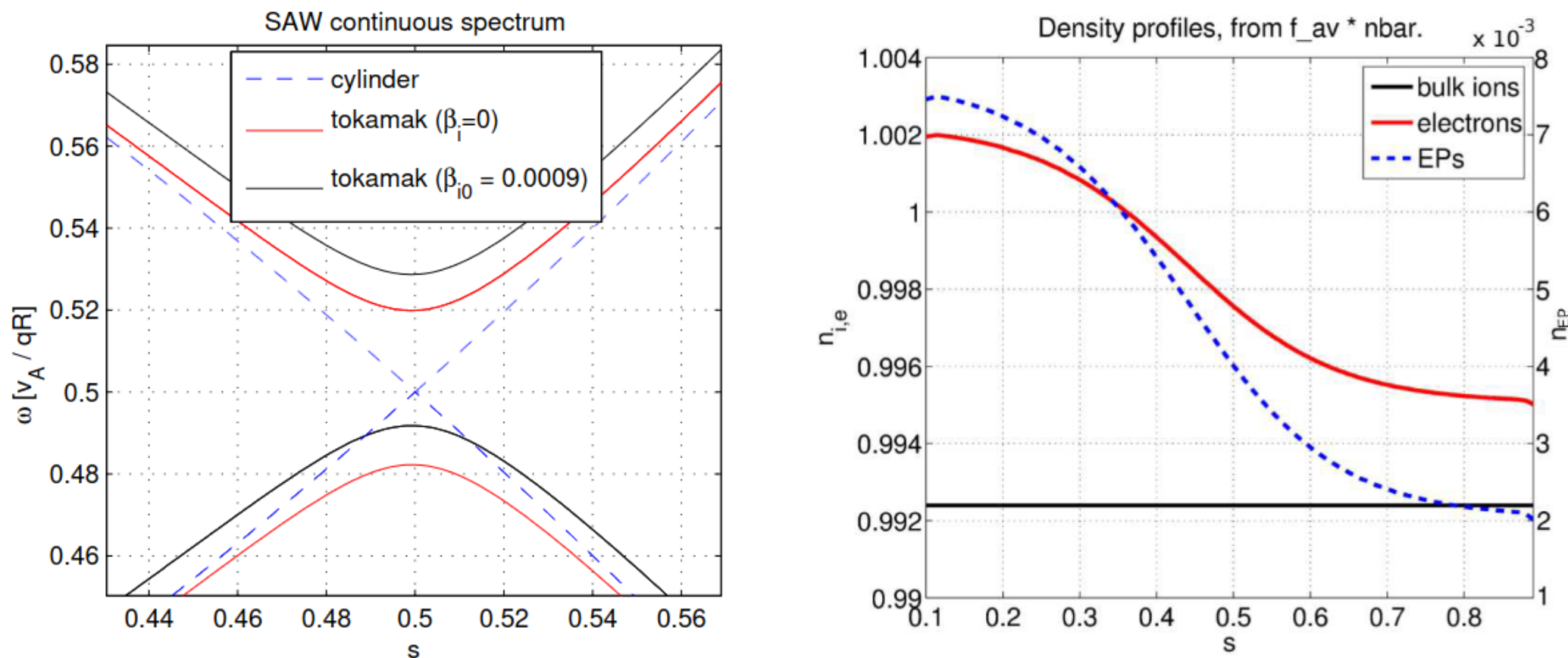
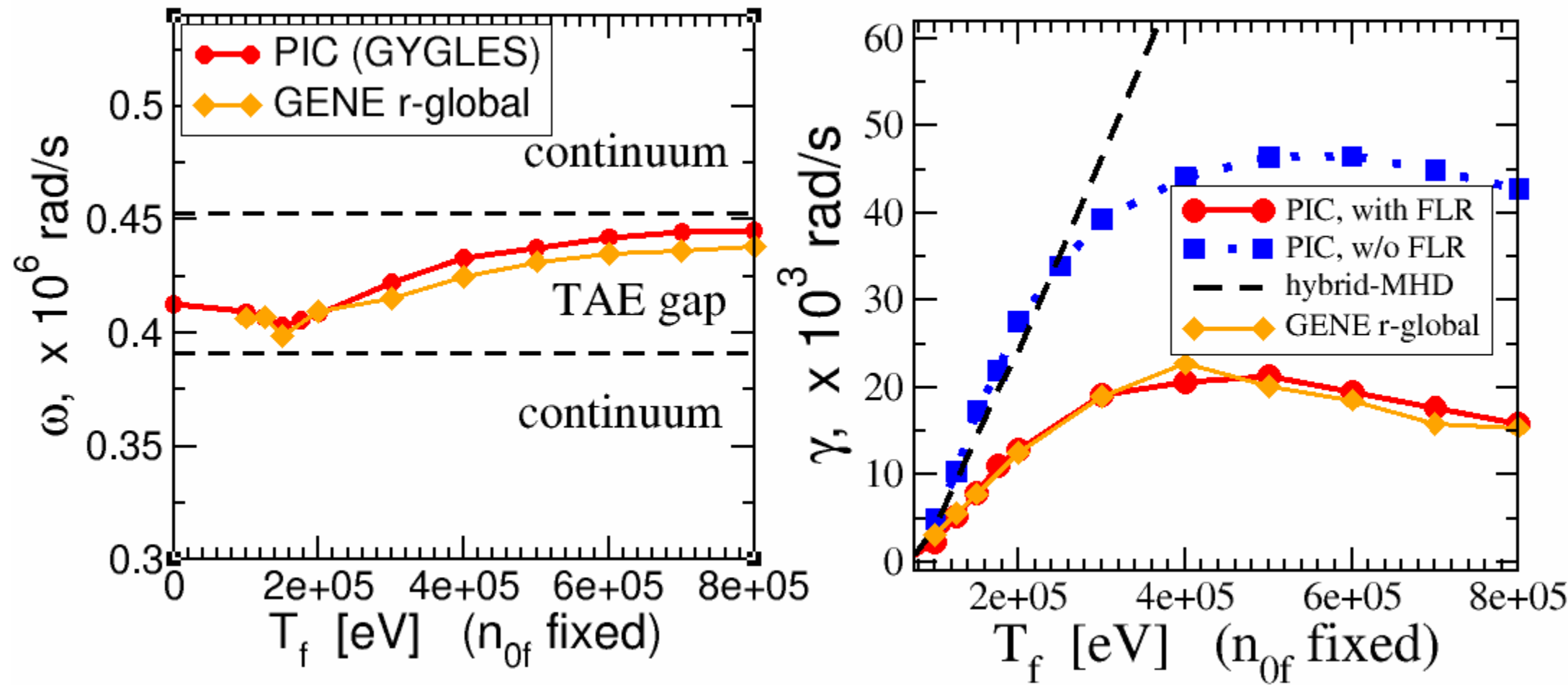
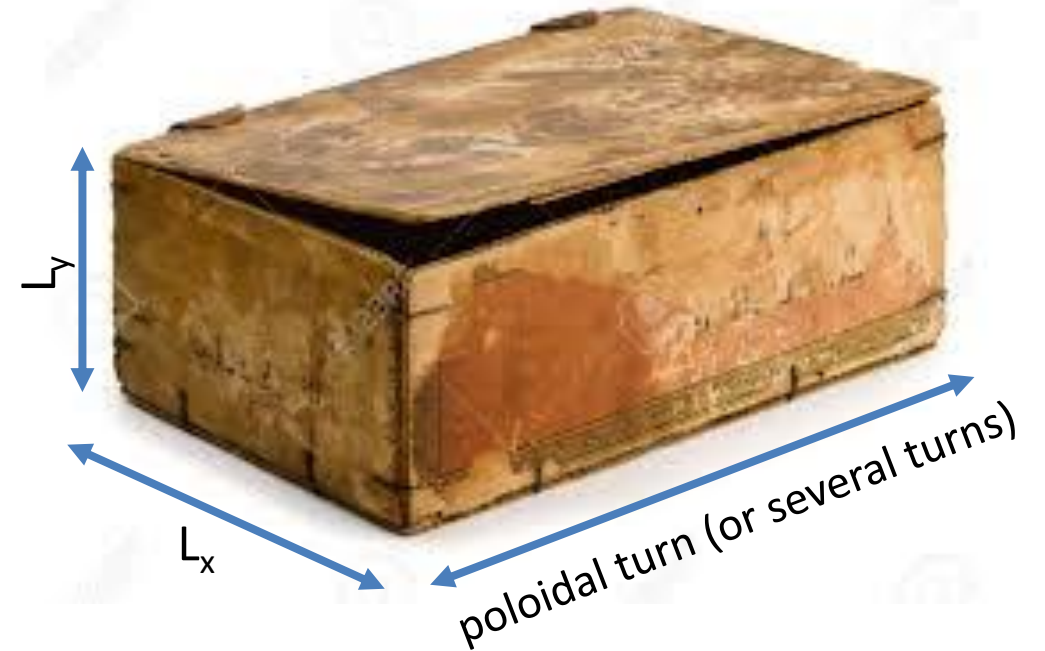


Figure 1: On the left, zoom of the continuous spectrum obtained by neglecting toroidicity and compressibility (blue dashed lines), or by keeping toroidicity only (red line) or by keeping both toroidicity and compressibility (black line), with approximated formula of Ref. [21]. On the right, density profiles in units of  $n_{e0} = 2 \cdot 10^{19} m^{-3}$ , for a case with  $\langle n_{EP} \rangle / n_{e0} = 0.004$ .



Radially-global GENE linear frequency and growth rate compare very well to GYGLES (even maybe slightly better than Davide’s comparison with ORB5 performed last year).

GYGLES simulations: [Phys. Plasmas 16, 082105 \(2009\)](#) ; ITPA benchmark: [A. Könies et al 2018 Nucl. Fusion 58 126027](#)



In the flux-tube approach, one replaces a torus with a box

Poloidal box size is determined by spectral resolution:  $L_y = 2 \pi / k_{y,\min}$

Radial box size is determined by „twist-and-shift“ condition (parallel bc):  $L_x = 1/(k_{y,\min} \text{ shear})$

ITPA geometry:

$k_{y,\min} \rho_s = 0.022$

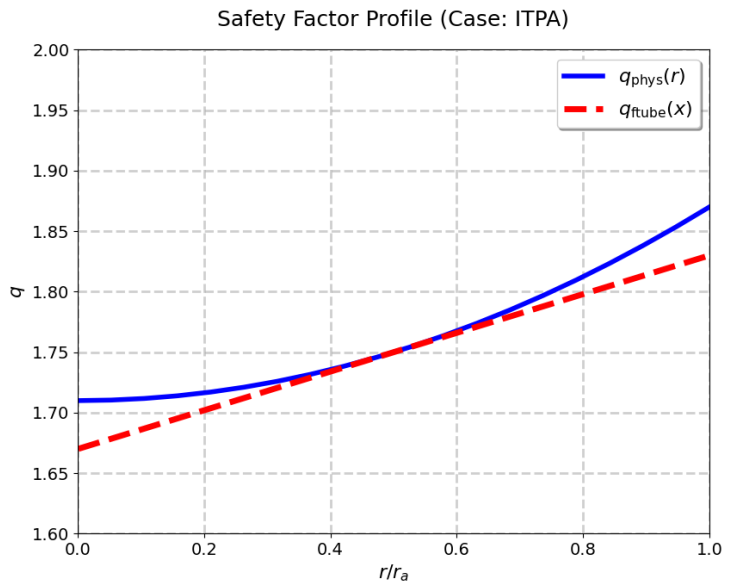
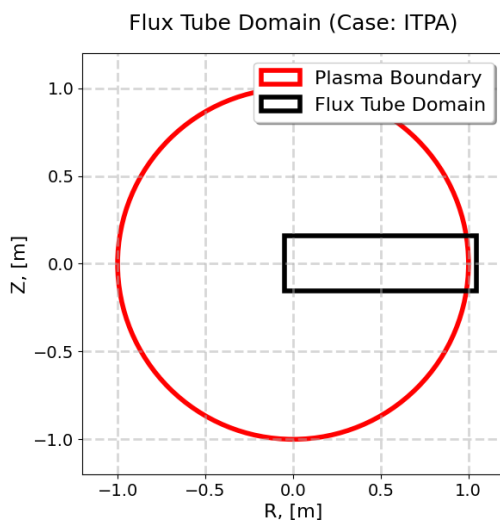
shear = 0.0457

$L_x = 507.9 \rho_s$ ,  $N_{\text{pol}} = 2$

$L_y = 291.7 \rho_s$

$r_a / \rho_s = 927.0$



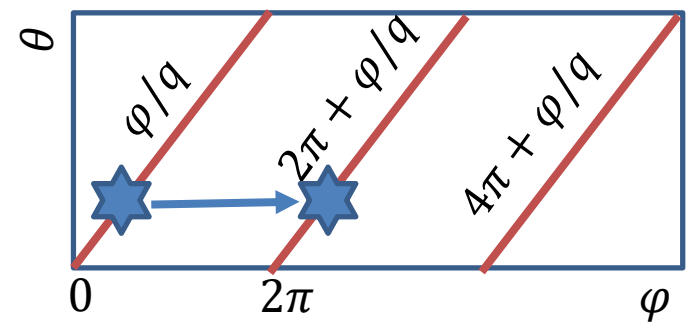


In the flux-tube approach, one replaces a torus with a box

Poloidal box size is determined by spectral resolution:

$$L_y = 2 \pi / k_{y,\min} \quad , \quad k_{y,\min} \rho_s = m_0 \rho_*$$

Radial box size is determined by „twist-and-shift“ condition (parallel bc):  $L_x = 1/(k_{y,\min} \text{ shear } N_{\text{pol}})$



*Twist-and-shift parallel boundary condition:*

ITPA geometry:

$$k_{y,\min} \rho_s = 0.022$$

$$\text{shear} = 0.0457$$

$$L_x = 1015.8 \rho_s \quad (N_{\text{pol}} = 1)$$

$$L_y = 291.7 \rho_s$$

$$r_a / \rho_s = 927.0$$

Shift:  $F(x, \alpha, \theta + 2\pi N_{\text{pol}}) = F(x, \alpha - 2\pi q N_{\text{pol}}, \theta) ; \alpha = q(x)\theta - \varphi ; k_y r_0 = m_0$

$$\exp(i k'_x x) \exp \left[ i k_y \left( y - \frac{2\pi q(x) r_0}{q_0} N_{\text{pol}} \right) \right] \sim \exp(-2\pi i N_{\text{pol}} k_y r_0) \exp(i k'_x - 2\pi i N_{\text{pol}} k_y \hat{s} x)$$

Twist:  $k'_x = k_x + 2\pi k_y \hat{s} N_{\text{pol}} \Rightarrow k_{x,\min} = 2\pi N_{\text{pol}} k_{y,\min} \Rightarrow L_x = 1/(k_{y,\min} \hat{s} N_{\text{pol}})$



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ITPA geometry:

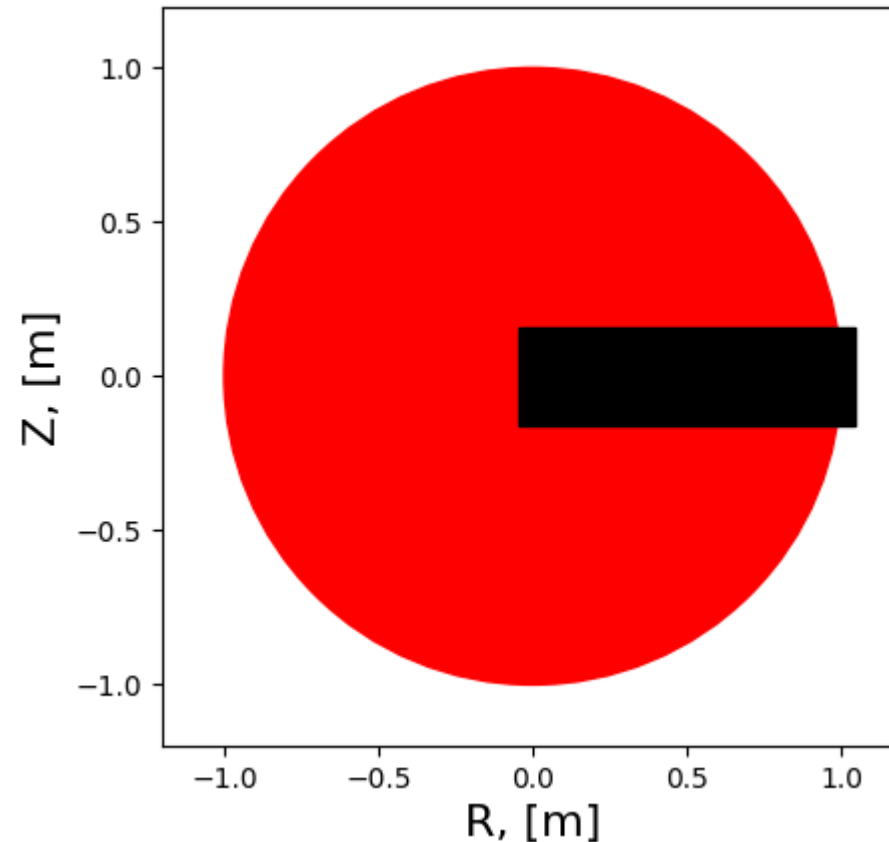
$$k_{y,\min} \rho_s = 0.022$$

$$\text{shear} = 0.0457$$

$$L_x = 1015.8 \rho_s (N_{\text{pol}} = 1)$$

$$L_y = 291.7 \rho_s$$

$$r_a / \rho_s = 927.0$$



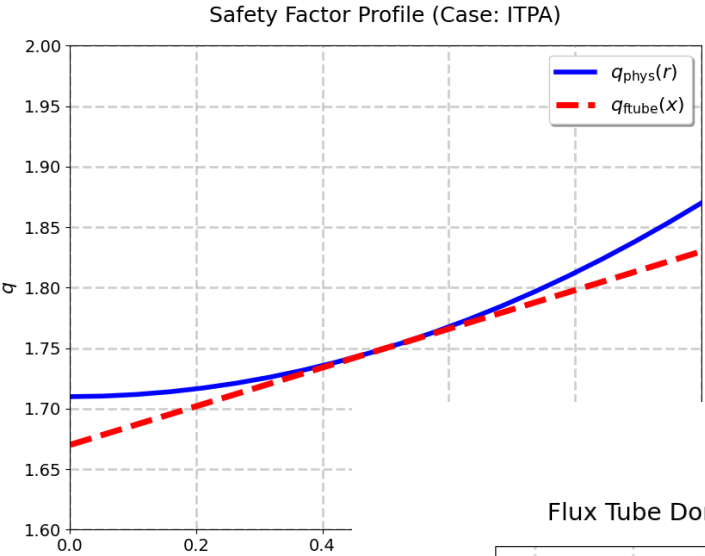
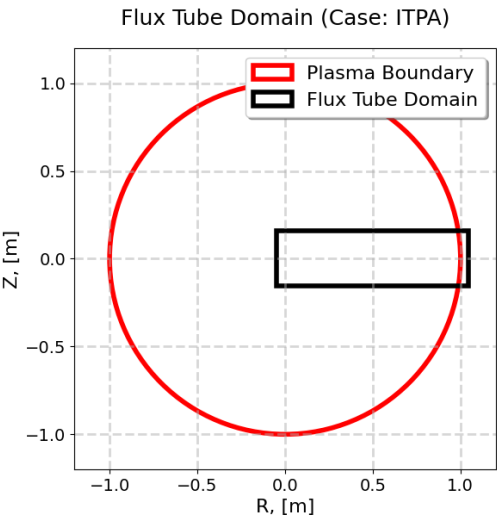
Only shear (twist-and-shift bc) and density gradient computed at  $r_0$  for radial (perpendicular) dependencies

Parallel variation same for all „x“

Minimalistic „adapted“  $L_x$  for linear run (nonlinear  $L_x$  can be much larger)

EP orbits are „smaller“ than the flux tube

The flux tube is not „small“



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ITPA geometry:

$$k_{y,\min} \rho_s = 0.022$$

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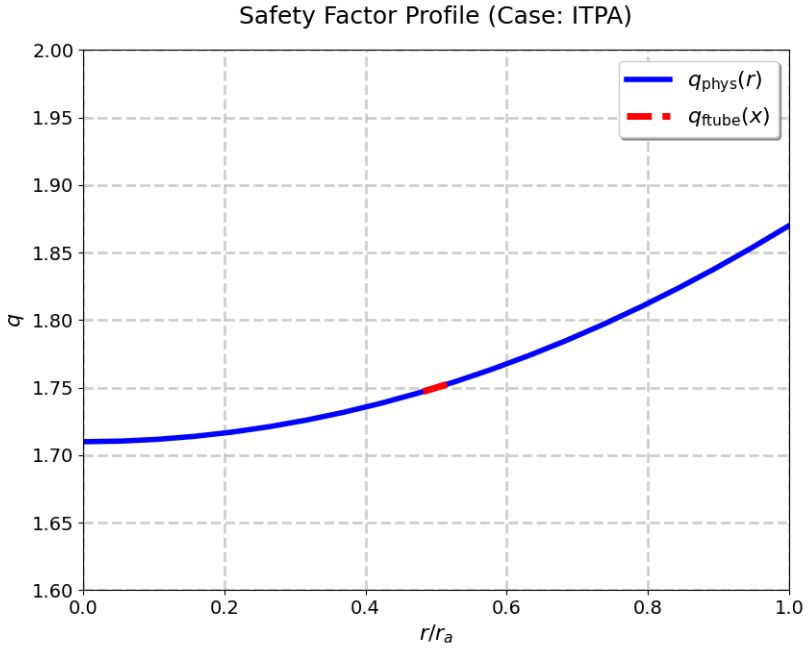
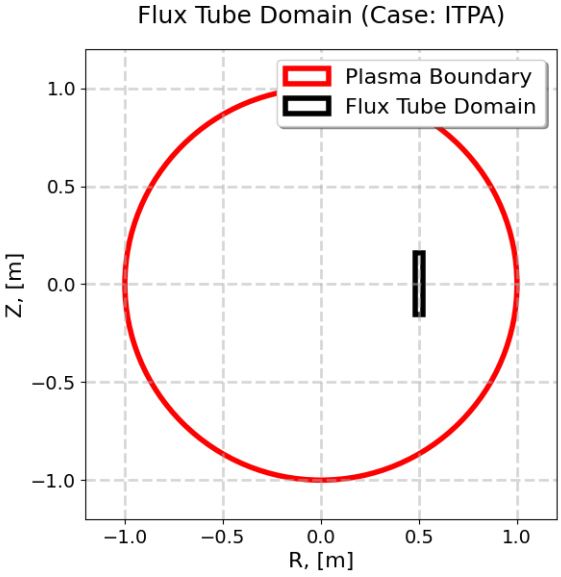
$$k_{y,\min} \rho_s = 0.022$$

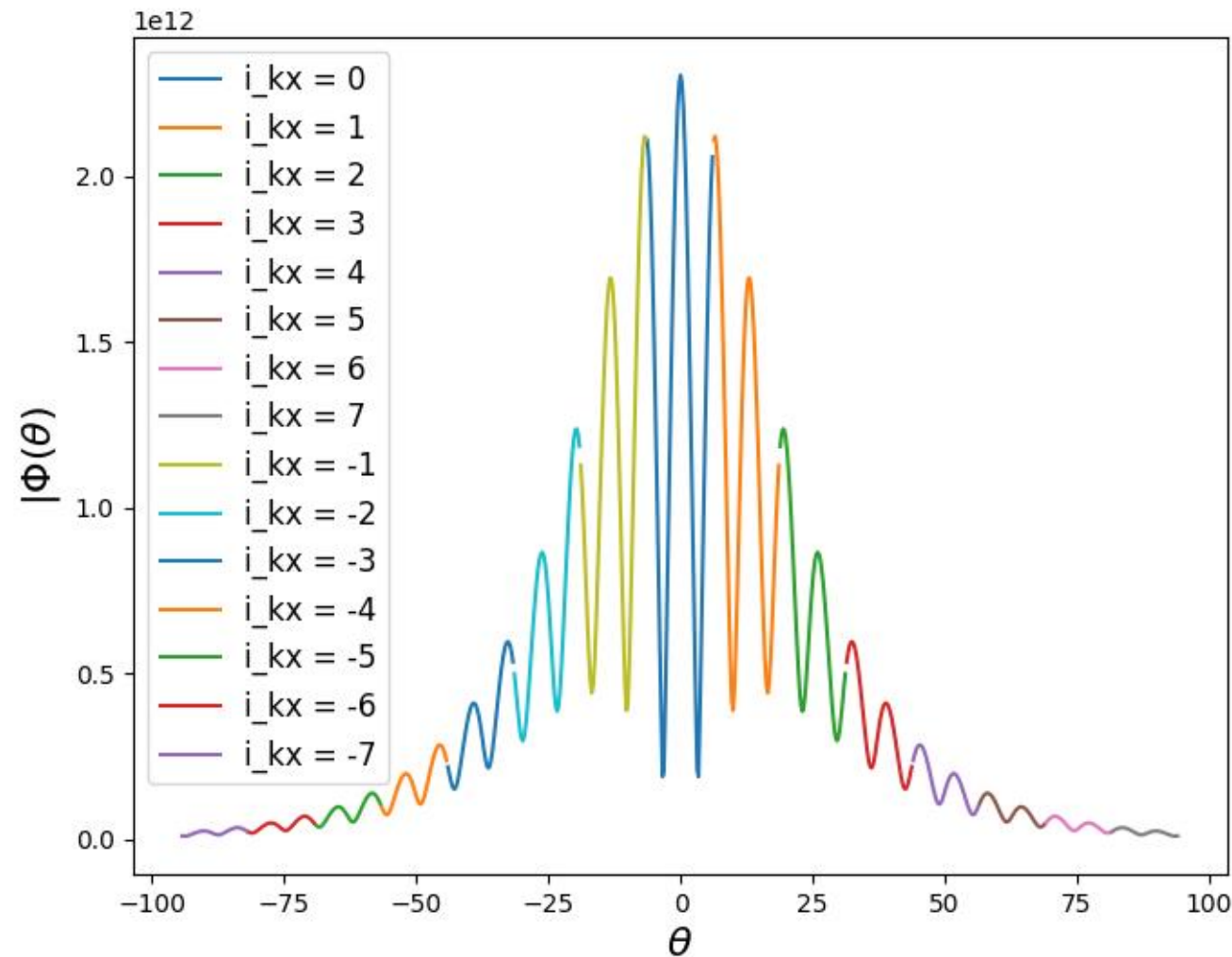
$$\text{shear} = 0.0457$$

$$L_x = 33.8 \rho_s \quad (N_{\text{pol}} = 30)$$

$$L_y = 291.7 \rho_s$$

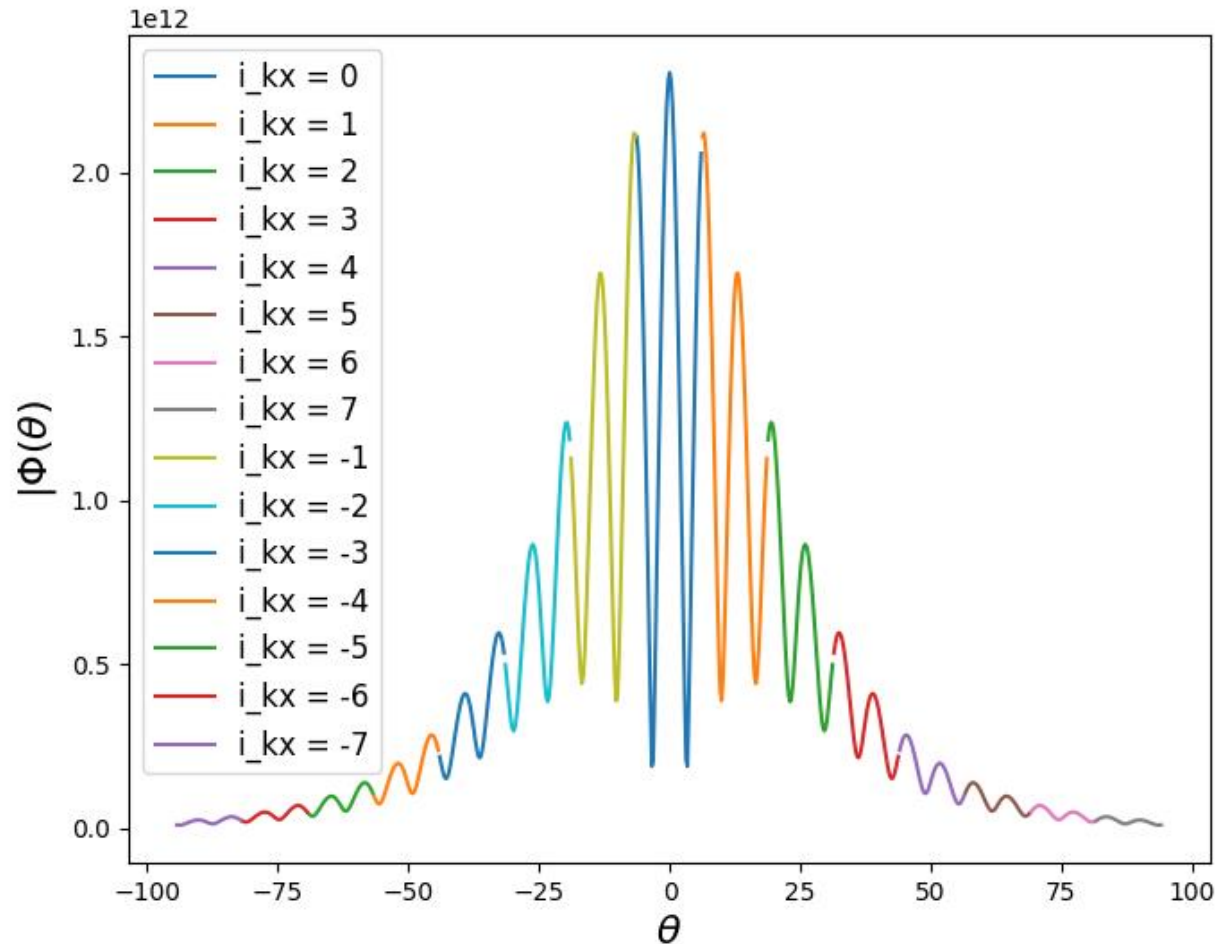
$$r_a / \rho_s = 927.0$$





TAE instability (physical gap mode)  
Around **31 poloidal turns** are needed to resolve the mode:  $-31\pi < \vartheta < -31\pi$





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Around **31 poloidal turns** are needed to resolve the mode:  $-31\pi < \vartheta < 31\pi$

TAE mode results from a specific interaction of two shear Alfvén waves

Multiple-scale parallel mode structure

The width of the continuum mode has to be resolved (implying high  $k_x$  or, alternatively, large  $L_x$ )

For ideal MHD, SAW width is a delta function resulting in infinite domain in the ballooning space

Kinetically, gyro-radius introduces a finite scale  $\sim \rho_s$



Extraction of the poloidal mode structure:

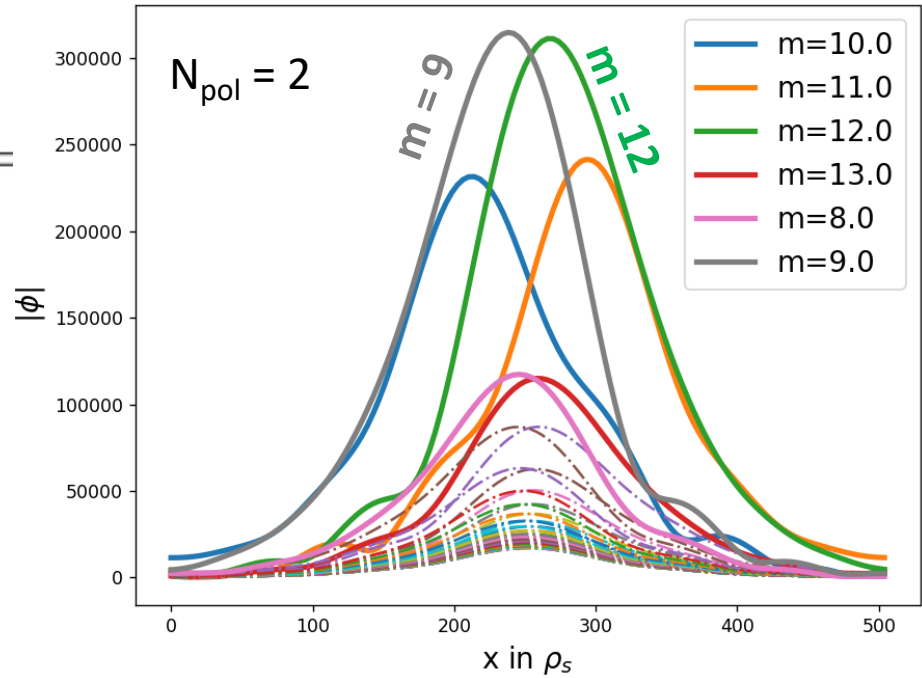
$$x = q_0/(B_0 r_0)(\psi - \psi_0) = r - r_0 ; y = (r_0/q_0)[q(x)\theta - \zeta] ; z =$$

$$q = q_0(1 + \hat{s} x/r_0) ; \hat{s} = (r_0/q_0)dq/dr$$

$$\phi_n = \Phi_n(x, z) \exp(ik_y r_0 \theta) \exp(ik_y x \hat{s} \theta) \exp[-ik_y (r_0/q_0)\zeta]$$

$$n_0 = k_y r_0/q_0 ; m_0 = k_y r_0 = k_y \rho_s/\rho_* ; \rho_* = \rho_s/r_0$$

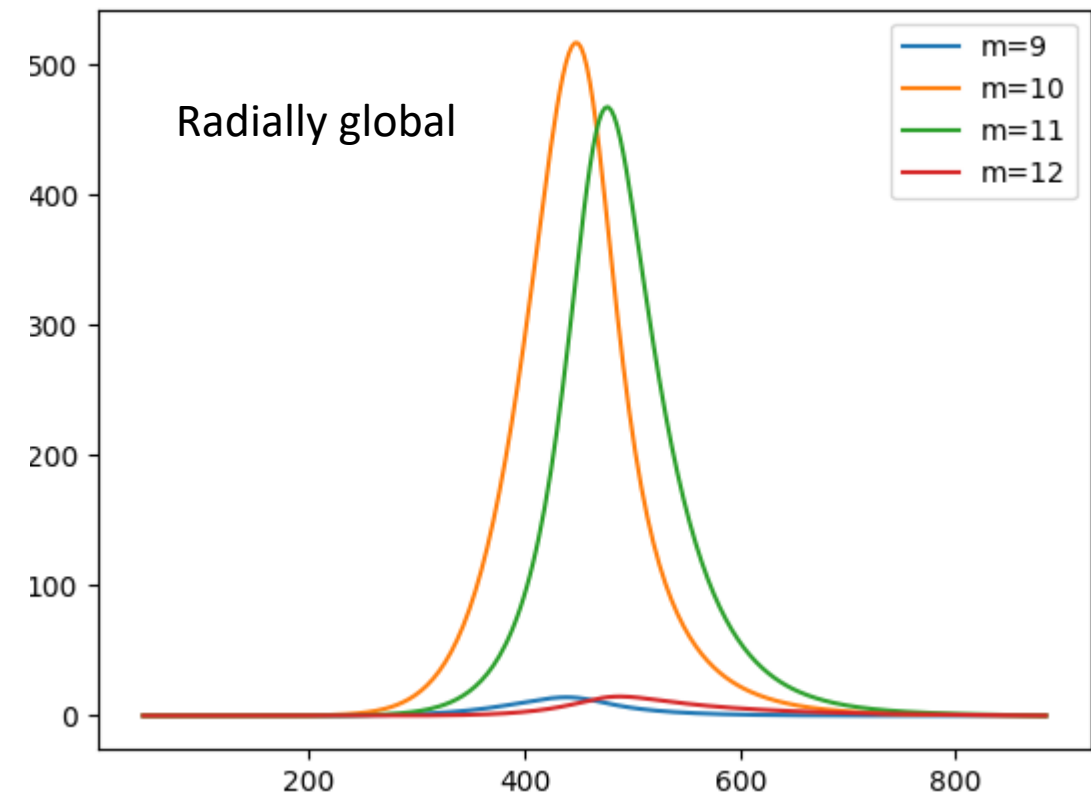
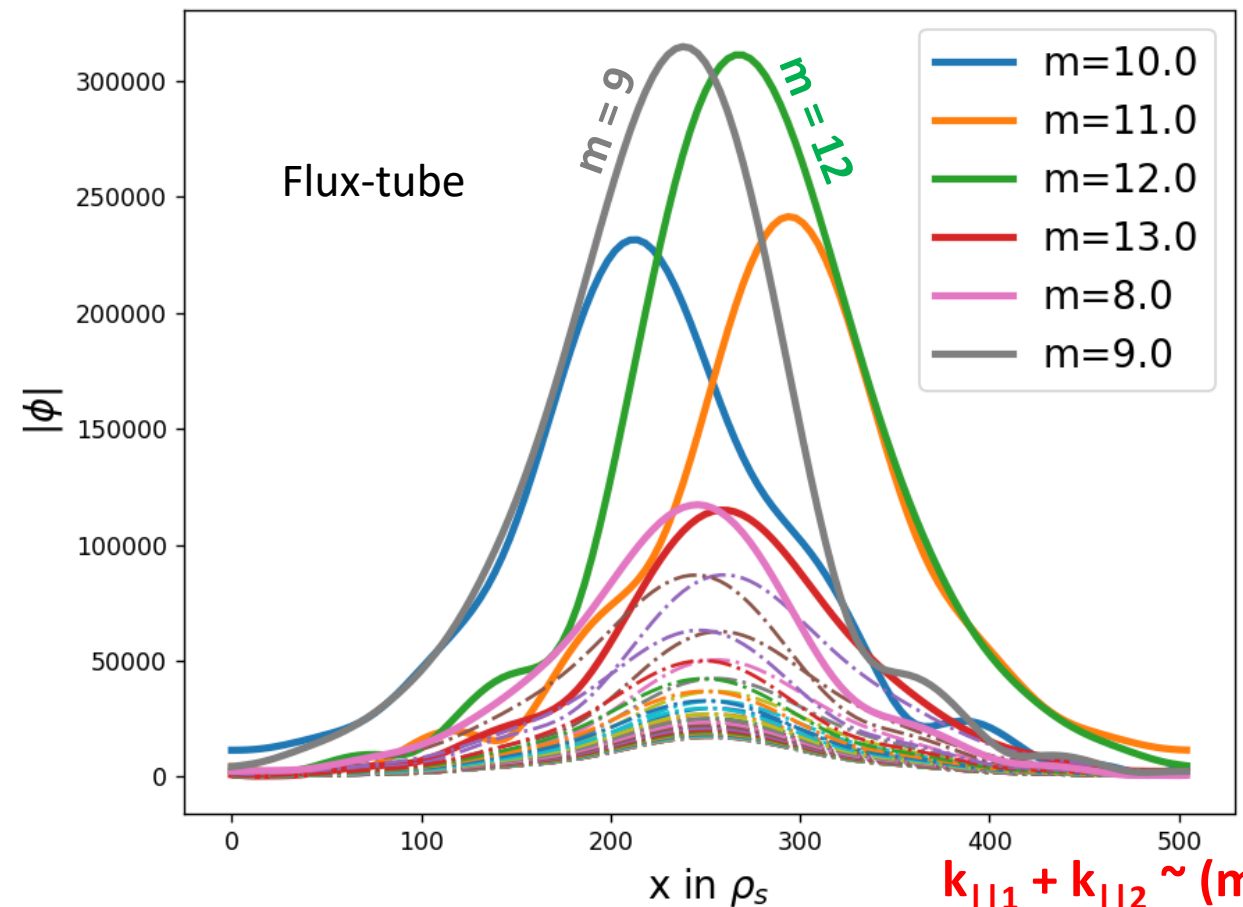
$$\phi_n = \sum_{m'} \left[ \sum_{k_x} \hat{\Phi}_{k_x m' n} \exp(ik_x x) \right] e^{im' \theta} \exp(ik_y x \hat{s} \theta) \exp(im_0 \theta) \exp(-in_0 \zeta), \theta \in [-N_{pol} \pi, N_{pol} \pi]$$



$$k_{||1} + k_{||2} \sim (m_1 - n_0 q_0) + (m_2 - n_0 q_0) = 0$$

- Mode structure of coupled poloidal harmonics (a “standing wave”) is reproduced:  
 $m_1 + m_2 = 2 n_0 q(s_0) = 21$  for  $n_0 = 6$  and  $q(s_0) = 1.75$
- **In contrast to global result, harmonics with  $m_2 - m_1 > 1$  are present and strong!**
- The radial structure of the mode is different : **periodic BC in radial direction (box „multiplied“)**

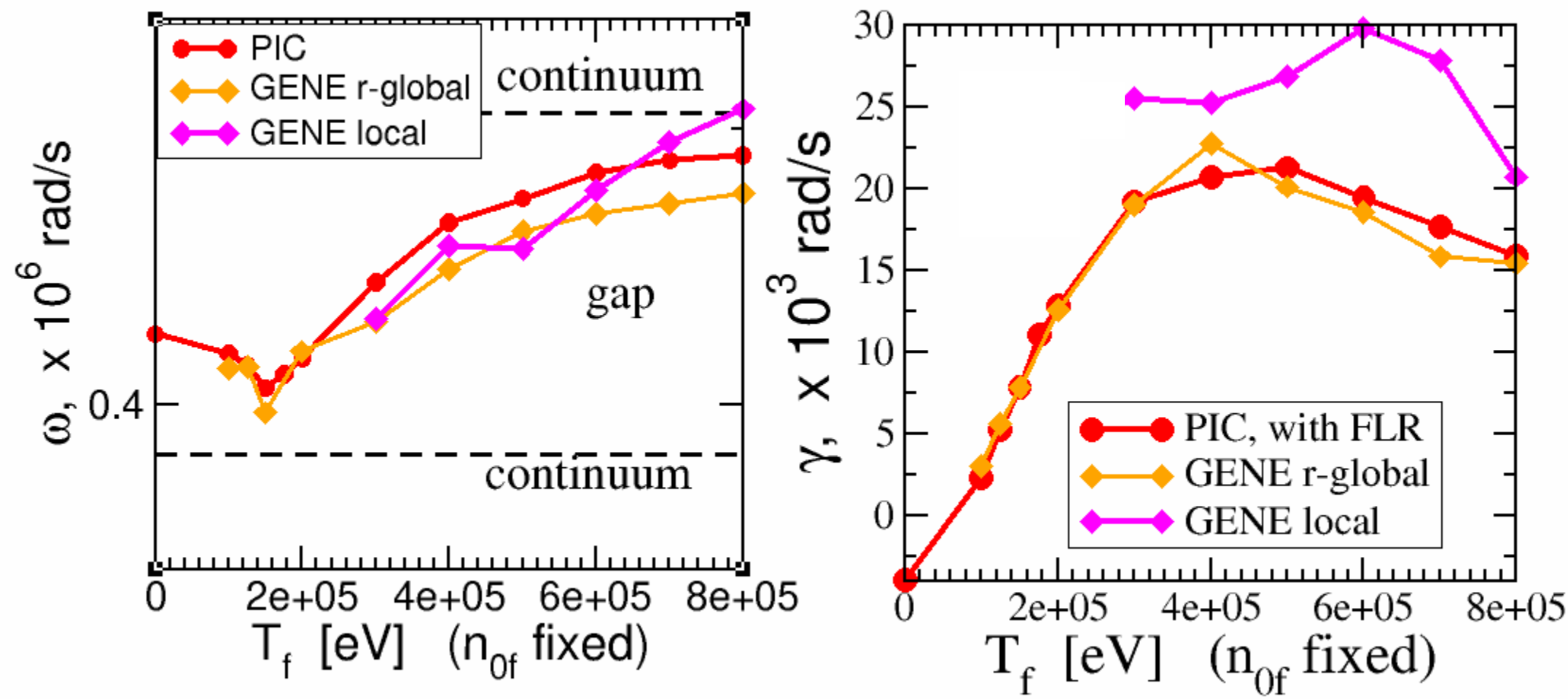
Toroidal mode number:  
 $n_0 = k_y r_0/q_0$   
Poloidal mode number:  
 $m_0 = k_y r_0$



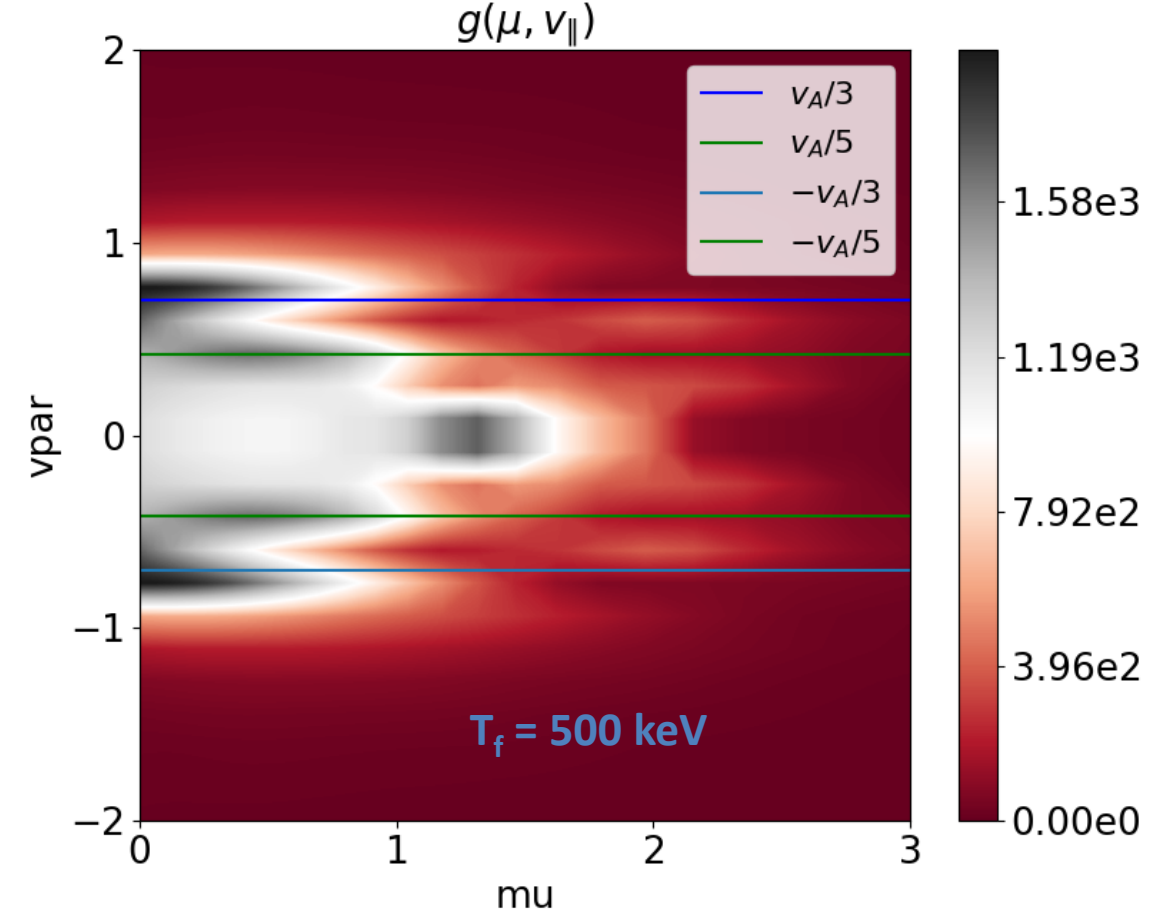
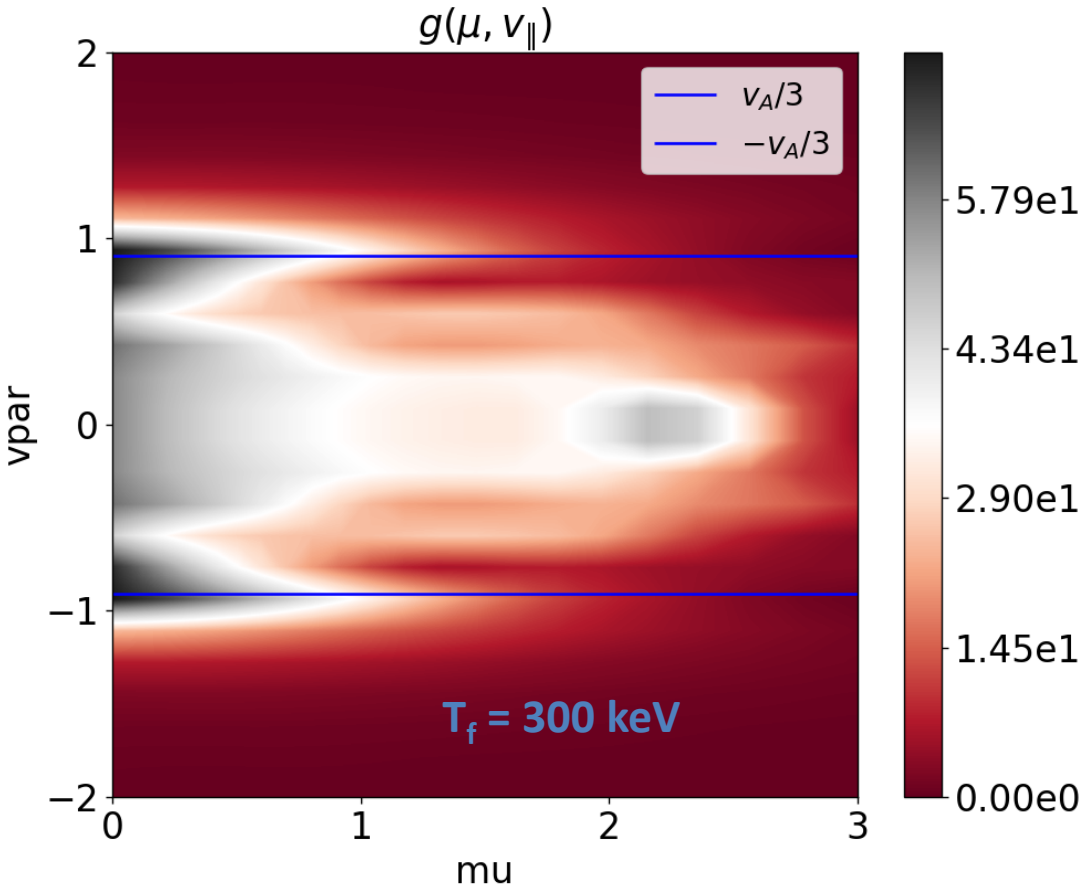
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Toroidal mode number:  
 $n_0 = k_y r_0 / q_0$   
Poloidal mode number:  
 $m_0 = k_y r_0$



Good agreement in the TAE frequency; comparable numbers in the growth rate  
Parallel mode structure at the accumulation point is the key for the mode; included (to some extent) in flux tube



Energetic-particle velocity space: non-adiabatic distribution function

Resonances with energetic passing particles: the main resonance  $v_A$  and the „side-band“ resonances

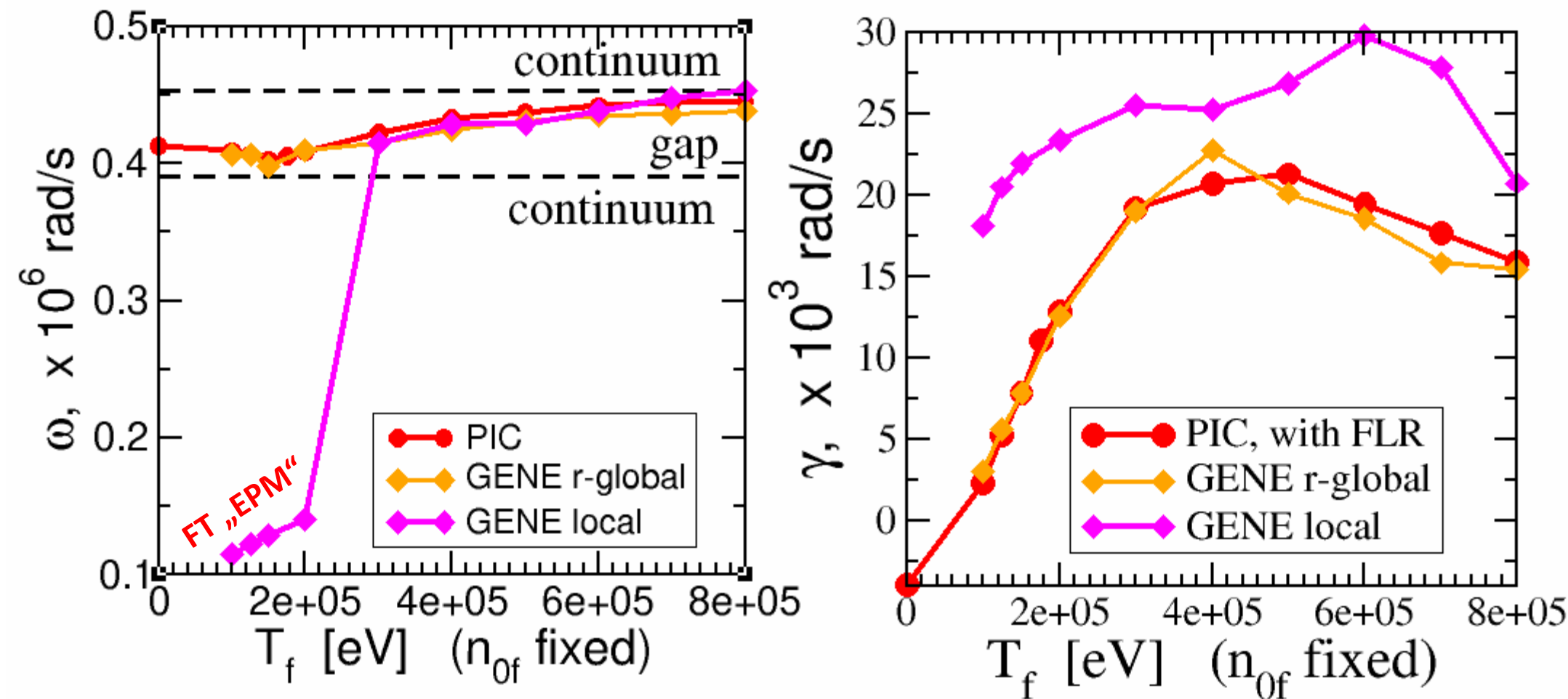
$v_A/(2s - 1)$ :  $v_A/3$ ,  $v_A/5$  (?),  $v_A/7$  (?) ...

In global simulations: only  $v_A/3$  or  $v_A$  are usually relevant

**Resonance condition:**

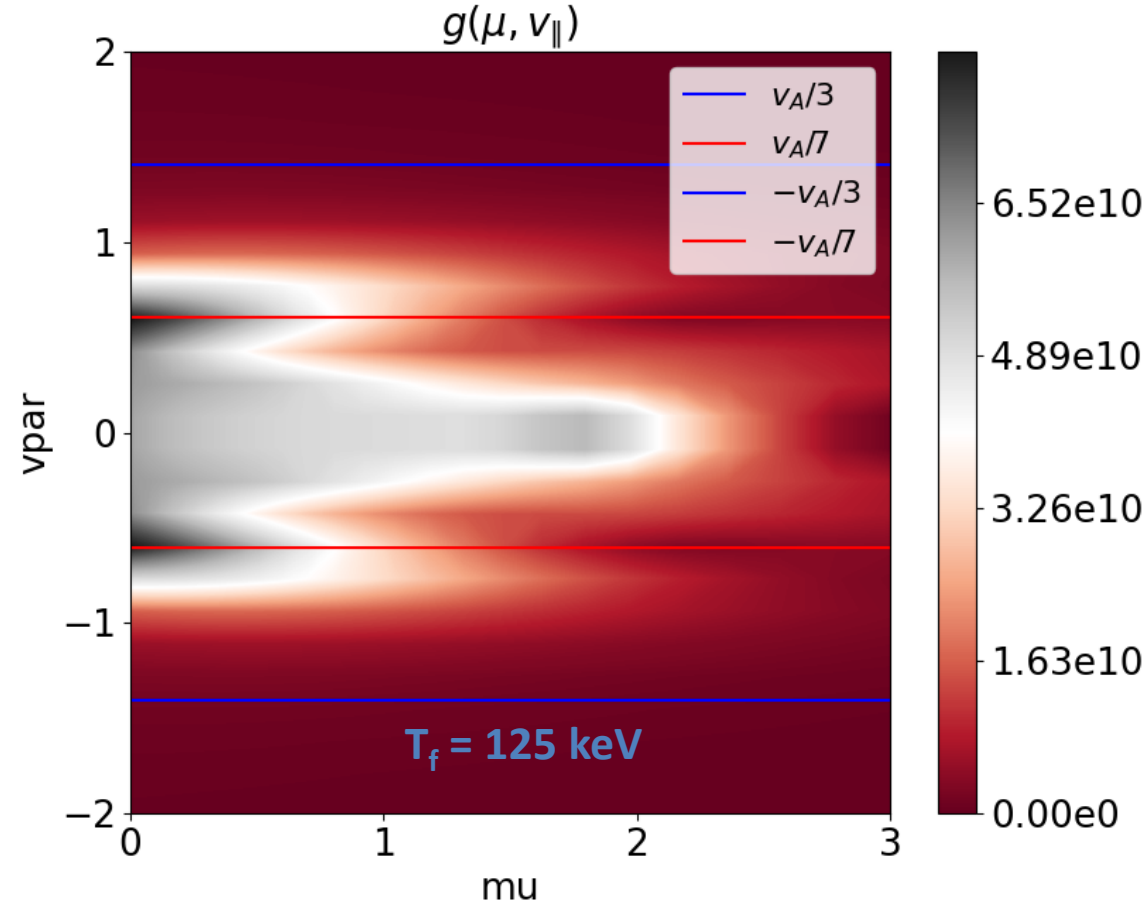
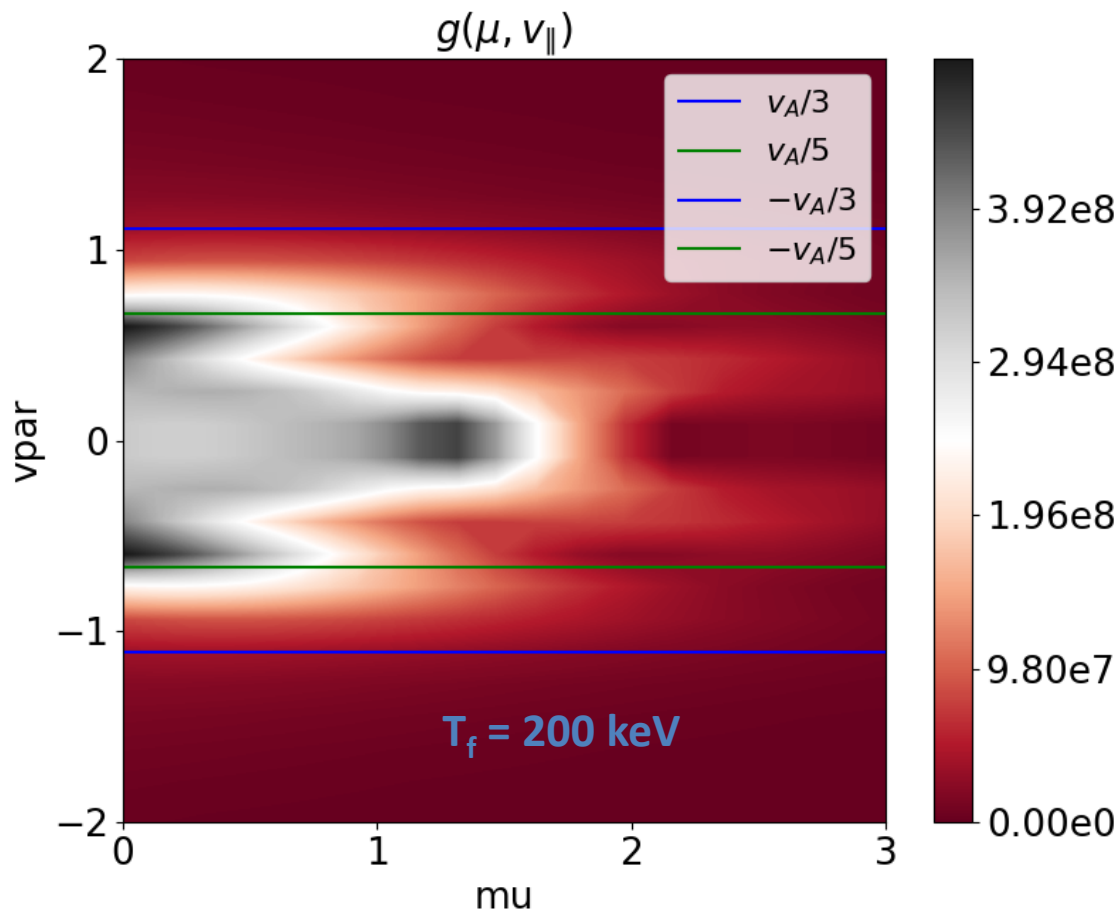
$$k_z v_A = k_{z1} v_{EP} , k_{z1} \neq k_z$$

**(generally)**



For smaller temperatures: jump to an „EPM“ in the flux-tube case (deep into the continuum)  
The „flux-tube EPM“ appears to be well resolved (codewise). Is it real?

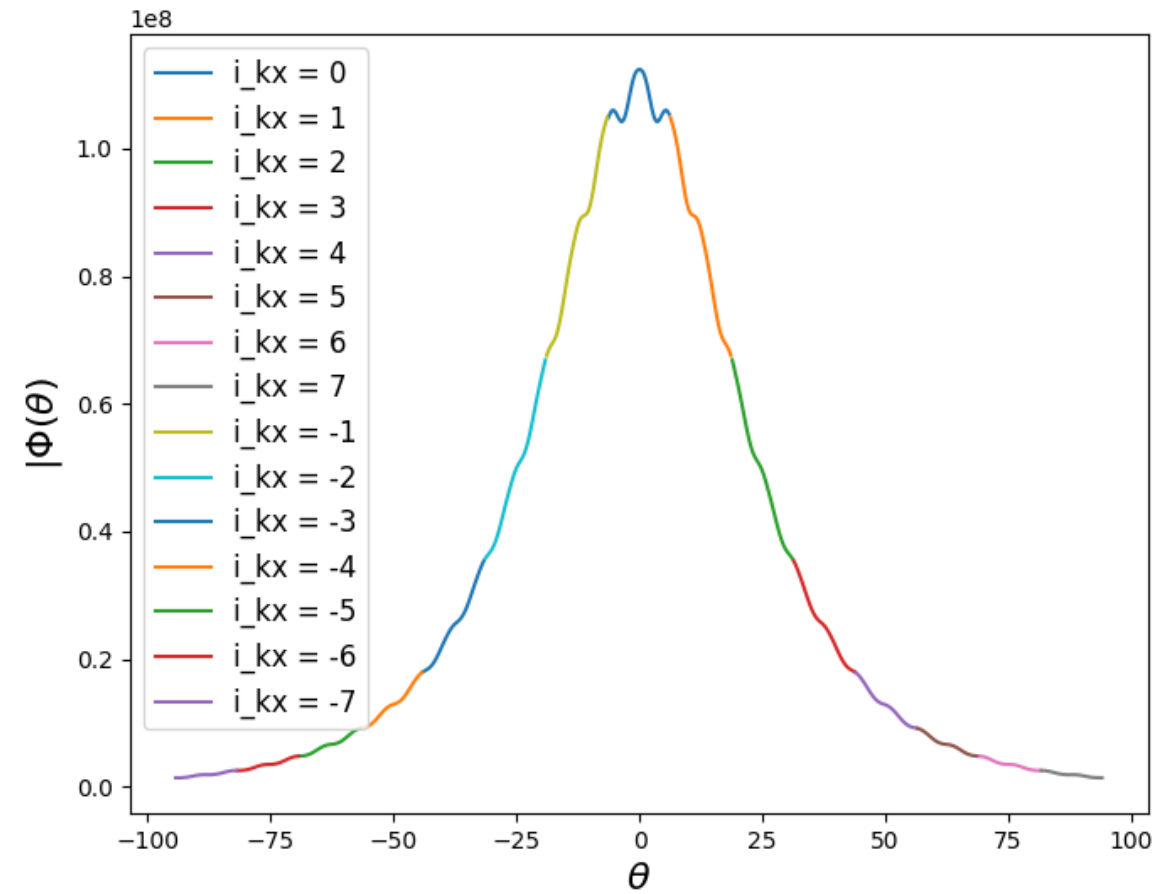
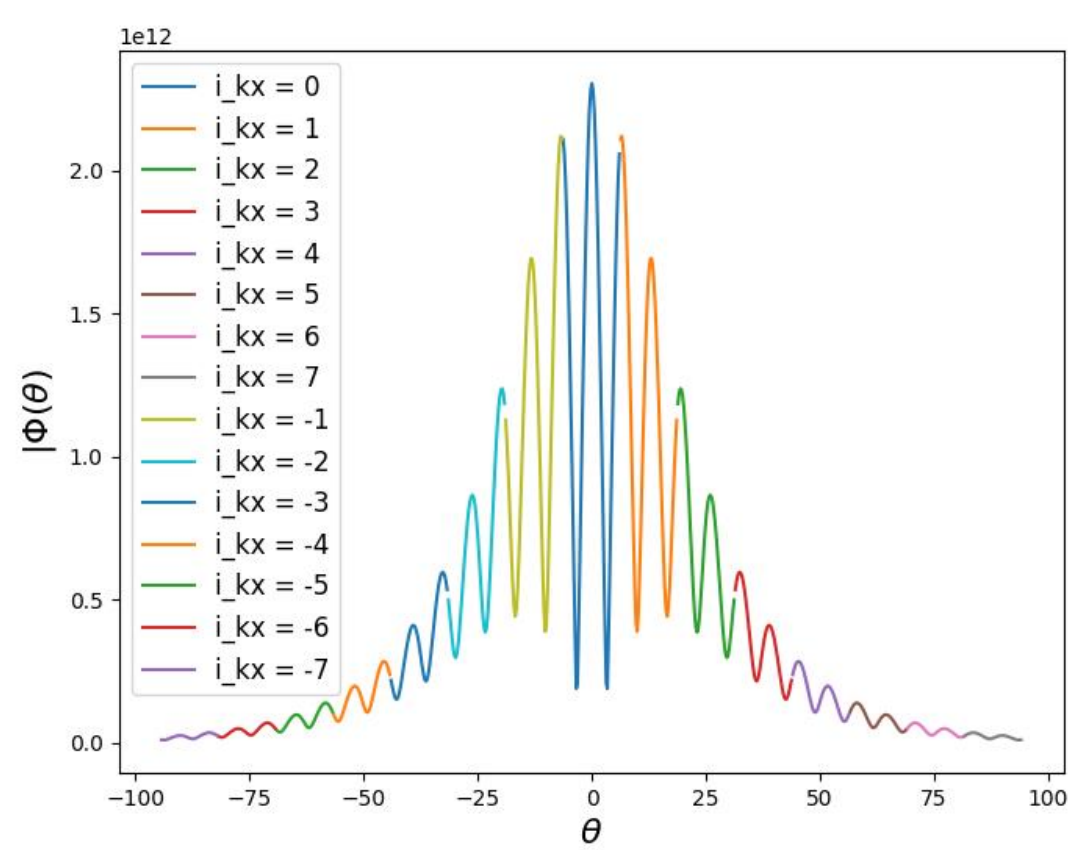




Energetic-particle velocity space: non-adiabatic distribution function  
Higher-order „side-band“ resonances are dominant:  $v_A/5$ ,  $v_A/7$  ; lower frequencies better for drive  
In global simulations: only  $v_A/3$  is usually relevant ; frequency remains in the TAE gap

**Resonance condition:**  
 $k_z v_A = k_{z1} v_{EP}$  ,  $k_{z1} \neq k_z$   
(generaly)

Additional „FT“ resonance because of the FT distortions in the mode structure?



TAE instability (physical gap mode)  
More than **31 poloidal turns** are needed to  
resolve the mode:  $-31\pi < \vartheta < -31\pi$

Flux-tube „EPM“. Is it physical?

**Parallel and radial mode structures  
result from the same dataset**



- Reasonable agreement in frequency for higher  $T_f$ ; similar order of growth rates
- Mode structure is very broad in parallel direction; multiple scales (fine-scale component + envelope)
- The radial mode structure is different (additional couplings): may be a consequence of the periodic boundary conditions (radially) effectively „multiplying“ the box in the radial direction
- Side-band resonances appear ( $v_A/5$ ,  $v_A/7$ ) not so present globally: may be a consequence of the issues with the parallel FT mode structure (failure of the „outgoing wave“ boundary condition; reflections)
- These additional resonances drive an „EPM“ at smaller  $T_f < 300$  keV