

2021.5.18 (Tue)

IFERC-CSC Workshop on JFRS-1 Projects

Virtual Event



# Global Gyrokinetic simulation for High-Beta plasma [Project Name: GGHB]

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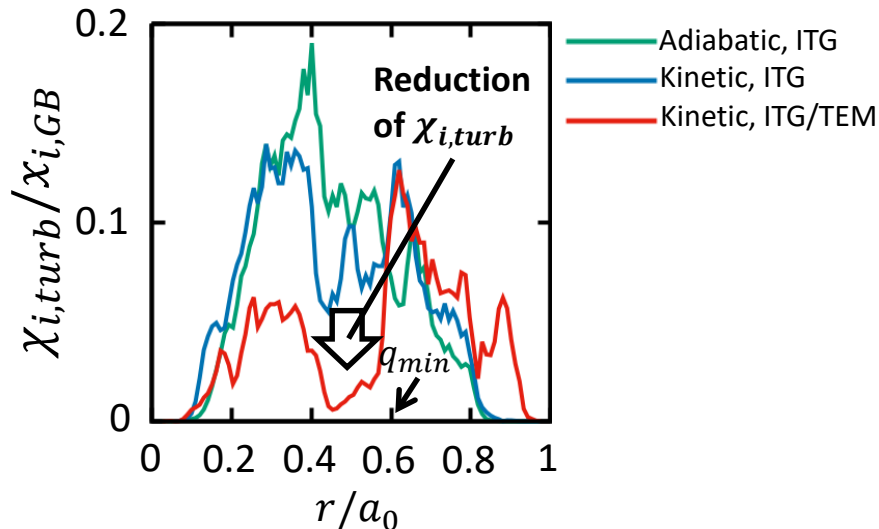
# Background & Motivation

- ✓ The plasma beta value given by  $\beta = 8\pi p/B^2$  is one of the important parameters for fusion plasmas because it is linked to the fusion reaction rate and also related to the production of bootstrap current.
- ✓ To achieve and sustain high-beta states, our group focuses on the following two topics;

## (1) Study of ITB formation

### for the achievement of high-beta plasma

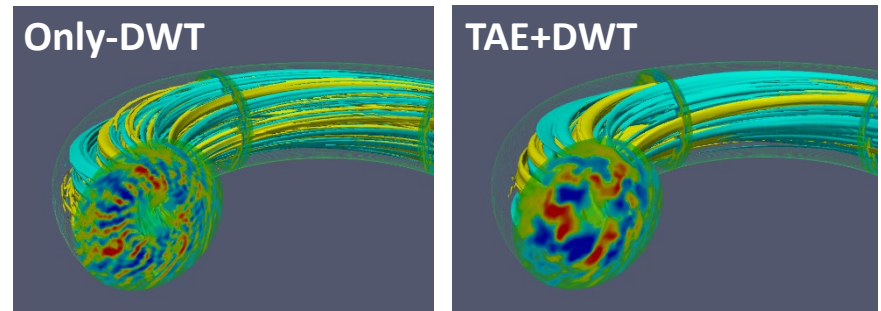
Based on **full- $f$  electrostatic model with hybrid electron**, we have studied **the spontaneous ITB formation in reversed magnetic shear configuration**.



## (2) Study of TAE-KBM interaction

### for the sustainment of high-beta plasma

Based on **delta- $f$  electromagnetic model with kinetic electron**, we have studied **the interaction between energetic-particle-driven MHD mode (TAE) and drift-wave turbulence**.





# Contents

1. Background & Motivation (1/20)

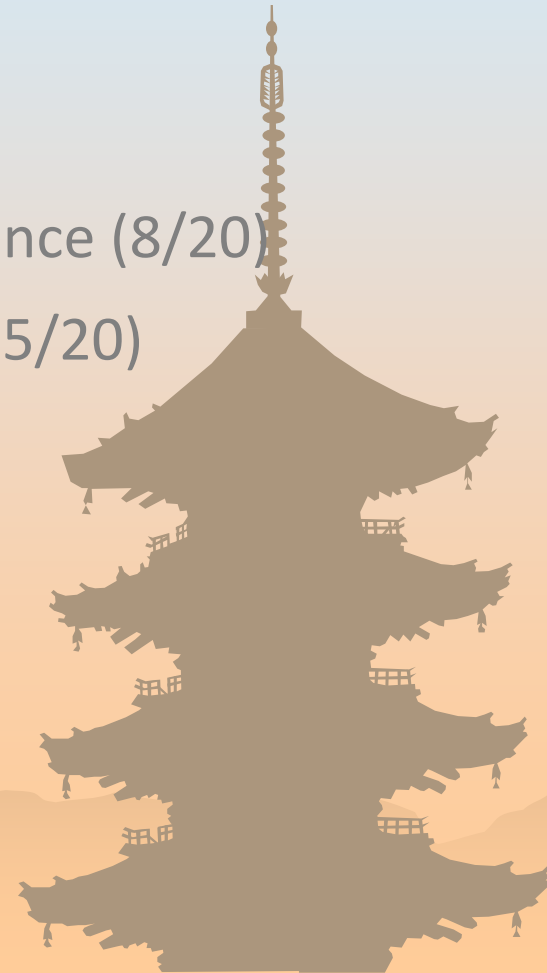
**2. Full- $f$  Gyrokinetic Code *GKNET* (4/20)**

- \* **Original *GKNET***
- \* **Extension of *GKNET***
- \* **Parallelization of *GKNET***

3. Spontaneous ITB formation in ITG/TEM Turbulence (8/20)

4. Interaction between TAE and KBM Turbulence (5/20)

5. Summary & Future plans (2/20)



# Original *GKNET*

Gyrokinetic Vlasov equation (ion) & Gyrokinetic quasi-neutrality condition

$$\frac{\partial}{\partial t}(Jf) + \frac{\partial}{\partial \mathbf{R}} \left( J \frac{d\mathbf{R}}{dt} f \right) + \frac{\partial}{\partial v_{\parallel}} \left( J \frac{dv_{\parallel}}{dt} f \right) = S_{src} + S_{snk} + C_{coll}$$

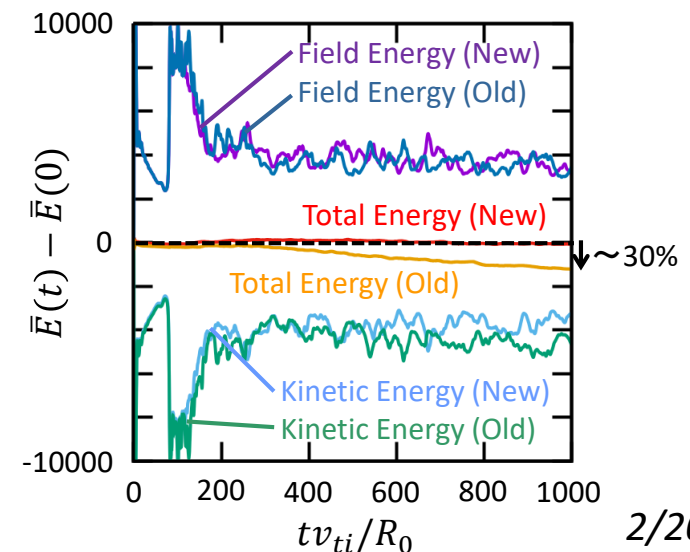
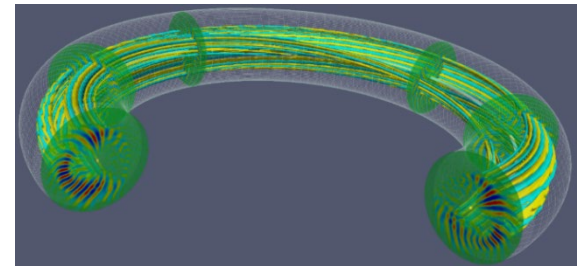
$$\frac{d\mathbf{R}}{dt} = \{\mathbf{R}, H\} = \frac{\mathbf{B}_{\parallel}^*}{B_{\parallel}^*} \frac{\partial H}{\partial v_{\parallel}} + \frac{1}{B_{\parallel}^*} \mathbf{b} \times \nabla H$$

$$-\nabla_{\perp} \cdot \frac{\rho_{ti}^2}{\lambda_{Di}^2} \nabla_{\perp} \phi(\mathbf{R}) + \frac{1}{\lambda_{De}^2} [\phi(\mathbf{R}) - \langle \phi(\mathbf{R}) \rangle_f] = 4\pi e \iint \langle \delta f_i(\mathbf{R}) \rangle \frac{B_{\parallel}^*}{m_i} dv_{\parallel} d\mu$$

$$\frac{dv_{\parallel}}{dt} = \{v_{\parallel}, H\} = -\frac{\mathbf{B}_{\parallel}^*}{B_{\parallel}^*} \cdot \nabla H$$

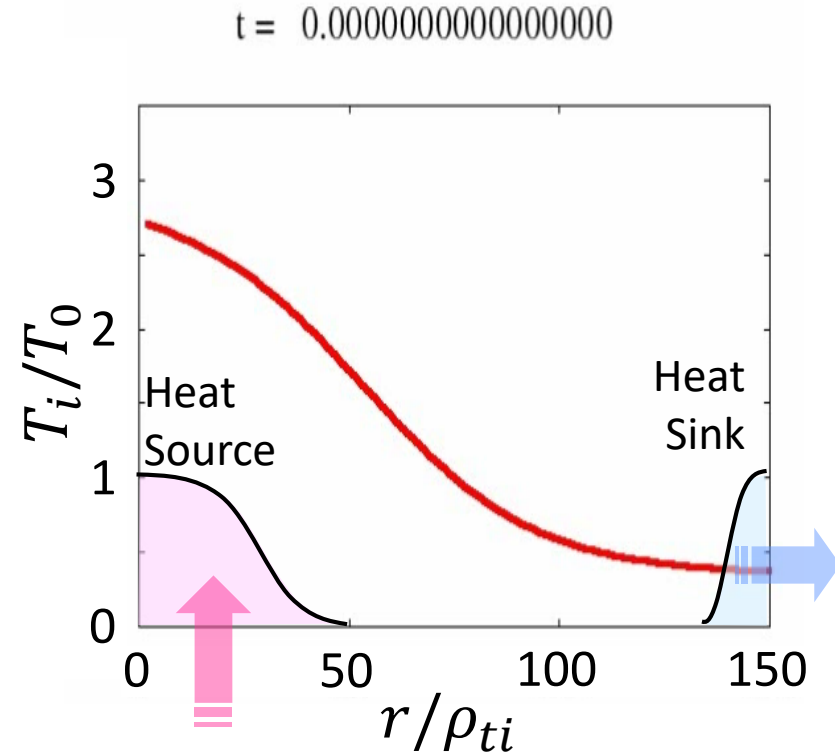
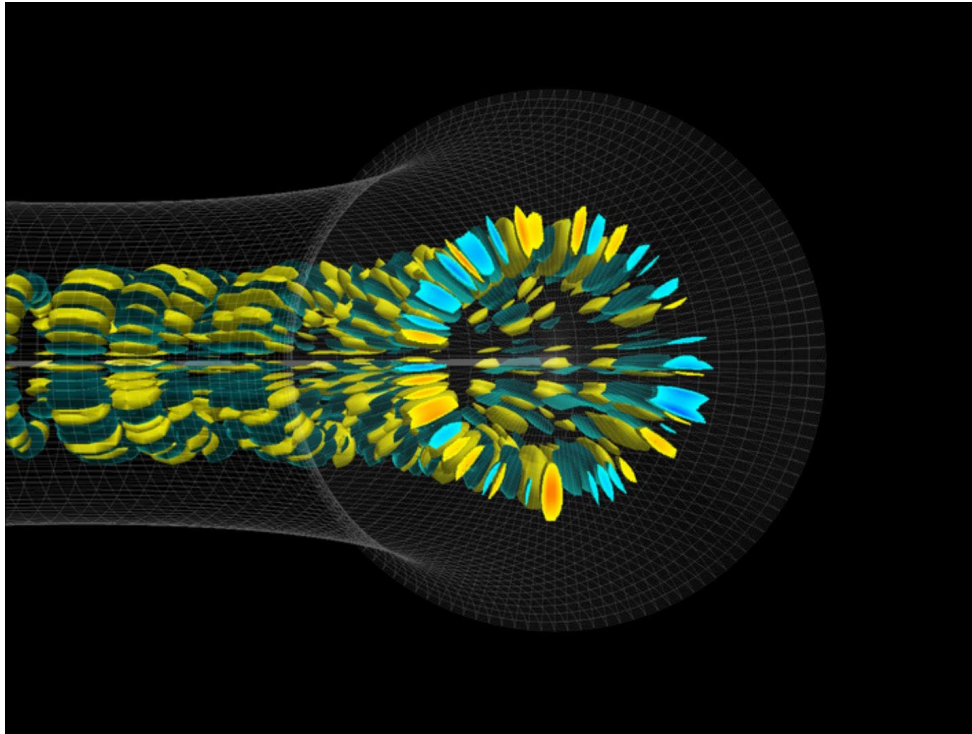
Adiabatic electron response

- ✓ Original *GKNET* is based on **full-*f* gyrokinetic model**, which trace turbulence and background profiles self-consistently [Imadera, IAEA-2014].
- ✓ We use the Morinishi scheme, which was developed for fluid simulation and introduced to rectangular gyrokinetic code, [Morinishi, JCP-2004, Idomura, JCP-2007] **to polar coordinate with new flux-conservative scheme**.
- ✓ **Field equation is solved in real space** (not k-space) and **full-order FLR effect** is taken into account by using 20 point average on gyro-ring [Obrejan, PFR-2015, CPC-2017].



# Flux-driven ITG Simulation by Original *GKNET*

Time evolution 3D electrostatic potential & 1D ion temperature



- ✓ By using external heat source/sink, we can evaluate quasi-steady state of turbulence, not decaying turbulence.
- ✓ Typical scale length  $L_T$  is tied to be globally constant, it does not largely change even if heat input power is increased (profile stiffness) -> L mode plasma

# Extension of *GKNET*

Gyrokinetic Vlasov equation (ion) & Gyrokinetic quasi-neutrality condition

$$\frac{\partial}{\partial t}(Jf) + \frac{\partial}{\partial \mathbf{R}} \left( J \frac{d\mathbf{R}}{dt} f \right) + \frac{\partial}{\partial v_{\parallel}} \left( J \frac{dv_{\parallel}}{dt} f \right) = S_{src} + S_{snk} + C_{coll}$$

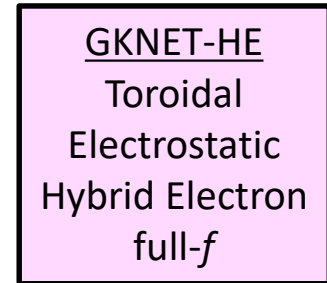
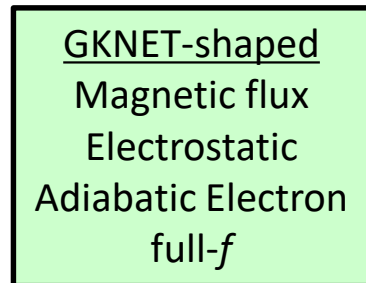
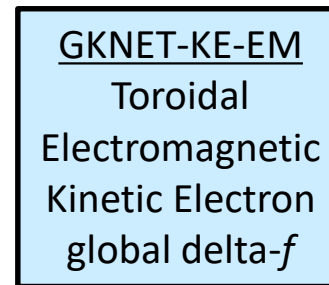
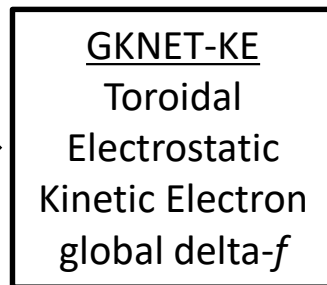
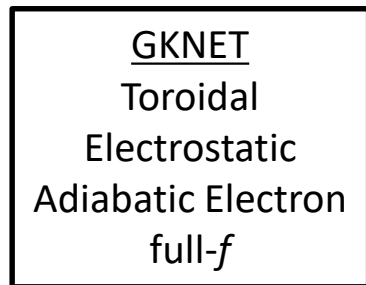
$$\frac{d\mathbf{R}}{dt} = \{\mathbf{R}, H\} = \frac{B_{\parallel}^*}{B_{\parallel}^*} \frac{\partial H}{\partial v_{\parallel}} + \frac{1}{B_{\parallel}^*} \mathbf{b} \times \nabla H$$

$$-\nabla_{\perp} \cdot \frac{\rho_{ti}^2}{\lambda_{Di}^2} \nabla_{\perp} \phi(\mathbf{R}) + \frac{1}{\lambda_{De}^2} [\phi(\mathbf{R}) - \langle \phi(\mathbf{R}) \rangle_f] = 4\pi e \iint \langle \delta f_i(\mathbf{R}) \rangle \frac{B_{\parallel}^*}{m_i} dv_{\parallel} d\mu$$

$$\frac{dv_{\parallel}}{dt} = \{v_{\parallel}, H\} = -\frac{B_{\parallel}^*}{B_{\parallel}^*} \cdot \nabla H$$

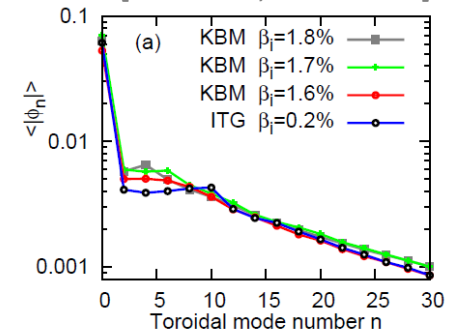
Adiabatic electron response

[Imadera, FEC-2014]



[Imadera, PFR-2020]

[Ishizawa, PoP-2019]

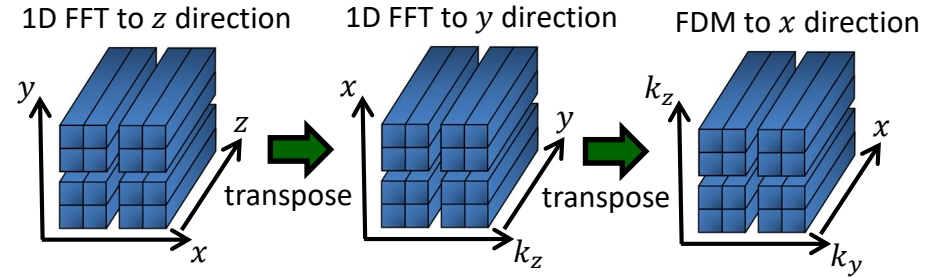


- ✓ Original *GKNET* is a full-*f* electrostatic gyrokinetic code with adiabatic electron, which is extended to
- (1) Magnetic flux coordinate version
  - (2) Electromagnetic delta-*f* version
  - (3) Hybrid Electron full-*f* version

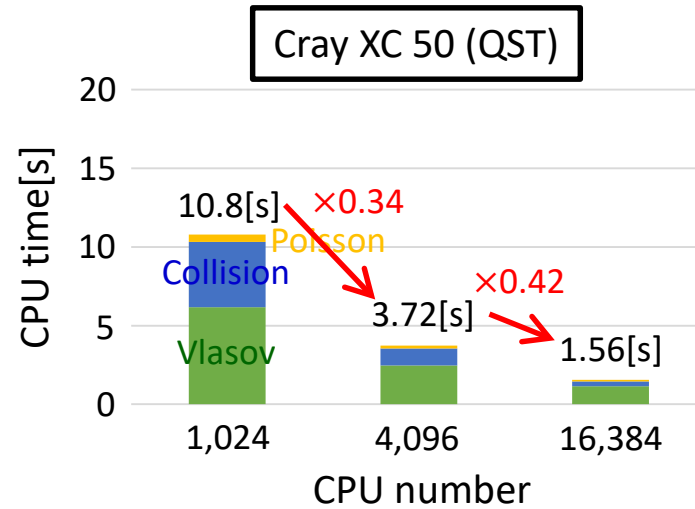
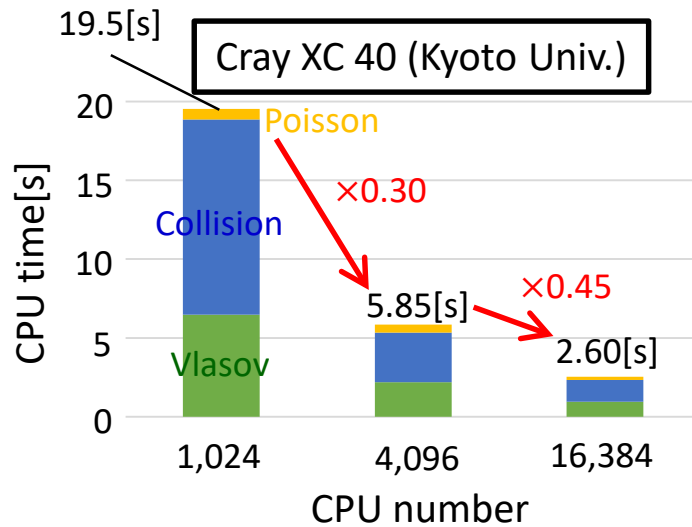
# Parallelization of *GKNET*

## High-efficient parallelization technique

✓ High efficient 2D FFT is installed by utilizing 1D FFT and MPI\_ALLtoALL transpose technique with the aid of communication and computation hiding optimization.



## Parallelization rate on super computer

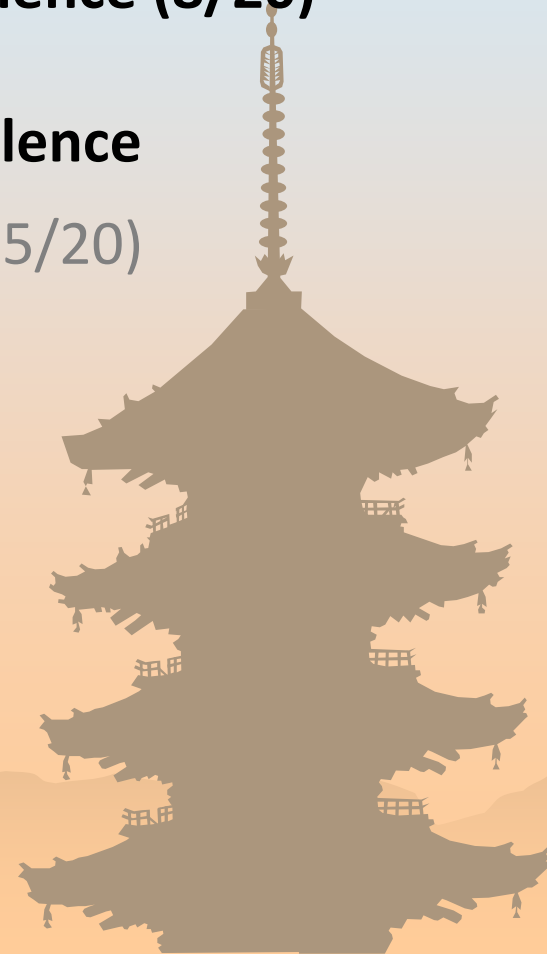


	Processor	Flops/Node	Memory/Node	Network
Cray XC40 (Kyoto Univ.)	Intel Xeon Phi 7250 (68[core])	3.046[Tflops]	16+96[GB]	Aries (Dragonfly)
Cray XC50 (QST)	Intel Xeon Gold 6148 (20[core] × 2)	3.072[Tflops]	192[GB]	InfiniBand EDR (Dragonfly)



# Contents

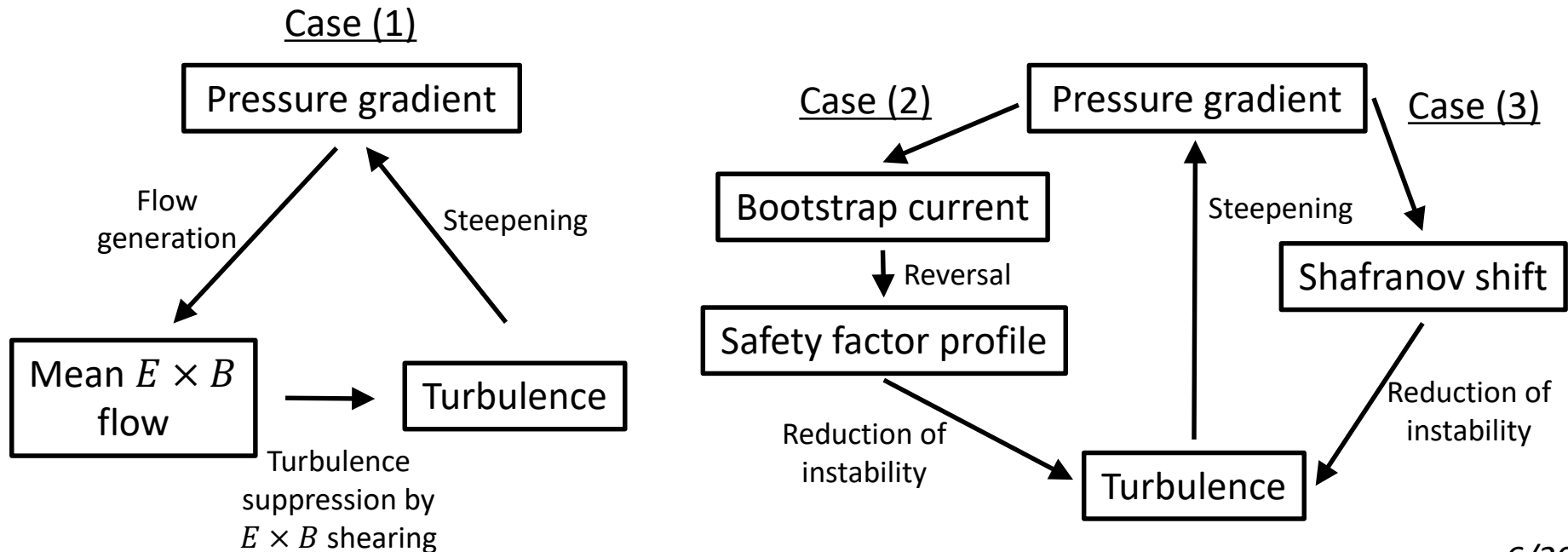
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# Background: Possible Mechanism of ITB Formation

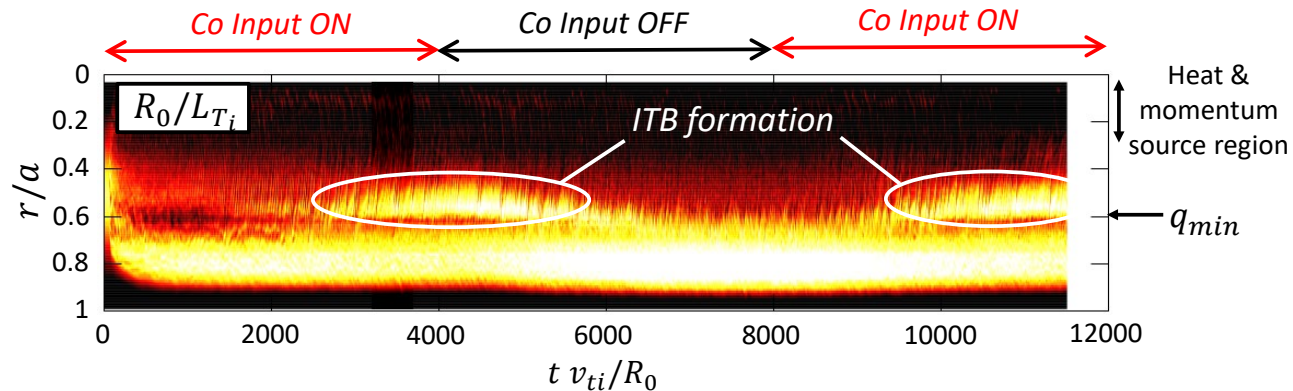
- ✓ Internal Transport barrier (ITB) has a crucial key to achieve a high-performance plasma confinement.
- ✓ Some possible mechanism for ITB formation are proposed [Ida, PPCF-2018] as
  - (1) Positive feedback loop via  $E \times B$  mean flow [Sakamoto, NF-2004] [Yu, NF-2016]
  - (2) Positive feedback loop via safety factor profile (BS current) [Eriksson, PRL-2002]
  - (3) Positive feedback loop via Shafranov shift + EM stabilization [Staebler, NF-2018]



# Motivation of This Research

- ✓ By our full- $f$  gyrokinetic code *GKNET*, we found that **momentum injection can change mean  $E \times B$  flow through the radial force balance, which can break the ballooning symmetry of turbulence, leading to ITB formation.** [Imadera, IAEA-2016]

- ✓ Such a mechanism can also benefit the ITB formation around  $q_{min}$  surface in reversed magnetic shear plasma.



- ✓ However, in our previous study based on the original *GKNET* with adiabatic electron, **enough large co-momentum injection is required for ITB formation in flux-driven ITG turbulence.** In addition, some experiments indicate the importance of counter-intrinsic rotation. [Sakamoto, NF-2001]

- ✓ In this study, **we have introduced hybrid kinetic electron model** [Lanti, JP-2018] **and investigated spontaneous ITB formation in flux-driven ITG/TEM turbulence.**

GKNET-HE  
Toroidal  
Electrostatic  
Hybrid Electron  
full- $f$

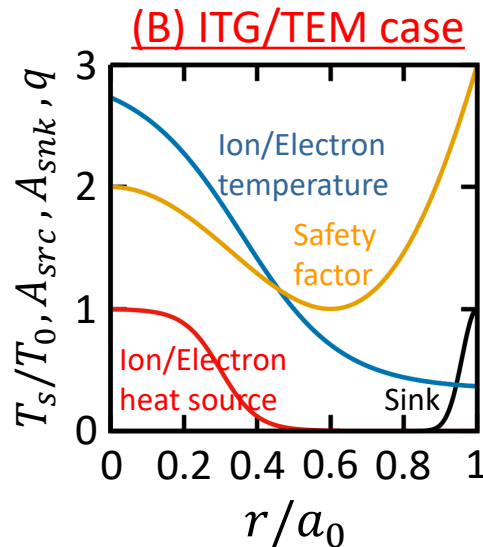
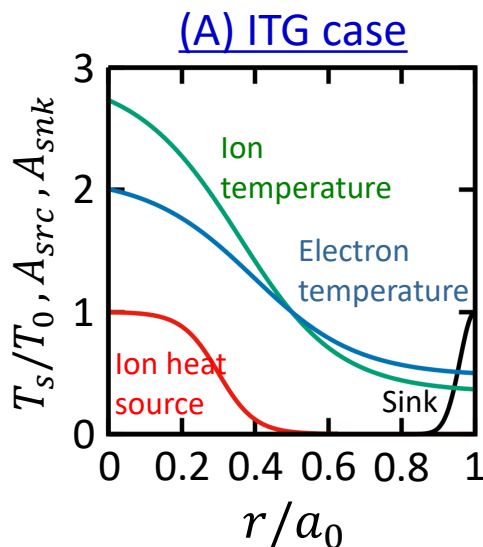
# ITB Formation in Flux-driven ITG/TEM Turbulence - 1

## Simulation condition

Parameter	Value
$a_0/\rho_i$	150
$a_0/R_0$	0.36
$(R_0/L_n)_{r=a_0/2}$	2.22
$(R_0/L_{Ti})_{r=a_0/2}$	10
$(R_0/L_{Te})_{r=a_0/2}$	(A) 6.92 (B) 10
$\Delta_r$	45
$\sqrt{m_i/m_e}$	10

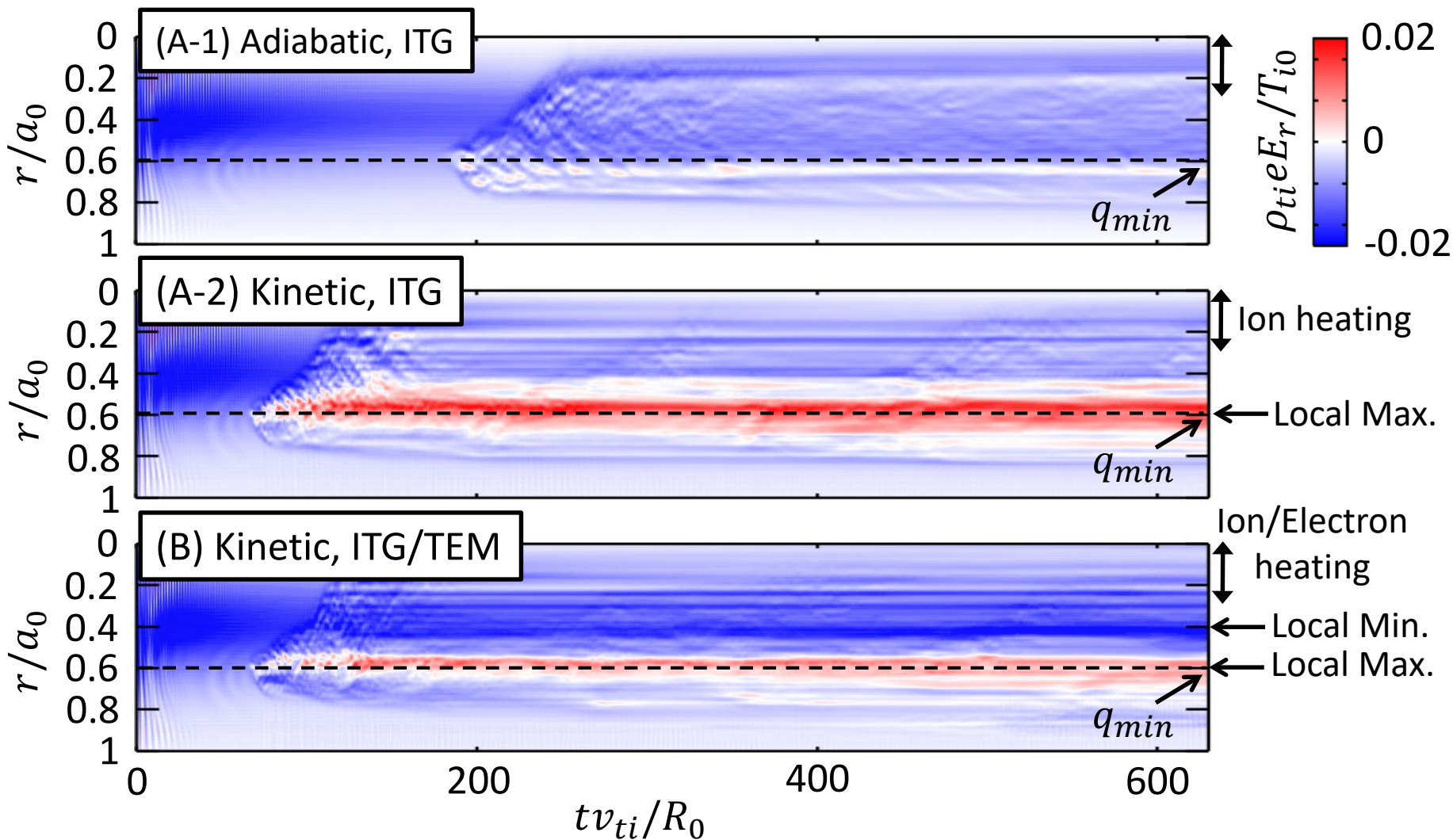
Parameter	Value
$\nu_i^*$	0.1
$\nu_e^*$	0.1
$\tau_{src,i}^{-1}$	0.02 -> 4[MW]
$\tau_{src,e}^{-1}$	(A) 0 -> 0[MW] (B) 0.02 -> 4[MW]
$\tau_{snk}^{-1}$	0.1/0.36

Parameter	Value
$N_r$	96
$N_\theta$	240
$N_\varphi$	48
$N_{v_{\parallel}}$	96
$N_\mu$	16
$\Delta t$	$3.125 \times 10^{-4}$



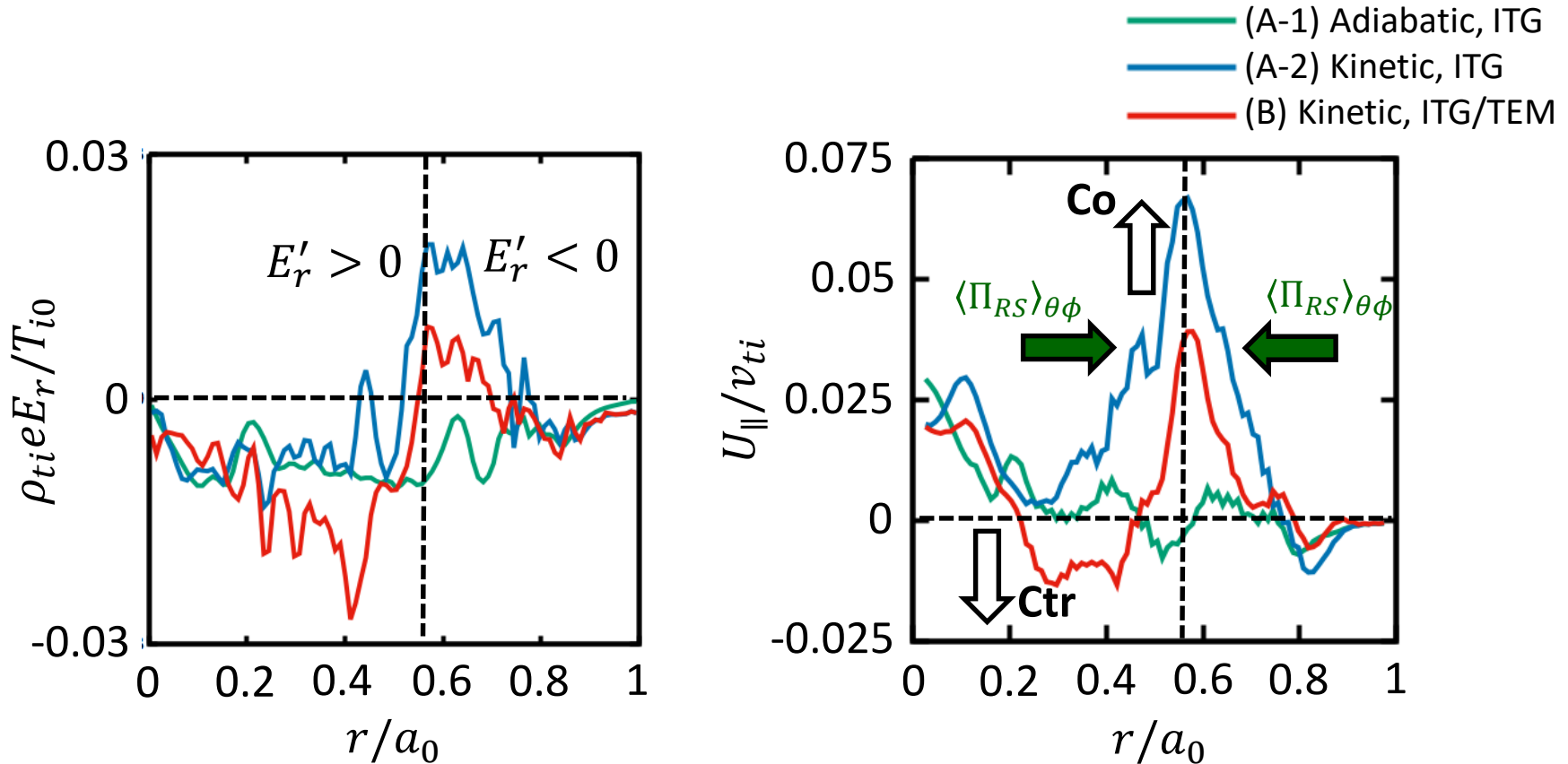
- ✓ We consider (A) ITG dominant and (B) ITG/TEM dominant cases.
- ✓ Safety factor profile is reversed, which local minimum is located at  $r = 0.6a_0$ .
- ✓ Only heat source is applied, which does not provide particle and momentum.

# ITB Formation in Flux-driven ITG/TEM Turbulence - 2



- ✓ Stable local maximum of mean  $E_r$  are formed near  $q_{min}$  surface only in kinetic electron cases.

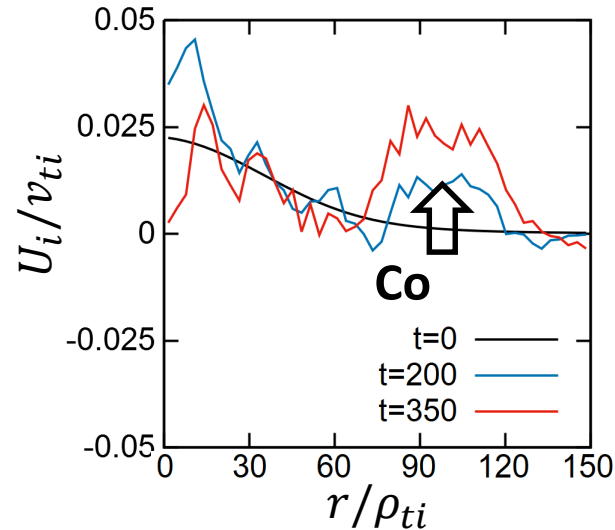
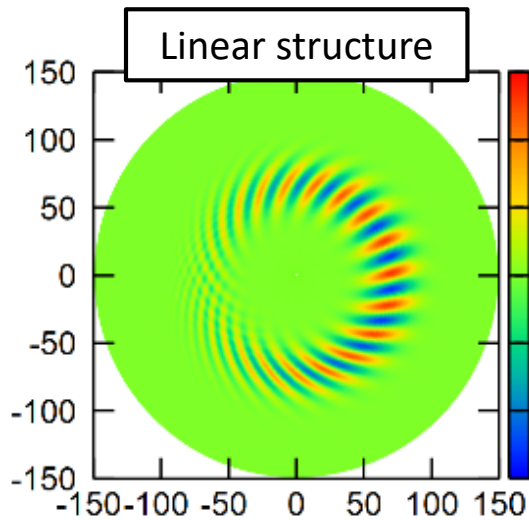
# ITB Formation in Flux-driven ITG/TEM Turbulence - 3



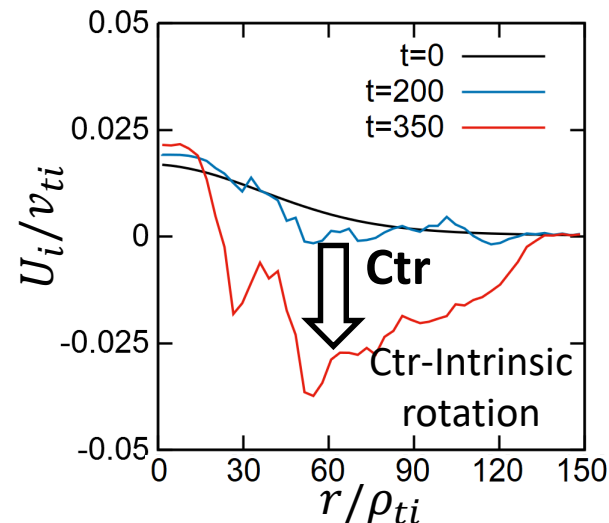
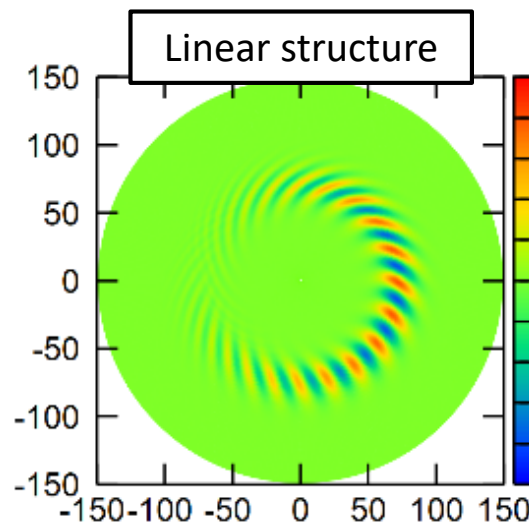
- ✓ Large co-rotation is driven around  $q_{min}$  surface in case (A-2) and (B).
- ✓ According to the momentum transport theory,  $\langle \Pi_{RS} \rangle_{\theta\phi} = \alpha I E_r' + \beta I' + \gamma \langle k_{\theta} k_{\phi} \phi_k^2 \rangle_{\theta\phi}$  [Kwon, NF-2012], the first and second terms can reduce momentum diffusion in this case, **which can keep the stable local maximum of mean  $E_r$**  through the radial force balance.
- ✓ Counter-rotation is also observed in negative magnetic shear region in case (B).

# What is the Origin of Co-/Counter-Rotation?

Decaying ITG turbulence in CBC case

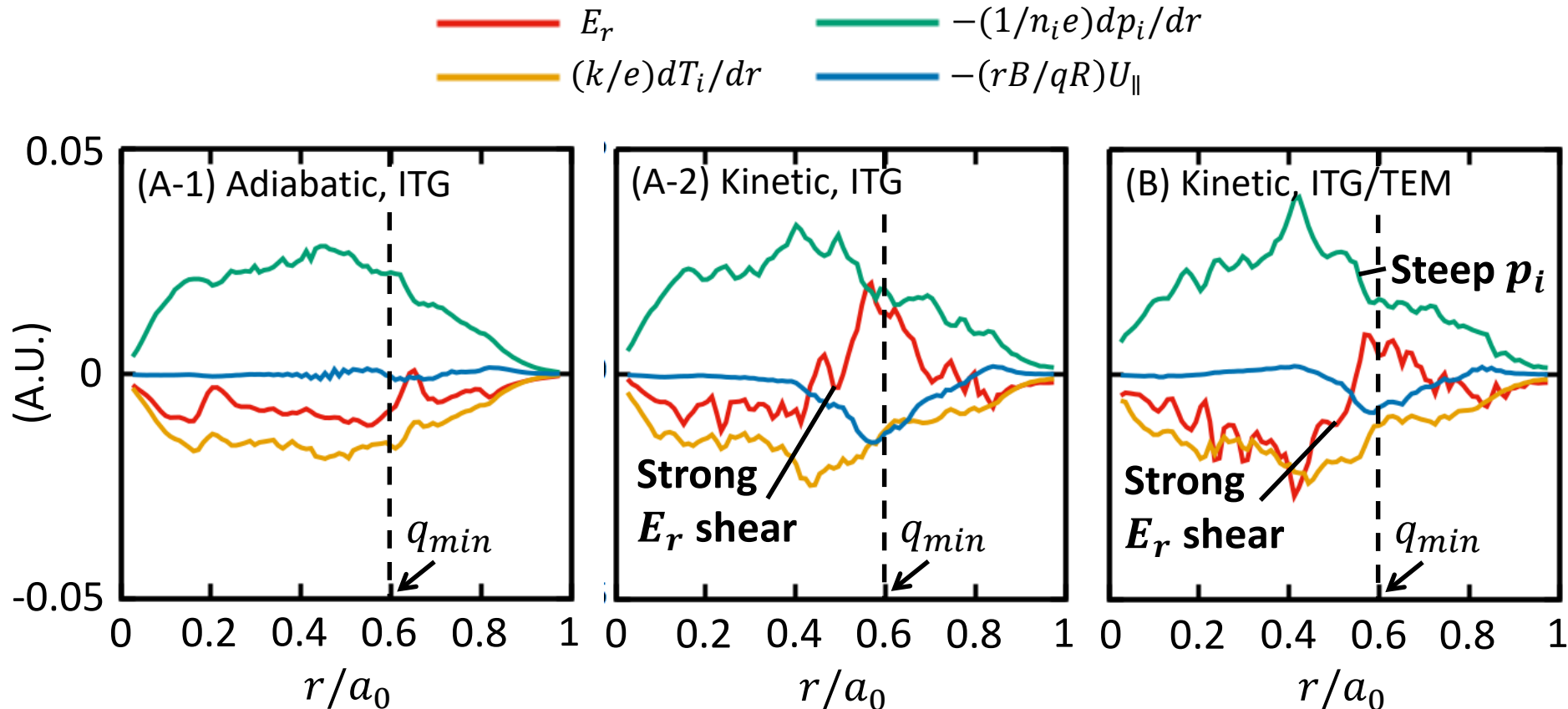


Decaying TEM turbulence in CBC case



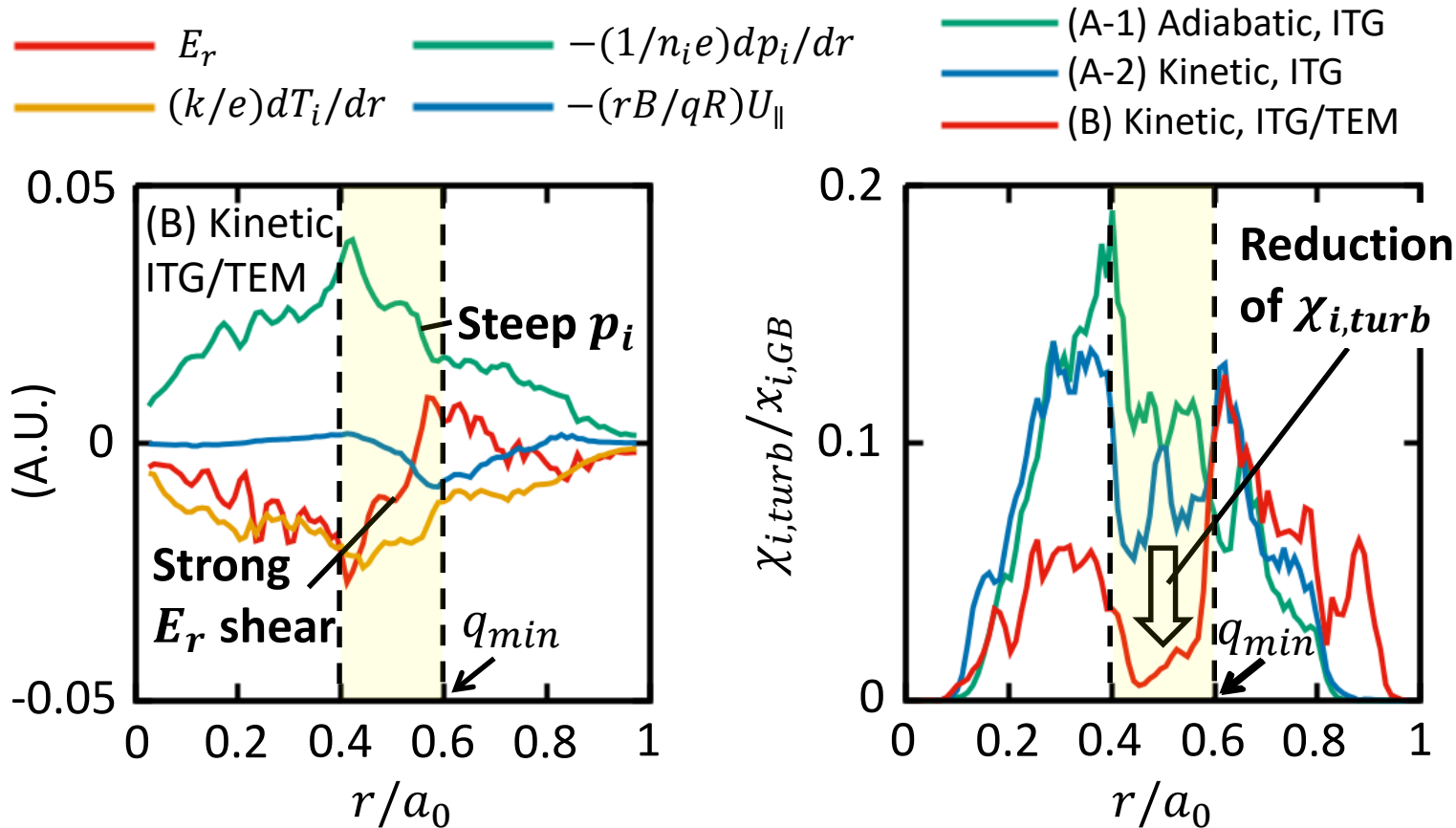
- ✓ The finite ballooning angle of the global mode structure arising from the profile shearing effect [Kishimoto, PPCF-1998] induces the residual stress part of momentum flux [Camenen, NF-2011].
- ✓ The sign of the ballooning angle between ITG and TEM turbulence is opposite (left figures) so that **the direction of intrinsic rotation is reversed.**
- ✓ The steep electron temperature gradient is considered to destabilize TEM in the negative magnetic shear region.

# ITB Formation in Flux-driven ITG/TEM Turbulence - 4



- ✓ In flux-driven ITG turbulence with kinetic electrons, the co-current toroidal rotation can balance with  $E_r$ , of which shear becomes strong just inside of  $q_{min}$  surface.
- ✓ On the other hand, in ITG/TEM turbulence with kinetic electrons,  $E_r$  is reversed in negative magnetic shear region, which makes its shear stronger and pressure gradient steeper.

# ITB Formation in Flux-driven ITG/TEM Turbulence - 5



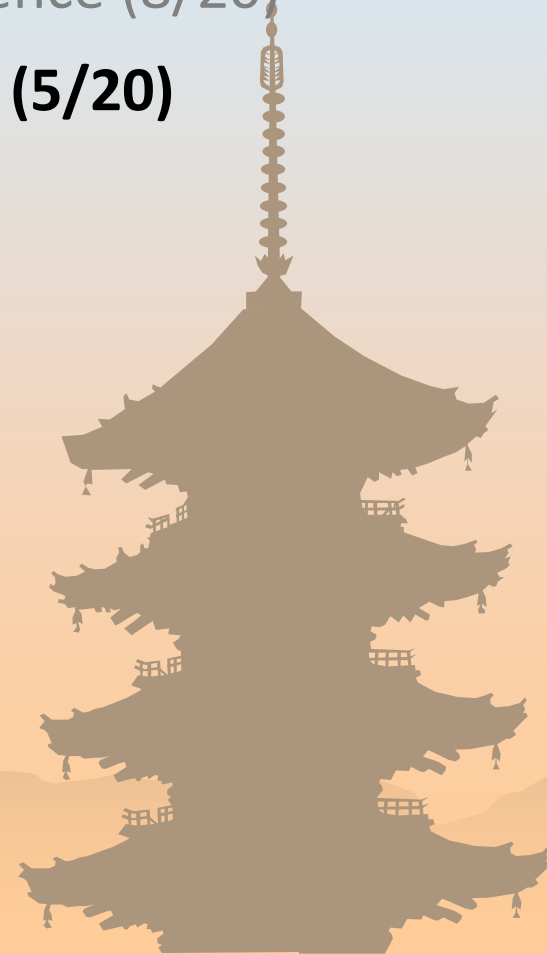
- ✓ As the result, ion turbulent thermal diffusivity in flux-driven ITG/TEM case spontaneously decreases to the neoclassical transport level among  $0.4a_0 < r < 0.6a_0$ , where  $E_r$  shear becomes steep.
- ✓ These results indicate that the co-existence of different modes can trigger the discontinuity near  $q_{min}$ , leading to the spontaneous ITB formation.





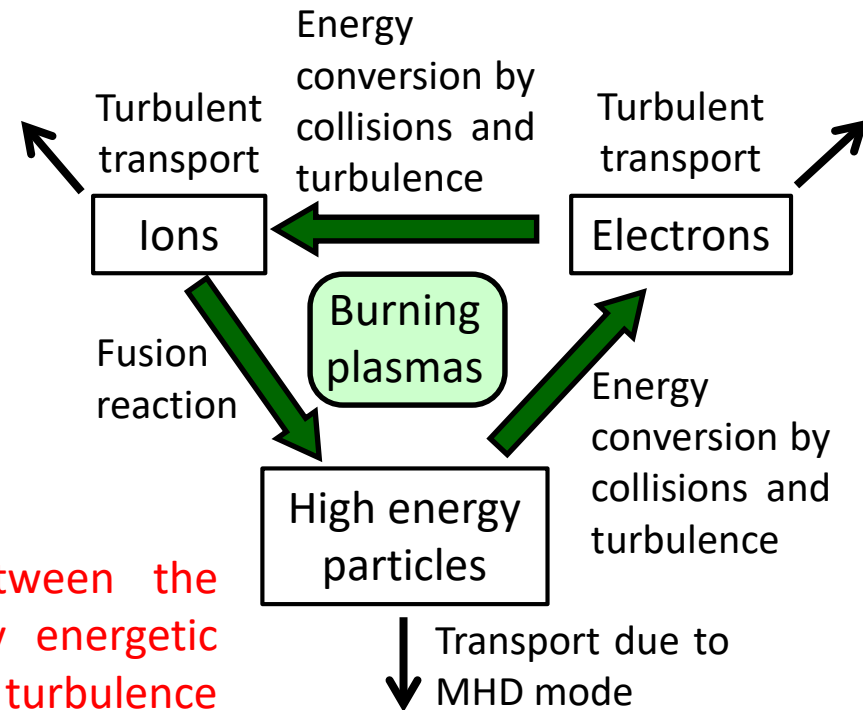
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  - \* Linear Analysis of TAE and KBM**
  - \* Nonlinear Analysis of TAE and KBM**
5. Summary & Future plans (2/20)

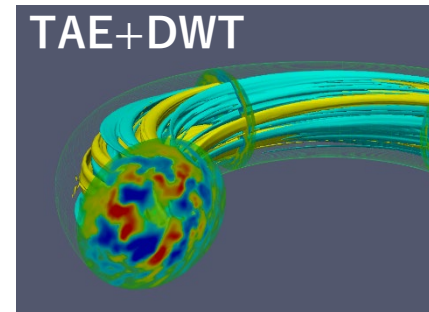
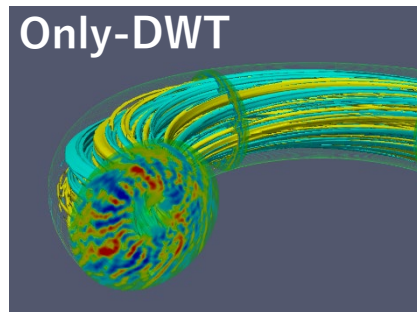


# Background & Motivation

- ✓ In order to realize high performance burning plasmas, it is necessary to reduce both energetic alpha-particle transport and bulk plasma transport simultaneously.
- ✓ Drift-wave turbulence and MHD modes driven by energetic-particles coexist in burning plasmas, thereby the interaction between them is expected to take place and lead to new transport phenomena.
- ✓ We investigate **nonlinear interactions between the toroidal Alfvén eigenmode (TAE) driven by energetic particles and electromagnetic drift-wave turbulence (KBM) by using the global delta- $f$  electromagnetic version of GKENT.**

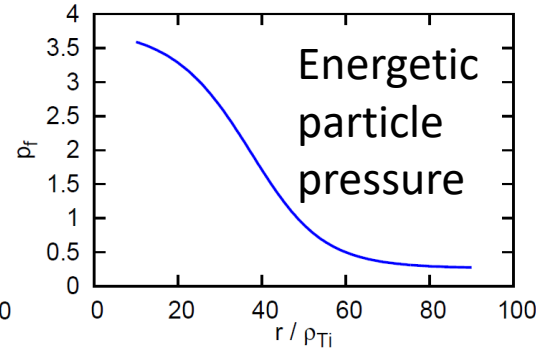
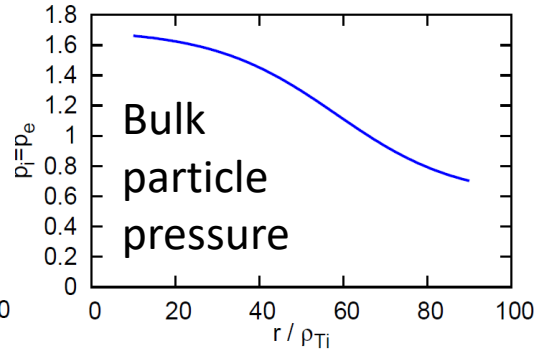
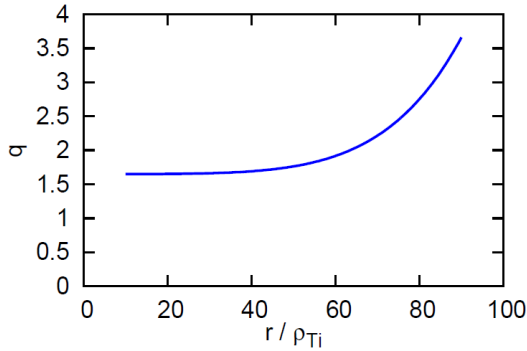


GKNET-KE-EM  
Toroidal  
Electromagnetic  
Kinetic Electron  
global delta- $f$



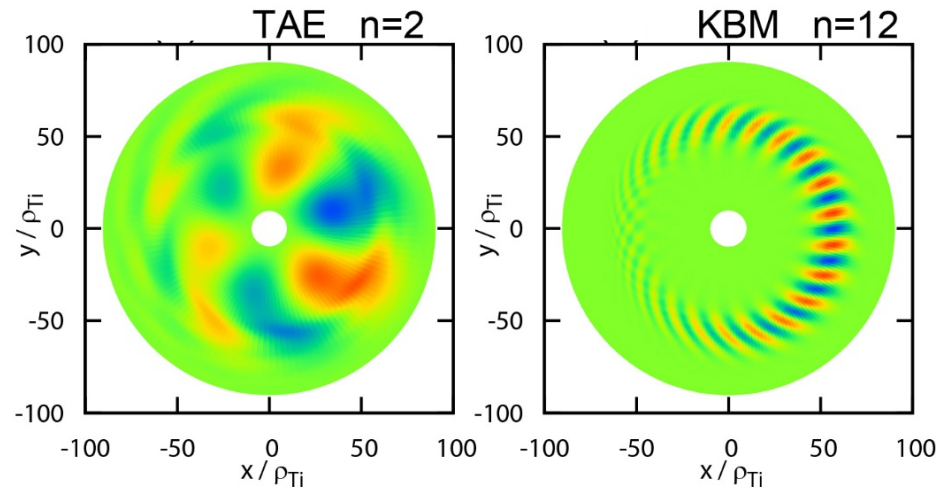
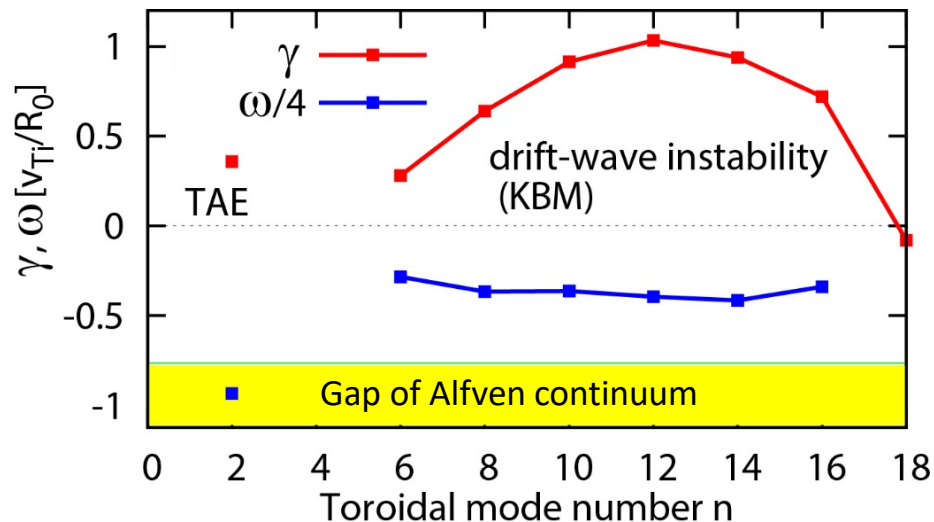
# Simulation Set-up & Linear Analysis of TAE & KBM

## Simulation set-up



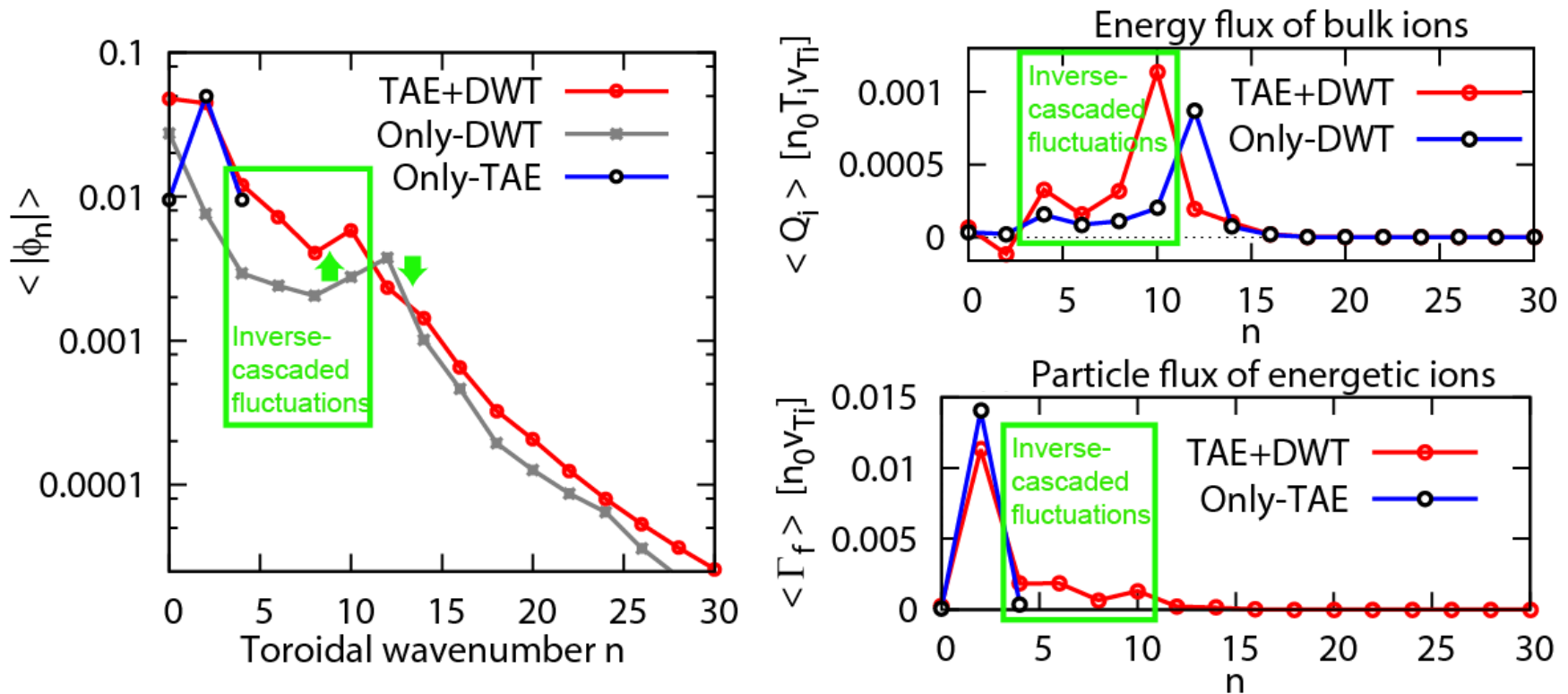
$\beta = 1.28\%$   
 $\rho_* = 1/100$   
 $M_i/M_e = 100$   
 $T_f/T_i = 25$   
 $n_f/n_i = 0.025$

## Linear dispersion relation & eigenfunctions



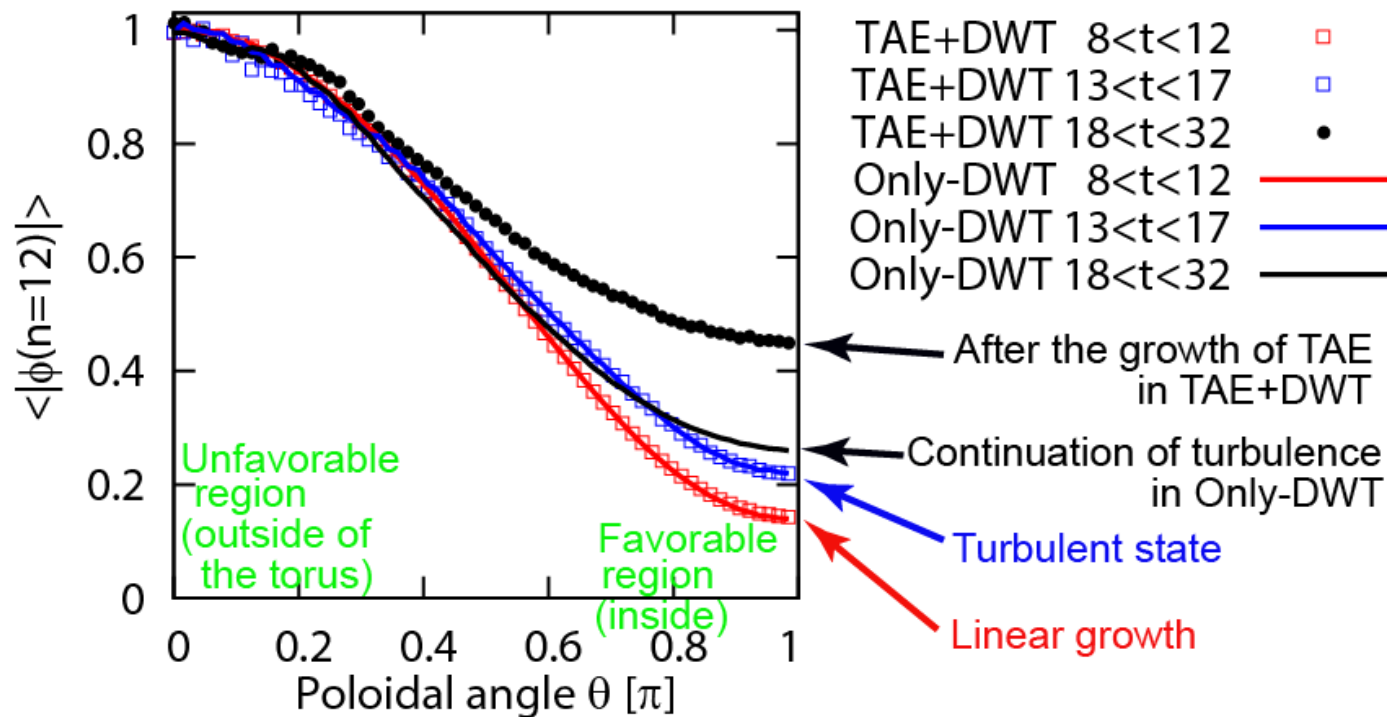
- ✓ The plasma is unstable against a TAE at low toroidal mode number  $n = 2$  while kinetic ballooning mode (KBM) is unstable at high toroidal mode number  $n > 6$ .

# Nonlinear Analysis of TAE & KBM -1



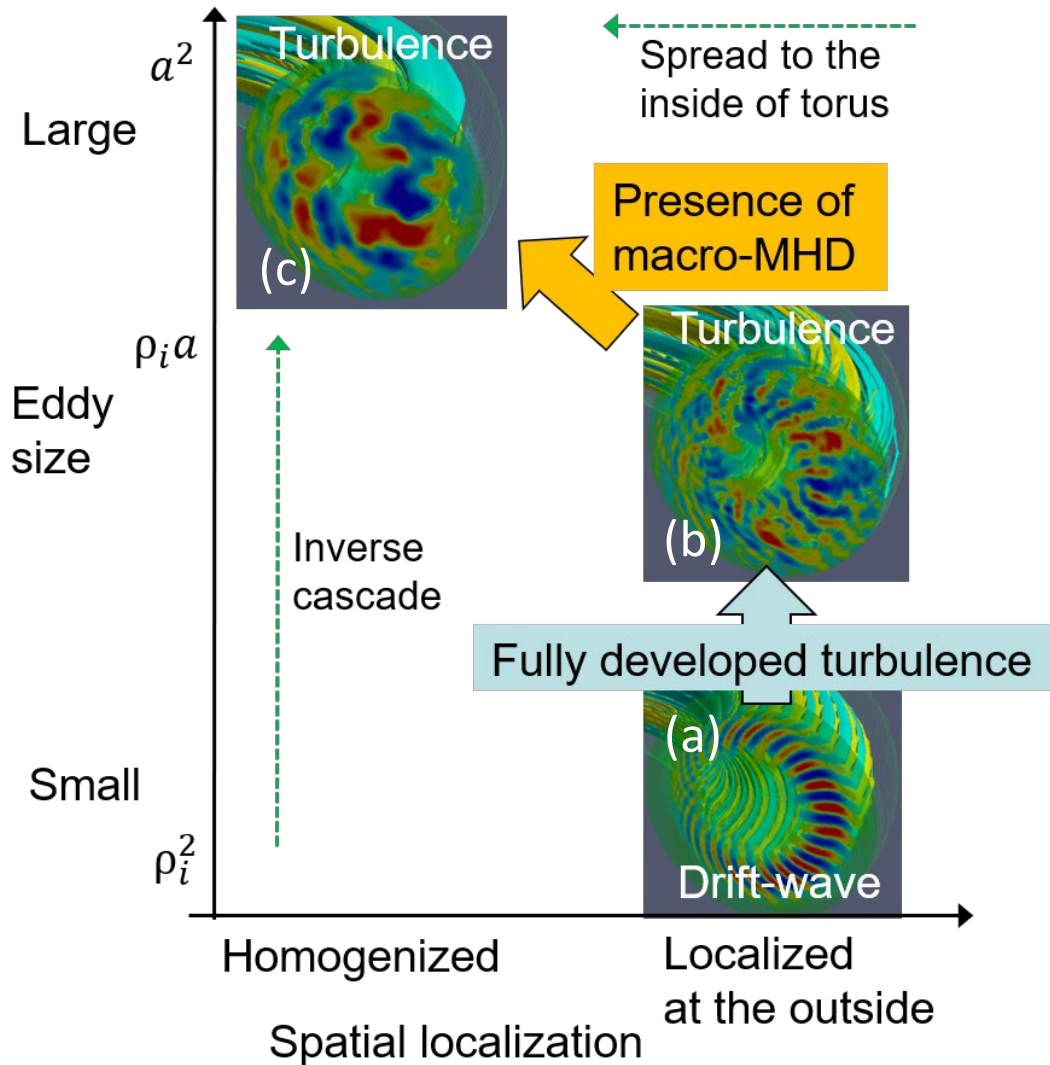
- ✓ TAE suppresses the most unstable drift-wave mode but enhances a smaller toroidal wavenumber mode, causing the inverse cascade.
- ✓ Due to the inverse-cascaded fluctuations, the energy flux of bulk ions  $Q_i$  in TAE+DWT is enhanced at middle wavenumbers ( $4 < n < 10$ ).
- ✓ The interaction slightly suppresses the particle flux of energetic ions  $\Gamma_f$  at  $n < 2$  but enhances  $\Gamma_f$  by the inverse-cascaded fluctuations.

# Nonlinear Analysis of TAE & KBM -2



- ✓ Before the growth of the TAE, the drift-wave turbulence is poloidally localized in the unfavorable curvature region.
- ✓ Then, after the development of the TAE, the turbulence spreads to the favorable curvature region because of the global structure of the TAE, suppressing the most unstable drift-wave mode through the geometrical damping effect.

# Nonlinear Analysis of TAE & KBM -3

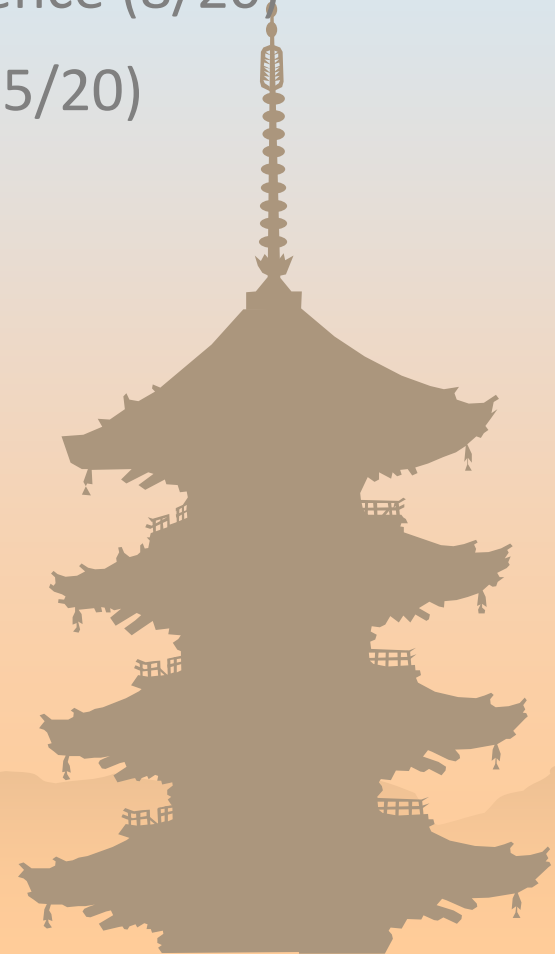


- ✓ The drift-wave grows at the outside of the torus at the frame (a).
- ✓ Then becomes turbulence with the inverse cascade at the frame (b).
- ✓ The nonlinear mode coupling of turbulence with the macro-scale MHD instability, by contrast, transfers the energy to the homogenized and large-scale structure at the frame (c).



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- 5. Summary & Future plans (2/20)**



# Summary & Future Plans - 1

## Summary: Interaction between TAE and KBM turbulence

- ✓ We have performed the global electromagnetic simulation to study multi-scale nonlinear interactions between micro-scale drift-wave turbulence and the toroidal Alfvén eigenmode, which is a macro-scale MHD instability driven by energetic particles.
- ✓ As a result of the interactions, the TAE transfers the energy of turbulence from high  $n$  modes to low  $n$  modes, causing the inverse cascade.
- ✓ The inverse-cascaded fluctuations enhance both the bulk ion energy transport and fast ion particle transport.
- ✓ Before the growth of the TAE, the drift-wave turbulence is poloidally localized in the unfavorable curvature region. Then, after the development of the TAE, the turbulence spreads to the favorable curvature region, suppressing the most unstable drift-wave mode through the geometrical damping effect.



# Summary & Future Plans - 2

## Summary: Spontaneous ITB formation in ITG/TEM turbulence

- ✓ We have performed the flux-driven ITG/TEM simulation in reversed magnetic shear configuration by using hybrid kinetic electron model.
- ✓ In the presence of both ion and electron heating, **a counter-intrinsic rotation by TEM turbulence is driven in negative magnetic shear region, leading to stronger  $E_r$  shear and the resultant spontaneous larger reduction of ion turbulent thermal diffusivity.**
- ✓ An increase of counter intrinsic rotation in the narrow region of the ITB located just inside of  $q_{min}$  is also observed in JT-60U reversed magnetic shear discharge with balanced momentum injection [Sakamoto, NF-2001]. -> Qualitative agreement!

## Future Plans

- ✓ By reflecting bootstrap current and Shafranov shift effects to the analytical magnetic equilibrium [Imadera, PFR-2020] in time, we can take them into account, which can help us to understand the overall positive feedback loop.

GKNET-shaped  
Magnetic flux  
Electrostatic  
Adiabatic Electron  
full- $f$

