



# Runaway electron beam suppression using impurity flushing and large magnetohydrodynamic instabilities

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- Runaway Electrons (RE) are created in thunderstorms and during disruptions of tokamak devices
- They may reach energies up to several 10s of MeV and pose threats to Plasma Facing components
- Shattered Pellet Injection (SPI) as the present disruption mitigation system for ITER aims at:
  - Suppressing the primary RE generation mechanisms (Dreicer & Hot Tail)
  - Suppressing the avalanche multiplication early in the disruption
- But most of state-of-the-art models predict finite primary populations subsequently avalanched [*Vallhagen JPP 2020, Fülöp IAEA 2016*]
- → A second line of defense is needed: in-flight suppression of a fully-accelerated RE beam

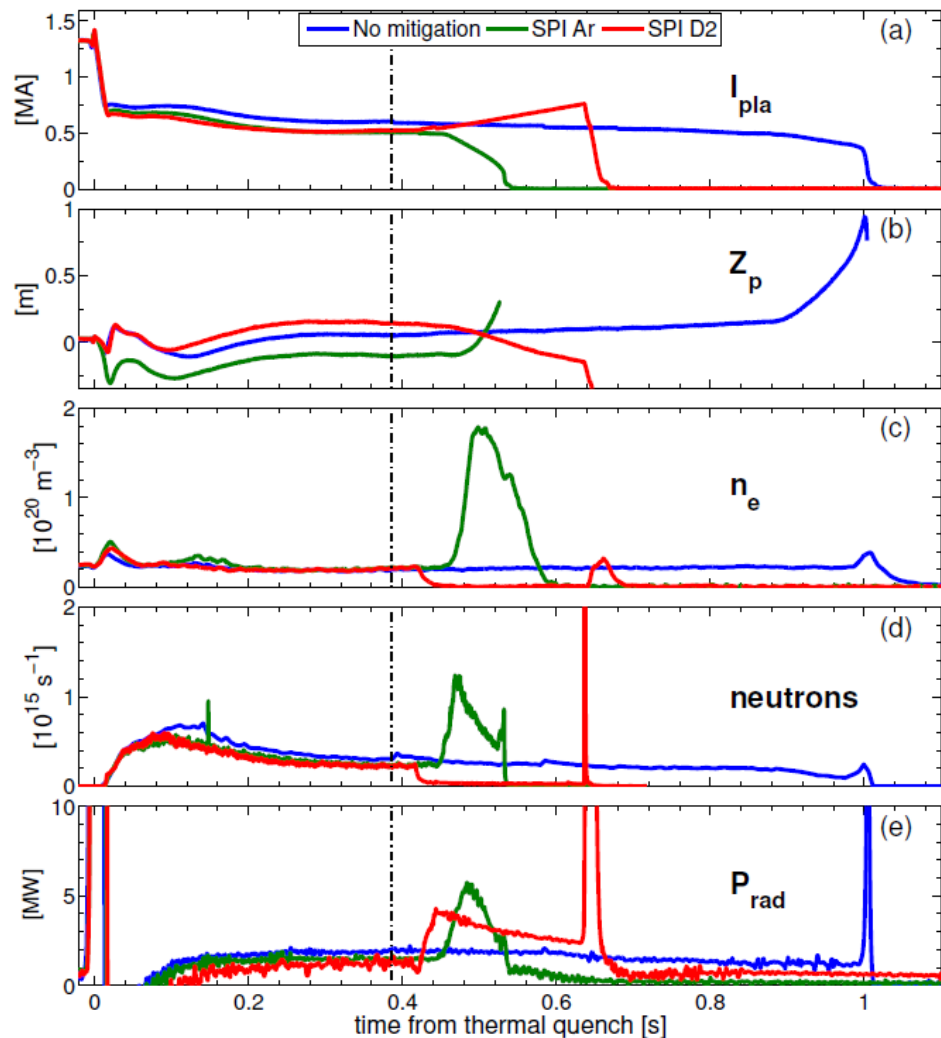


- High-Z SPI (or MGI) species have been tested for RE beam mitigation in the past on DIII-D, JET, AUG, Tore Supra
  - Some success on DIII-D, Tore Supra, AUG, [[Whyte PRL 2002](#), [Shiraki NF 2018](#), [Saint-Laurent EPS 2009](#), [Pautasso NF 2020](#)]
  - Very limited success under certain conditions in JET [[Reux NF 2015](#)]
- Increasing evidence from theoretical models that high-Z mitigation is not enough for large ITER-class RE currents [[Martin-Solis NF 2017](#), [Vallhagen JPP 2020](#), [Hesslow NF 2019](#)]
- Main topic of the present article: use of deuterium to mitigate a RE beam
  - Up to 1.27 MA of runaways dissipated **without measurable heat loads**
  - Builds upon similar deconfinement events observed at DIII-D [[Paz-Soldan PPCF 2019](#)]

# Runaway beam scenarios : no mitigation



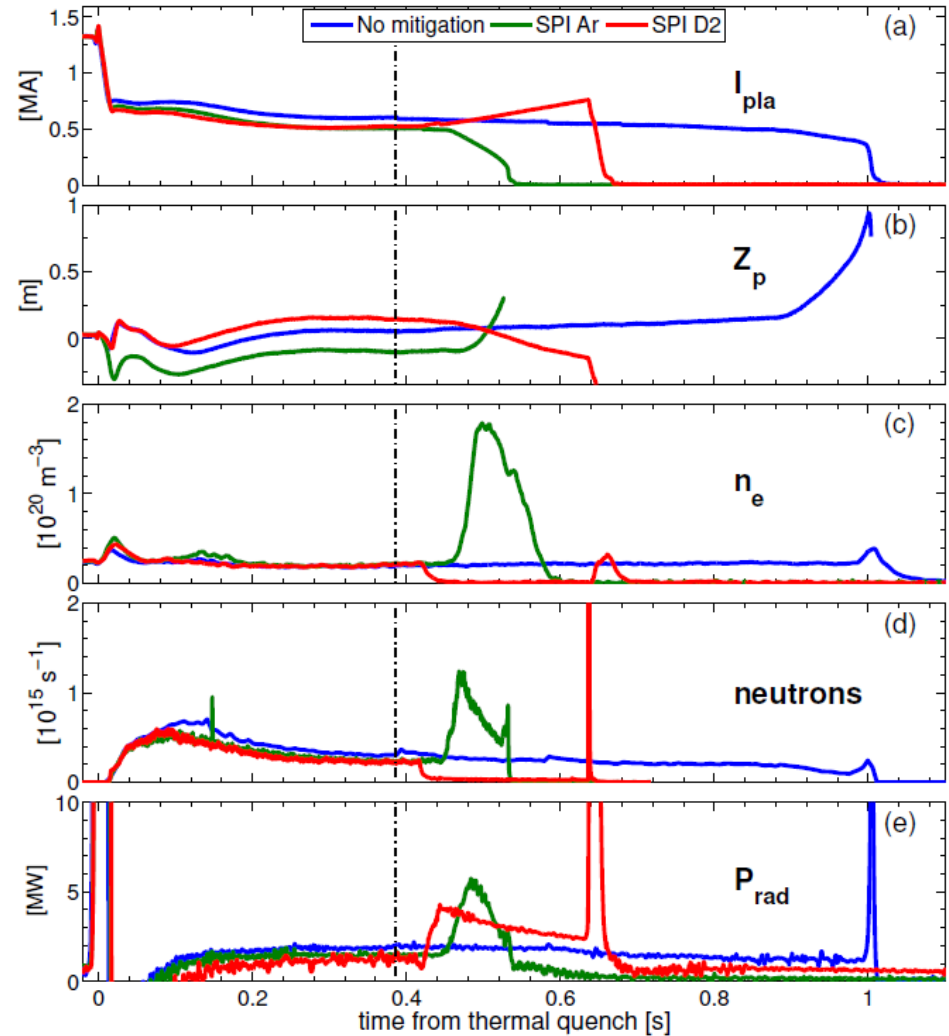
- JET « no-mitigation » beam: **(blue curve)**
- Limiter configuration @1.5 MA/3.0 T
- $2.38 \cdot 10^{21}$  atoms of argon used to trigger the disruption
- 750 kA RE beam slowly decaying in current.
- $n_{e, \text{line-av}} \sim 10^{19} \text{ m}^{-3}$ ,  $T_e \sim 5\text{-}15 \text{ eV}$  [*Sridhar NF 2020*]
- $P_{\text{rad}} \sim 2 \text{ MW}$  due to the argon companion plasma
- Impact at termination on localized areas, with measurable heat loads



# Runaway beam scenarios: high-Z mitigation



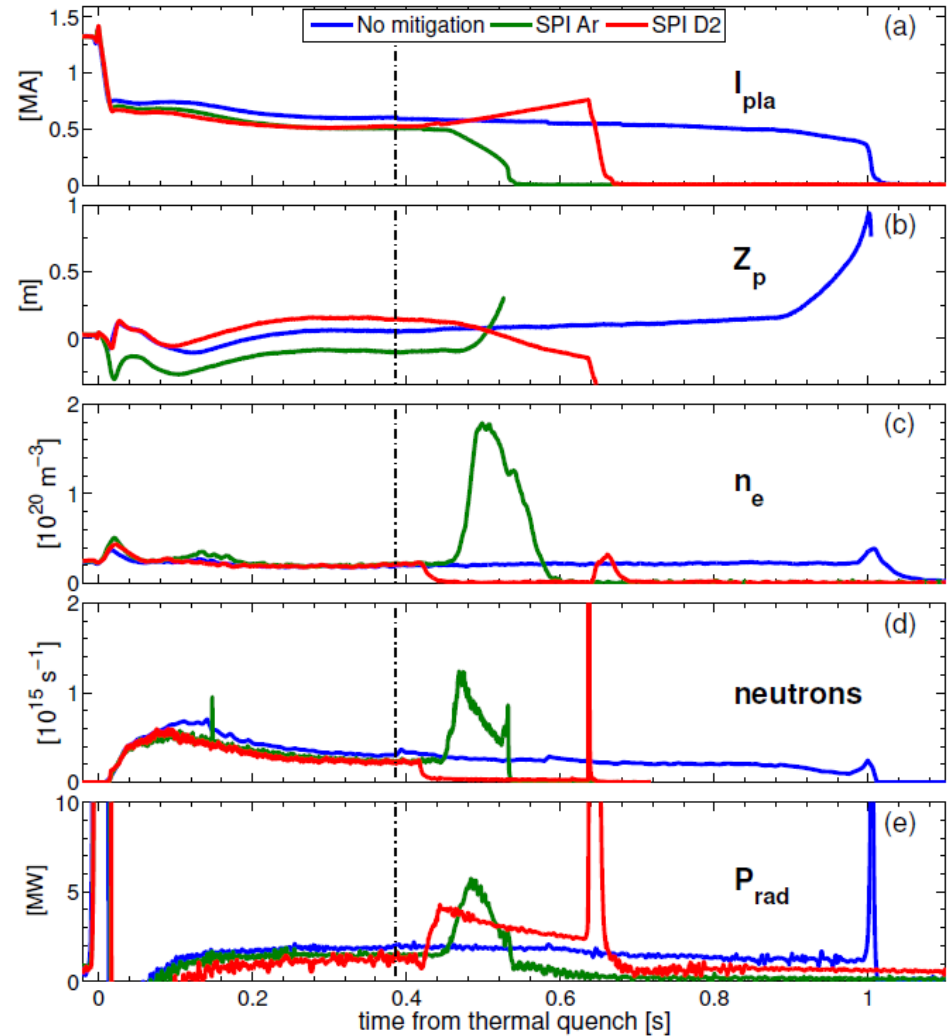
- JET Argon SPI scenario (green curve)
- Faster current decay
- Increase of free electron density,  $P_{\text{rad}}$  increase due to the density and impurity content increase
- Vertical destabilization
- Localized impact on the wall, with significant heat loads



# Runaway beam scenarios: D<sub>2</sub> mitigation



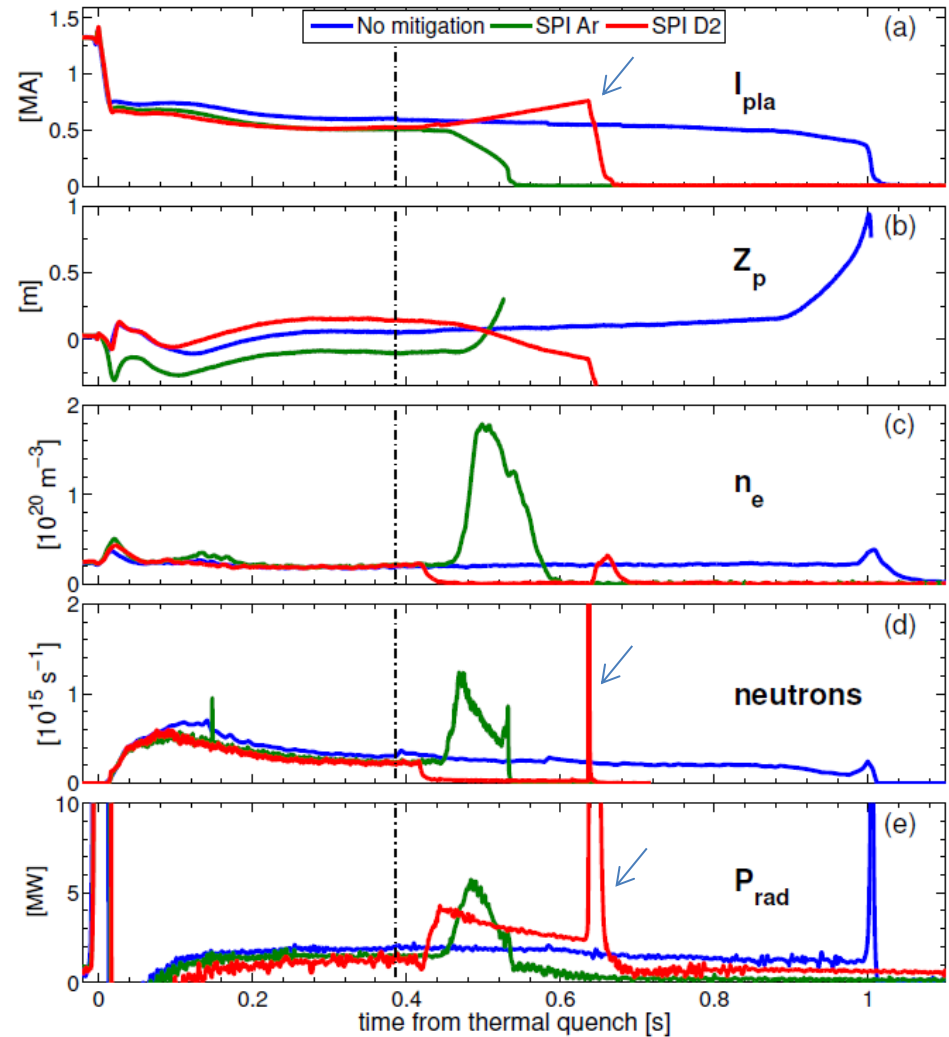
- JET D<sub>2</sub> SPI scenario (red)
- $1.6 \cdot 10^{23}$  atoms injected by SPI
- Current increases → decrease of the RE+companion plasma resistivity
- $n_e$  drops to non-measurable values ( $<10^{18} \text{ m}^{-3}$ ) → Plasma recombination
  - Argon expelled from the plasma (VUV spectroscopy)
  - Process not yet clear
- Neutron rate drops by a factor 10
- $P_{\text{rad}}$  increases to 4 MW.
- Measured loop voltage in good agreement with voltage derived from Bethe Stopping power for a neutral gas

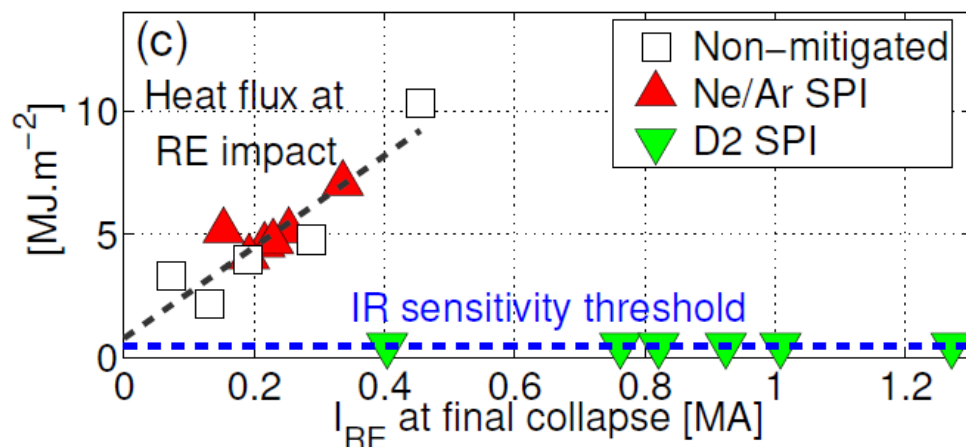


# Runaway beam scenarios: D<sub>2</sub> mitigation



- 220 ms after the D<sub>2</sub> shard plume arrival ( $I_{RE} \sim 770$  ms)
- neutron spike indicating RE losses
- Complete disappearance of RE synchrotron emission (IR cameras) in  $< 3$  ms
- Current decays over  $\sim 12$  ms
- Radiated power spike
- Rest of the current decay: purely ohmic
  - In some cases: current spike similar to the one from a normal disruption





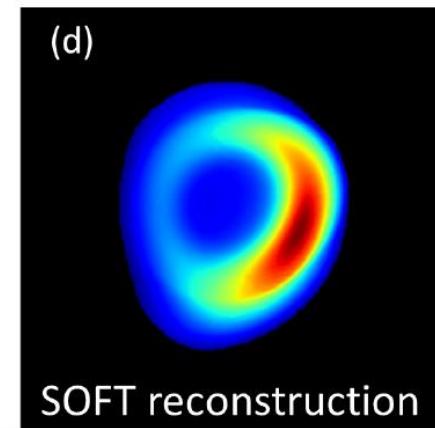
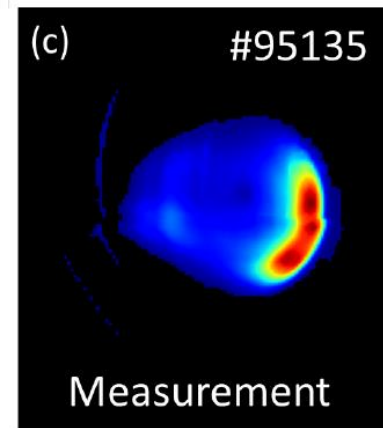
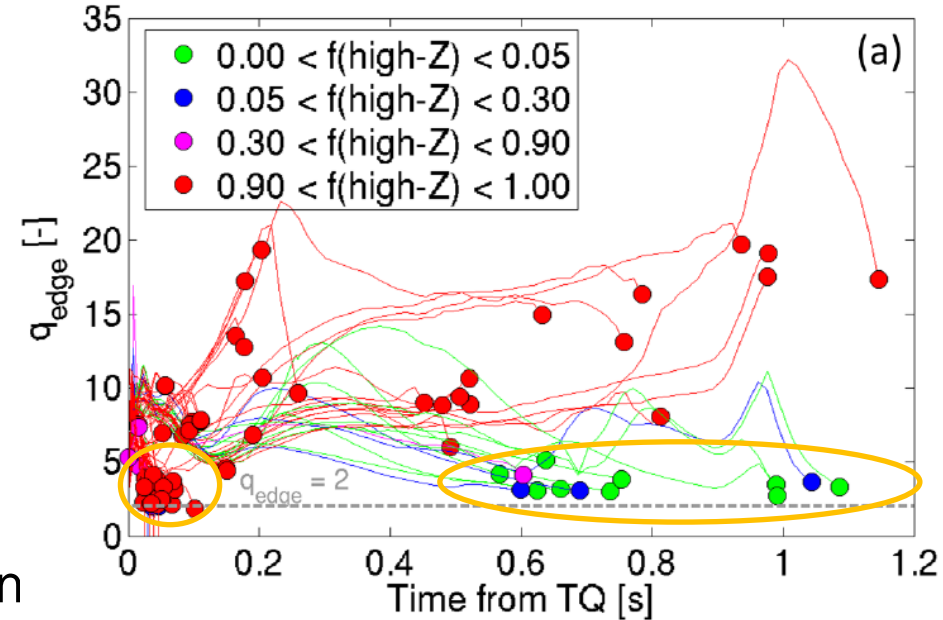
- Most prominent feature: absence of measurable heat loads at beam termination
  - <0.5 MJ.m<sup>-2</sup>
  - High-Z heat loads up to 10 MJ.m<sup>-2</sup> despite lower currents
- **Two mechanisms at play:**
  - **A large and brief MHD instability dissipating the runaway electrons**
  - **The absence of RE regeneration (absence of conversion from magnetic to kinetic energy) during the final collapse**



# Development of the MHD instability : j-profile



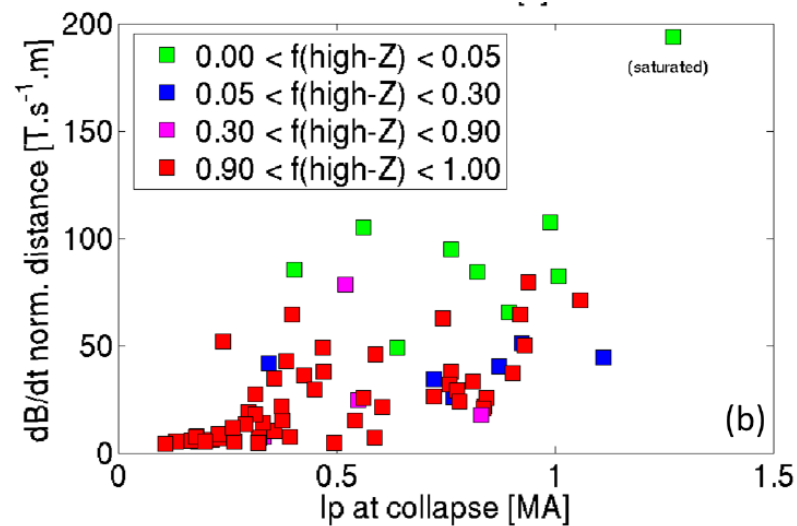
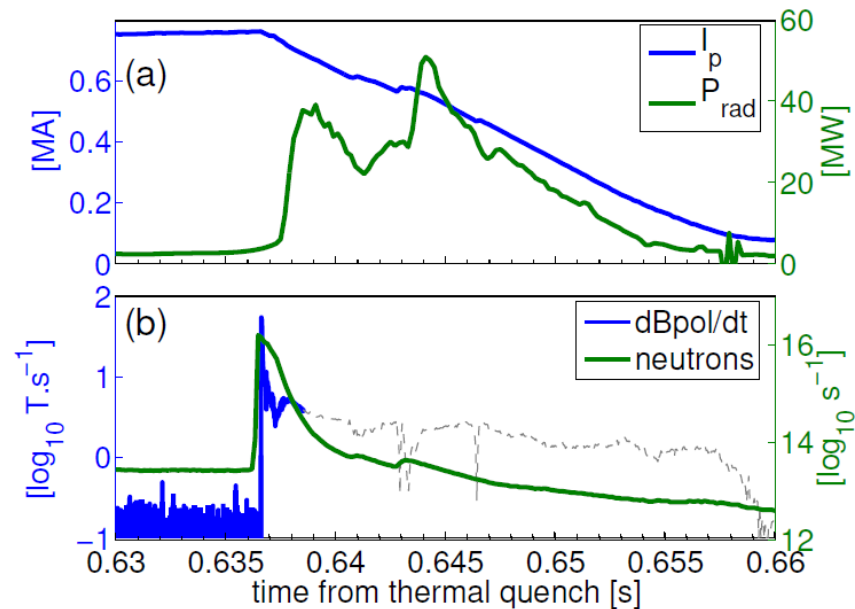
- Current rise leads to low  $q_{\text{edge}}$ .
- For D<sub>2</sub>-mitigated scenarios, final collapse happens when  $q_{\text{edge}}$  reaches 2-5
- Not specific to D<sub>2</sub> cases
- Probably not a current-limiting instability as suggested in [Paz-Soldan PPCF 2019]
- Reconstruction of the synchrotron emission using the SOFT code [Hoppe NF 2018]
- Best match obtained when current profile is hollow (peak at  $a/2$ ). Energy <15 MeV, pitch angle 0.1-0.3. No  $f(E)$  and pitch can match the observation with a peaked profile



# MHD Instability characterization

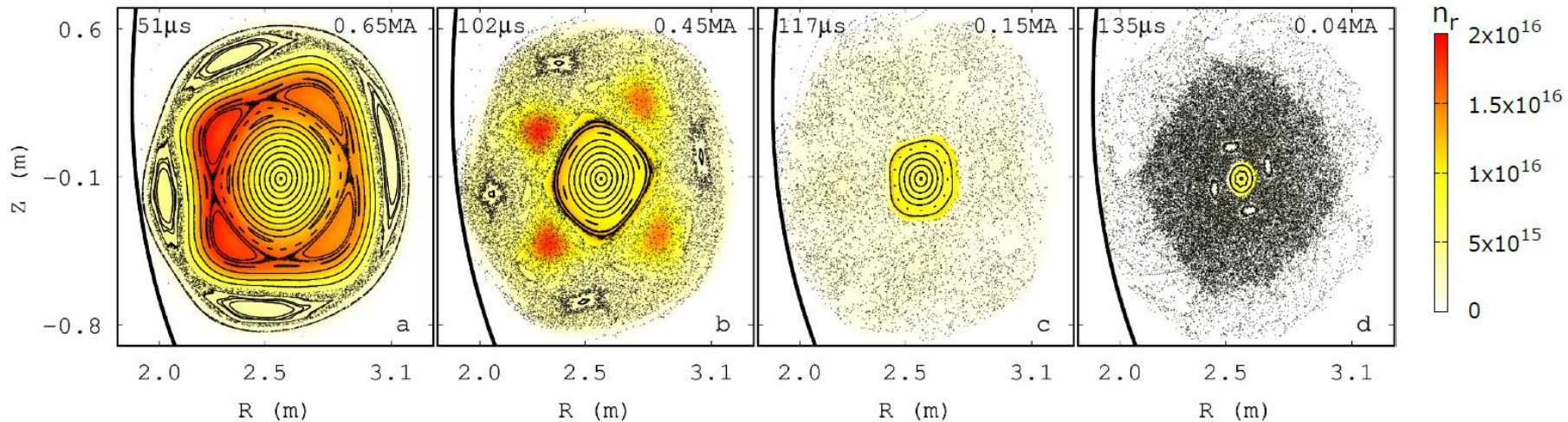


- Islands visible in the IR pattern
  - $m=4$  visible at  $a/3$
- $n=1$  from Mirnov coils (toroidal array)  $\rightarrow q=4 \rightarrow$  further evidence for a hollow profile
- The instability itself develops on a 10-20  $\mu\text{s}$  timescale (peak dB/dt)
- Experimental magnitude of the instability:
  - $(\text{dB}/\text{dt})_{\text{max}}$  weakly correlated with the high-Z fraction ( $D_2$  cases vs. High-Z cases): all  $D_2$  cases have large  $(\text{dB}/\text{dt})_{\text{max}}$ , but some high-Z cases too.
  - dB/B even less correlated
  - $\rightarrow$  growth rate rather than dB/B appears to be the key feature.





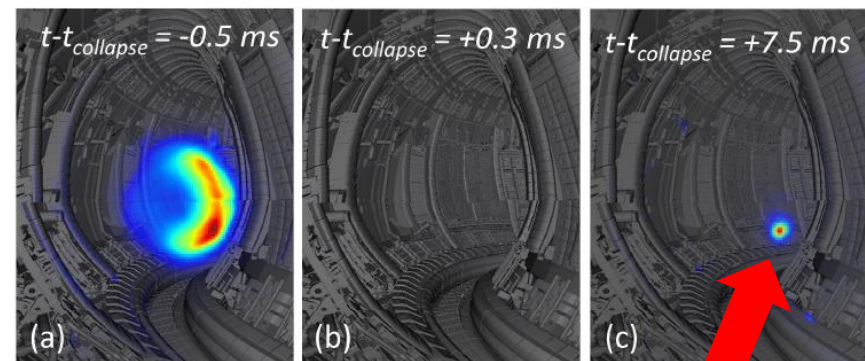
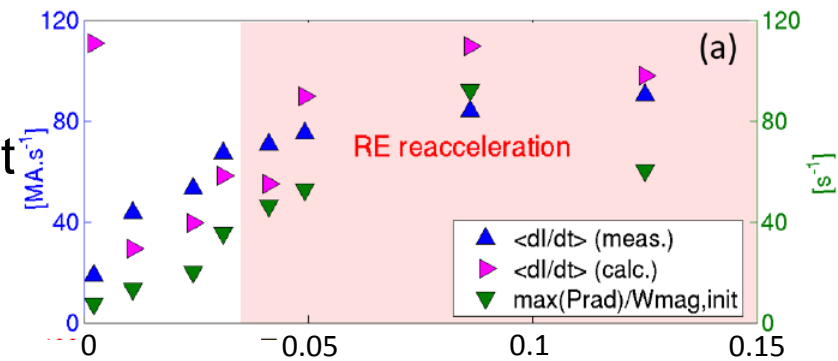
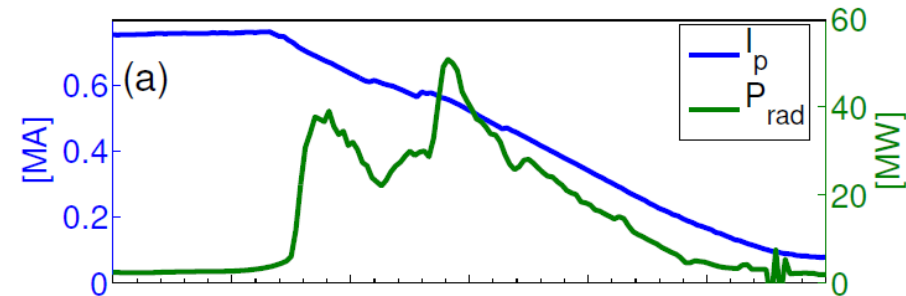
- JOREK simulations have been made using a RE fluid model [[Bandaru Phys Rev E 2019](#)]
- Initial current profile determined from SOFT simulations
- Result: MHD dominated by tearing mode from the outside  $q=4$  surface.
  - Stochastization starts in the edge
  - Confinement destroyed in  $100 \mu\text{s}$   $\rightarrow$  timescale compatible with measurements
  - 95% of runaways are lost. RE losses deposited near the limiter contact point but over a wider area.



# RE regeneration during the final collapse



- After the collapse : current carriers shift from RE to thermal bulk
- Radiated power and free density increases
- Argon line emission is back → argon not completely « purged »
- Maximum radiated power and current quench rate during final collapse correlated with the  $Ar/D_2$  ratio in the neutral species injected into the vessel (triggering MGI + mitigating SPI) → further evidence of the incomplete purge
- When Ar fraction gets higher: reappearance (regeneration) of a small RE beam during collapse





- Behaviour investigated using a 0D model (lumped circuits + RE population evolution)

$$\frac{3}{2} \frac{\partial}{\partial t} n_f T_e = \frac{(I - I_{RE})^2}{\sigma S^2} - n_f n_Z L(T_e) \quad \text{Energy balance for the companion plasma.}$$

$$\frac{d}{dt} (LI + L_v I_v) = -2\pi RE \quad \text{Ohm's law for the plasma}$$

$$\frac{d}{dt} (L_v I + L_v I_v) = -I_v R_v \quad \text{Ohm's law for passive conductive structures}$$

$$\frac{1}{I_{RE}} \frac{\partial I_{RE}}{\partial t} \approx \frac{n_f + n_b}{n_f \ln \Lambda_f(p_c) + n_b \ln \Lambda_b(p_c)} \frac{1}{\sqrt{Z_{RE}(p_c) + 5}} \frac{e(E - E_{crit})}{m_e c}$$

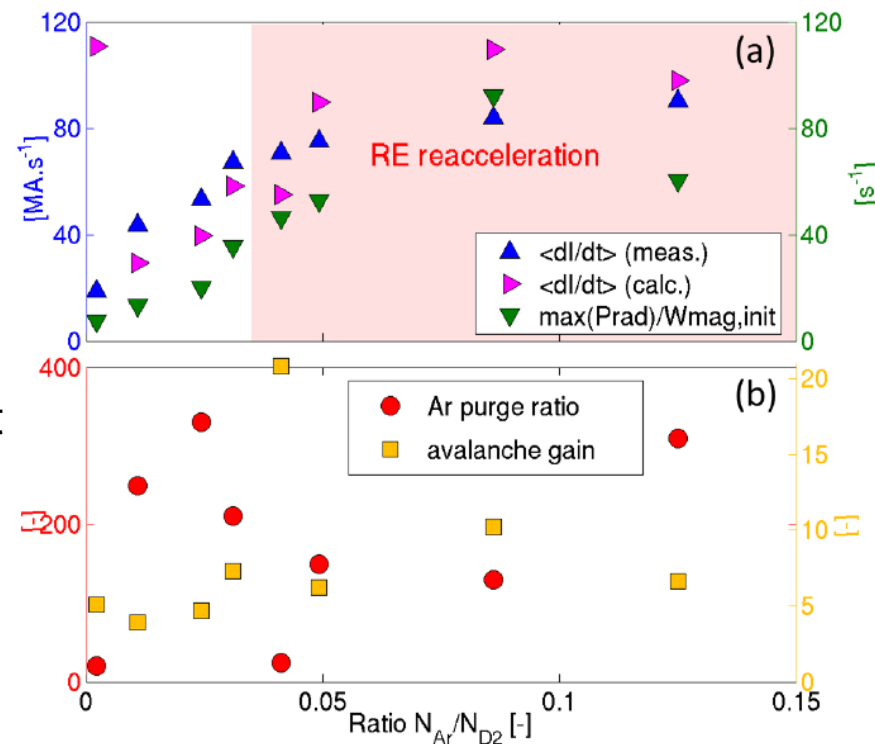
Ruanway population evolution (avalanche)  
 [Aleynikov NF 2017, Martin-Solis PoP 2015,  
 Hesslow NF 2019, Breizman NF 2019]

- Takes into account partial screening of impurities
- Critical momentum obtained from acceleration-friction force balance
- Computes how fast the plasma reheats compared to the avalanche timescale → self consistent temperature, current and RE generation

# RE regeneration modelling



- Numerical parameters:
  - Vessel resistive time  $L_v/R_v = 5$  ms,  $L_v = 2\mu\text{H}$  as in [Strauss PoP 2017]
  - $L_z$  taken from ADAS
  - $I_i = 0.5$ , Spitzer resistivity
  - Plasma cross section and pre-collapse current from measurement
- Results:  $dI/dt$  in good agreement with the experiment
- Argon purge rate between 50 and 300, confirming the purge mechanism
- Avalanche gain too low to regenerate a full RE beam, but small correlation between avalanche gain and  $\text{Ar}/\text{D}_2$  ratio.
- If  $\text{Ar}/\text{D}_2$  ratio higher: continuous reacceleration during collapse

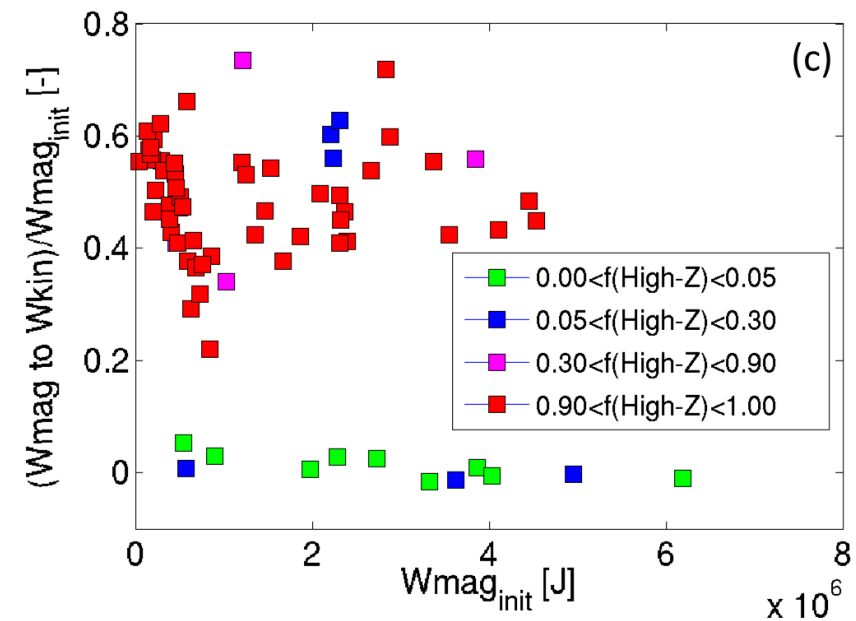


Global dynamics of final collapse well captured by the model

# Conversion of magnetic into kinetic energy



- Continuous RE regeneration during the collapse: plays a role in the conversion from magnetic to kinetic energy of the runaway beam.
- Typical case:  $W_{\text{mag}} = 2.2 \text{ MJ}$ ,  $W_{\text{kin}} = 0.4 \text{ MJ}$ .
  - $\rightarrow$  much larger damage if  $W_{\text{mag}}$  is converted into  $W_{\text{kin}}$
- Conversion rate calculated using the method proposed in [Loarte NF 2011] and adding radiated power
- Conclusions: conversion rate close to zero for cases where high-Z material  $< 30\%$
- High-Z and low-Z cases clearly distinguished





- D<sub>2</sub> RE beam mitigation was found to be efficient and reproducible at JET provided that the companion plasma is pure enough
  - Large enough D<sub>2</sub>/High-Z ratio
  - More experiments planned to better assess the boundary conditions
- Extrapolability to ITER and future devices is now the main question
  - Is a large and brief enough MHD instability accessible on ITER?
    - More MHD simulations (JOEUK)
  - Modeling predicts the avalanche gain will be larger for ITER
    - Can an arbitrarily high companion plasma purity level be reached?
- Even if runaways are re-accelerated: only a fraction of the initial RE current can be regenerated
  - → Repetitive D<sub>2</sub> SPI could be used to do a stepwise reduction of the RE current down to a tolerable level
- So far: one of the best hopes for RE beam mitigation on large machines



